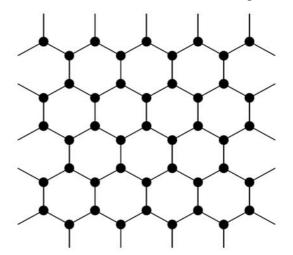
Problem 12: Band structure of graphene

[4 points]

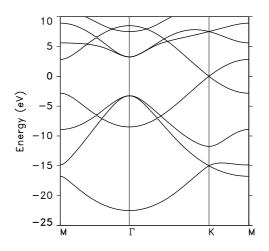
Graphene consists of a layer of carbon atoms that are arranged in a hexagonal structure. The lattice vectors are given by $\vec{a}_1 = (1, 0) a$ and $\vec{a}_2 = (-1, \sqrt{3}) \frac{a}{2}$.



- a) Give the primitive vectors \vec{b}_1 and \vec{b}_2 of the reciprocal lattice and construct the first Brillouin zone.
- b) Use the empirical tight-binding method with one p_z orbital per atom to calculate the band structure $E_n(k_x, k_y)$ of graphen. The hopping term t only acts between nearest neighbours.
- c) Plot the band structure for $E_0 = 0$ eV and t = -2.828 eV along the high-symmetry lines from Γ to K and from K to M.

$$K:$$
 $\left(\frac{2}{3}, 0\right) \frac{2\pi}{a}, \qquad M:$ $\left(\frac{1}{2}, \frac{-1}{2\sqrt{3}}\right) \frac{2\pi}{a}.$

d) The figure shows the band structure of graphen resulting from a calculation with s, p_x , p_y and p_z orbitals. Compare your result with this band structure.



e) Show that in the vicinity of K, i. e. for $\vec{k} = K + \vec{q}$ (with small \vec{q}), the tight-binding method yields isotropic bands with linear dispersion $E_{\pm}(\vec{k}) \approx \pm v \cdot |\vec{q}|$ ("Dirac cones of graphene").

Problem 13: Density of states for a linear chain

The band structure of a linear chain with one s-like orbital per atom is given within the framework of the empirical tight-binding method by $E_k = E_0 + 2t \cos(ka)$. Calculate the density of states $N(E) = \sum_{k} \delta(E - E_k)$ of the chain. *Hint*: Substitute the sum by an integral.

Problem 14: Homogeneous electron gas

Consider N interacting electrons in a volume Ω with a neutralizing background of a constant positive density $\rho_{\text{nucl}} = e n_{\text{nucl}} = e \frac{N}{\Omega}$. Within the Hartree-Fock approximation, the one-particle wave functions $\Psi_{\vec{k},\sigma}(\vec{r})$ are given by the solutions of

$$\left(-\frac{\hbar^2 \nabla^2}{2 m} + V_{EN}(\vec{r}) + V_{Coul}(\vec{r})\right) \psi_{\vec{k},\sigma}(\vec{r}) \\ -\sum_{\sigma'=-\frac{1}{2}}^{1/2} \sum_{\vec{k}'} \delta_{\sigma,\sigma'} e^2 \int_{\Omega} \frac{\Psi_{\vec{k}'}^*(\vec{r}',\sigma') \Psi_{\vec{k}}(\vec{r}',\sigma)}{|\vec{r}-\vec{r}'|} d^3 r' \Psi_{\vec{k}',\sigma'}(\vec{r}) = \varepsilon_{\vec{k},\sigma} \psi_{\vec{k}}(\vec{r},\sigma)$$

with

$$V_{EN}(\vec{r}) = -\frac{N}{\Omega} e^2 \int_{\Omega} \frac{1}{|\vec{r} - \vec{r'}|} d^3 r'$$

and

$$V_{\text{Coul}}(\vec{r}) = e^2 \int_{\Omega} \frac{n(\vec{r}')}{|\vec{r} - \vec{r}'|} d^3 r' , \qquad n(\vec{r}) = \sum_{\sigma} \sum_{\vec{k}} |\Psi_{\vec{k}}(\vec{r}, \sigma)|^2 .$$

The sums over \vec{k} and \vec{k}' include all occupied states, i. e. $|\vec{k}| \leq k_F$, $|\vec{k}'| \leq k_F$.

a) Show that the Hartree-Fock equations of this system are solved by plane waves

$$\Psi_{\vec{k}}(\vec{r},\sigma) = \frac{1}{\sqrt{\Omega}} e^{i\vec{k}\cdot\vec{r}} \quad \text{for both spins} \quad \left(\sigma = \pm \frac{1}{2}\right) \ .$$

Hint: Convince yourselves that V_{EN} is compensated by V_{Coul} .

b) Calculate the eigenvalues $\varepsilon_{\vec{k},\sigma}$. To this end, convert the sum over \vec{k}' into an integral. Useful integral:

$$\int x \ln \left| \frac{x+a}{x-a} \right| \, dx = \frac{1}{2} \left(x^2 - a^2 \right) \ln \left| \frac{x+a}{x-a} \right| + ax \, .$$

- c) Plot $\varepsilon_{\vec{k},\sigma}$ and discus its behaviour at $k = k_F$.
- d) The eigenvalues

$$\varepsilon_{\vec{k},\sigma} = \frac{\hbar^2}{2m}k^2 + \Sigma^{\text{ex}}(\vec{k},\sigma)$$

contain the self energy $\Sigma^{\text{ex}}(\vec{k}, \sigma)$ (here: only exchange). Within the Hartree-Fock approximation, it contributes the exchange energy

$$E^{\text{ex}} = \frac{1}{2} \sum_{\substack{\vec{k}, \sigma \\ k < k_F}} \Sigma^{\text{ex}}(\vec{k}, \sigma)$$

to the total energy. Show that $E^{\text{ex}} = -N \cdot \frac{3}{4} \left(\frac{3}{\pi}\right)^{\frac{1}{3}} \cdot n^{\frac{1}{3}}$ with $n = \frac{N}{\Omega}$.

[2 points]

[6 points]