As discussed in the lecture, electrostatic response can be expressed in terms of three potentials  $\delta\varphi_{\text{ext}}(\vec{r})$ ,  $\delta\varphi_{\text{ind}}(\vec{r})$ , and  $\delta\varphi(\vec{r})$ .  $\delta\varphi_{\text{ext}}$  is provided "from outside",  $\delta\varphi_{\text{ind}}$  constitutes the system's response, and  $\delta\varphi = \delta\varphi_{\text{ext}} + \delta\varphi_{\text{ind}}$  is the resulting, "total" potential. The charge-density response  $\delta\rho_{\text{ind}}(\vec{r})$  can be expressed as  $\delta\rho_{\text{ind}} = \chi_0 \circ \delta\varphi$  (with charge susceptibility  $\chi_0$ ), resulting in  $\delta\varphi_{\text{ind}} = V \circ \delta\rho_{\text{ind}}$  (with Coulomb interaction  $V(\vec{r} - \vec{r'}) = \frac{1}{|\vec{r} - \vec{r'}|}$ ).

(SS 2016)

For a homogeneous system, Fourier transform

$$f(\vec{r}) = \frac{1}{(2\pi)^3} \int f(\vec{q}) e^{i\vec{q}\vec{r}} d^3 q , \qquad f(\vec{q}) = \int f(\vec{r}) e^{-i\vec{q}\vec{r}} d^3 r$$

turns the convolutions into multiplications, i.e.

$$\delta\varphi_{\mathrm{ind}}\left(\vec{q}\right) = V\left(\vec{q}\right) \cdot \chi_{0}\left(\vec{q}\right) \cdot \delta\varphi\left(\vec{q}\right) \,.$$

By defining  $\varepsilon(\vec{r}, \vec{r}')$  such that  $\varepsilon \circ \delta \varphi = \delta \varphi_{\text{ext}}$  one arrives at  $\varepsilon = \delta(\vec{r} - \vec{r}') - V \circ \chi_0$  (or, for a homogeneous system,  $\varepsilon(\vec{q}) = 1 - V(\vec{q}) \cdot \chi_0(\vec{q})$ ).

Alternatively, the charge-density response can be expressed as  $\delta \rho_{\text{ind}} = \chi \circ \delta \varphi_{\text{ext}}$  (with  $\chi \neq \chi_0$ !), finally resulting (for a homogeneous system) in  $\frac{1}{\varepsilon(\vec{q})} = 1 + V(\vec{q}) \chi(\vec{q})$ .

The static dielectric function of a homogeneous, infinitely large semiconductor may be given by

$$\varepsilon(q) = 1 + \frac{1}{(\varepsilon_0 - 1)^{-1} + q^2/\bar{q}^2}$$

with  $\varepsilon_0 = \varepsilon (q = 0)$  the static dielectric constant and  $\bar{q}$  being some characteristic wave number (e.g.,  $q_{\rm TF}$  in a metal). Consider an external point charge Q at  $\vec{r} = 0$  (e.g. due to a defect atom), resulting in  $\delta \varphi_{\rm ext} (\vec{r}) = \frac{Q}{r}$ .

- a) Calculate and plot the induced charge density  $\delta \rho_{\text{ind}}(\vec{r})$ . Show that the entire induced charge amounts to  $(\varepsilon_0^{-1} 1)Q$ .
- b) Calculate and plot  $\delta \varphi_{\text{ind}}(\vec{r})$  and  $\delta \varphi(\vec{r})$ . Show that for  $r \to \infty$ ,

$$\delta \varphi \left( \vec{r} \right) \longrightarrow \frac{Q}{\varepsilon_0 r} .$$

c) What do you get for a simple metal?

## Problem 4: Commutator relations for fermions

a) The anticommutator of the operators  $\hat{A}$  and  $\hat{B}$  is given by

$$[\hat{A}, \,\hat{B}]_{+} = \hat{A}\,\hat{B} + \hat{B}\,\hat{A}$$

Show that the following relation holds for an additional operator D

$$[\hat{A}, \,\hat{B}\,\hat{D}]_{-} = [\hat{A}, \,\hat{B}]_{+}\,\hat{D} - \hat{B}\,[\hat{A}, \,\hat{D}]_{+}$$

(3 points)

## (5 points)

- b) The creation and annihilation operators  $\hat{c}_j^+$  and  $c_j$  have been introduced in the lecture. Use the anticommutator relations of these operators to calculate the following commutator  $[\hat{A}, \hat{B}]_- = \hat{A}\hat{B} - \hat{B}\hat{A}$ 
  - i)

