

Numerics for partial differential equations

Introduction

- Overview of 2. order PDEs
(classification, classical model problems)
- Overview of most important discretization concepts
(finite differences, finite elements, finite volumes)

Finite Differences

- Exm. heat equation
(stability concepts, convergence)
- Exm. transport equation
(stability, convergence, dissipation)

Finite Elements

- Exm. Poisson equation
(matrix assembly, a priori error estimates)
- Adaptivity
(grid refinement, a posteriori error estimates)

Detailed table of contents

I. Introduction

1. Partial differential equations (PDEs) of 2. order
What is a PDE?
Classification of 2. order PDEs
Elliptic PDEs: energy minimization
Parabolic PDEs: gradient flows
Hyperbolic PDEs: conservation laws
Reduction of hyperbolic PDEs of 2. to 1. order in 2D

2. Discretization concepts

- Finite Differences
- Finite Volumes
- Finite Elements/Galerkin methods

II. Finite Differences

1. Parabolic PDEs

- a) 1D model system: heat conduction
theta-method
Example simulation
Digression: Fourier transform & Plancherel's/Parseval's theorem
L2 stability
Digression: Semidiscrete Fourier transform
L2 stability
L2 von Neumann stability
Example simulation
L-infinity stability: maximum principle
l-infinity stability: discrete maximum principle
l-infinity stability wrt. data term
Consistency: truncation error
Convergence (consistency + stability = convergence!)
- b) Extension to higher dimensions - 2D model system: heat conduction
Summary 2D-domain
Alternating Direction Implicit (ADI) method
- c) Extension to nonlinear parabolic PDEs
Model system: Allen-Cahn equation
l-infinity stability
Example simulation

2. Hyperbolic PDEs

- a) 1D model system: transport equation
Upwind method
Example simulation
Digression: Lebesgue spaces
Lp stability
L2 stability
l-infinity stability
Domain of dependence
Convergence (consistency + stability)
Higher order methods
L2 stability of higher order methods
Example simulation
Numerical dissipation
- b) Extension to 2D and nonlinear PDEs
Necessary changes...

III. Finite Elements (for Elliptic PDEs)

1. Elliptic PDE theory
Introduction to Euler-Lagrange equations
Digression: weak derivative
Digression: Sobolev spaces
Linear elliptic boundary value problems
Weak solution
Associated bilinear form
Existence, uniqueness, equivalence to minimization
Example: weak solution to Poisson problem
Digression: Poincaré inequality
Weak solution to linear elliptic boundary value problems
2. Ritz-Galerkin method
 - a) Existence, uniqueness, minimization
 - b) Principle of error estimation
 - c) Linear Finite Elements
Setting and ansatz space
Matrix assembly
Computation of local system matrix
Mass and stiffness matrix for regular grid
Interpolation error (on reference triangle)
Quasiuniform triangulation
A priori H1 error estimate
Aubin-Nitsche duality trick: a priori L2 error estimate
Example simulation
PDE Regularity (Friedrichs' theorem)
 - d) General Finite Element Method
Finite Element examples
Interpolant
Bramble-Hilbert lemma
 - e) Variational Crimes
Deviation from exact calculation
Strang's first lemma
Example application
Strang's second lemma
Example application
Duality (extension of Aubin-Nitsche trick)
Example application
 - f) Digression: duality
Digression: dual spaces
Digression: H-1
Digression: elliptic differential operators
 - g) Adaptivity
Idea of adaptive FE solutions
A posteriori error estimation
A posteriori error estimation in L2
Local error estimator
Reliability of local error estimator
Efficiency of local error estimator
Clément interpolation
Refinement strategy
Properties of bisection method
Refinement procedure

3. Extension to nonlinear vectorial elliptic PDEs

- Model system: 3D elasticity
- Introduction to elasticity
- Energy density
- Principal stretches
- Variational formulation
Digression: techniques for nonlinear elliptic PDEs
Digression: weak convergence
Digression: substitute for Lax-Milgram lemma
- Well-posedness
- Discretization
- Gamma-convergence
- Convergence