 $[\hat{n}_j, \, \hat{c}_k]_-$  and  $[\hat{n}_j, \, \hat{c}_k^+]_-$  with  $\hat{n}_j = \hat{c}_j^+ \, \hat{c}_j$ .

ii)

$$[\hat{c}_i^+ \, \hat{c}_j, \, \hat{c}_l^+ \, \hat{c}_m]_- \, = \, \alpha \, \cdot \, \hat{c}_i^+ \, \hat{c}_m \, + \, \beta \, \hat{c}_l^+ \, \hat{c}_j \, \, .$$

Calculate  $\alpha$  and  $\beta$ .

iii)

$$[\hat{c}_{i}^{+} \hat{c}_{j} \hat{c}_{l}^{+} \hat{c}_{m}, \hat{c}_{n}^{+} \hat{c}_{p}]_{-} = (\alpha \cdot \hat{c}_{i}^{+} \hat{c}_{p} + \beta \cdot \hat{c}_{n}^{+} \hat{c}_{j}) \hat{c}_{l}^{+} \hat{c}_{m} + \hat{c}_{i}^{+} \hat{c}_{j} (\gamma \cdot \hat{c}_{l}^{+} \hat{c}_{p} + \zeta \cdot \hat{c}_{n}^{+} \hat{c}_{m}) .$$

 $[\hat{A}\hat{B},\hat{D}]_{-} = [\hat{A},\hat{D}]_{-}\hat{B} + \hat{A}[\hat{B},\hat{D}]_{-}$ 

Calculate  $\alpha$ ,  $\beta$ ,  $\gamma$  and  $\zeta$ .

Useful relation:

The eigenstates of the Hamilton operator

$$\hat{H} = \sum_{j=1}^{\infty} \varepsilon_j \hat{c}_j^+ \hat{c}_j$$
 have the form  $|\phi\rangle = \prod_{j=1}^{\infty} (\hat{c}_j^+)^{n_j} |0\rangle$ .

a) Calculate  $\hat{n}_l |\phi\rangle$  with  $\hat{n}_l = \hat{c}_l^+ \hat{c}_l$ .

b) Determine the expectation values

a) 
$$\langle \phi | \hat{c}_l^+ \hat{c}_m | \phi \rangle$$
,  
b)  $\langle \phi | \hat{c}_i^+ \hat{c}_l^+ \hat{c}_k \hat{c}_m | \phi \rangle$ 

Use the occupation numbers  $n_k$  and  $n_m$  to represent your result.

## Problem 6: Two-level system

The Hamilton operator of a system with two spin degenerate energy levels  $\varepsilon_a$  and  $\varepsilon_b$  has the form

$$\hat{H} = \varepsilon_a \left( \hat{c}^+_{a\uparrow} \hat{c}_{a\uparrow} + \hat{c}^+_{a\downarrow} \hat{c}_{a\downarrow} \right) + \varepsilon_b \left( \hat{c}^+_{b\uparrow} \hat{c}_{b\uparrow} + \hat{c}^+_{b\downarrow} \hat{c}_{b\downarrow} \right) .$$

- a) Show that the state  $|\phi_1\rangle = \hat{c}^+_{a\uparrow}\hat{c}^+_{b\uparrow}|0\rangle$  is an eigenstate of the system. Which energy has the system in this state?
- b) Show that the state  $|\phi_2\rangle = \frac{1}{\sqrt{2}} \left(\hat{c}^+_{a\uparrow} + \hat{c}^+_{a\downarrow}\right) \hat{c}^+_{b\uparrow} |0\rangle$  is normalized. Is  $|\phi_2\rangle$  an eigenstate of the system?
- c) Calculate for  $|\phi_1\rangle$  and  $|\phi_2\rangle$ , respectively, the expectation values for the spin operators

$$\hat{S}_z = \frac{\hbar}{2} \sum_j \left( \hat{c}_{j\uparrow}^+ c_{j\uparrow} - \hat{c}_{j\downarrow}^+ \hat{c}_{j\downarrow} \right) \quad \text{and} \quad \hat{S}_x = \frac{\hbar}{2} \sum_j \left( \hat{c}_{j\uparrow}^+ \hat{c}_{j\downarrow} + \hat{c}_{j\downarrow}^+ c_{j\uparrow} \right) \quad \text{with} \quad j = a, b.$$

## (3 points)

(2 points)