Introduction: PDEs of 2. order

Partial differential equation

Def: Let $\Omega \subset \mathbb{R}^d$ open, $d \geq 2$. A partial differential equation (PDE) is an equation that connects a function $u: \Omega \rightarrow \mathbb{R}$ and its partial derivatives,

$$F(D^k u(x), D^{k-1} u(x), \dots, Du(x), u(x), x) = 0 \quad \forall x \in \Omega,$$
$$F: \mathbb{R}^{d^k} \times \mathbb{R}^{d^{k-1}} \times \dots \times \mathbb{R} \times \Omega \rightarrow \mathbb{R}. k \text{ is the } \underline{\text{order}} \text{ of the PDE, } u \text{ the sought function.}$$

Exm: Laplace - equation $\partial^2 u / \partial x_1^2 + \partial^2 u / \partial x_2^2 = 0$

Def: A PDE problem is a PDE together with boundary conditions on $T \subset \partial\Omega$

- Dirichlet bc : $u = g$ on T (g given)
- Neumann bc : $n^T A \nabla u = g$ on T ($A \in \mathbb{R}^{d \times d}$ regular, $n = \text{unit normal to } T$)
- Robin bc : $u + n^T A \nabla u = g$ on T

The PDE problem is to be solved for $u: \Omega \rightarrow \mathbb{R}$

Exm: Laplace equation with $u(x) = x_1$ on $\partial\Omega$ has the solution $u(x) = x_1$.

Introduction: PDEs of 2. order

Classification of 2. order PDEs

Notation: $\cdot u_{x_n} = \partial u / \partial x_n$, $u_{x_n x_2} = \partial^2 u / \partial x_n \partial x_2$, ...

$\cdot Du = (u_{x_1}, \dots, u_{x_d})$, $\nabla u = Du^T$

$\cdot \Delta u = \operatorname{div} \nabla u = u_{x_1 x_1} + \dots + u_{x_d x_d}$ "Laplace operator"

$\cdot \alpha = (\alpha_1, \dots, \alpha_d) \in \mathbb{N}_0^d$ is a multiindex of order $|\alpha| = \alpha_1 + \dots + \alpha_d$.

$D^\alpha u = \partial^{|\alpha|} u / \partial x_1^{\alpha_1} \dots \partial x_d^{\alpha_d}$

$\cdot \Omega$ & x_1, \dots, x_d usually denote only spatial variables; if one variable represents time we will call it t and take $[0, T] \times \Omega$ as the domain of the PDE; additionally, $u_t = u_x$.

We will only consider semilinear PDEs: $\sum_{i,j=1}^d a_{ij}(x) u_{x_i x_j}(x) + c(\nabla u(x), u(x), x) = 0$

Def: The PDE is called elliptic, if $A = (a_{ij})_{i,j}$ is positive definite

\cdot parabolic, if A is pos. semidefinite with a 1D eigenspace for 0

\cdot hyperbolic, if A has one positive & $d-1$ negative eigenvalues

Introduction: PDEs of 2. order

Elliptic PDEs: energy minimization

Elliptic PDEs often derive from the physical principle of energy minimization:

A physical system in equilibrium attains the state of minimal energy.

Let u be the state, $E[u]$ the energy $\Rightarrow E(u+\varphi) \geq E(u) \quad \forall \varphi: \Omega \rightarrow \mathbb{R}$

Exm: The potential energy of an electric charge distribution $u: \Omega \rightarrow \mathbb{R}$

(or of an elastic membrane with transversal displacement u or of a temperature distribution u) in an electric field f is

$$E(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx + \int_{\Omega} f u dx. \quad \text{On } \partial\Omega, \quad u = g \text{ is fixed.}$$

$$E(u+\varphi) - E(u) = \int_{\Omega} \nabla u \cdot \nabla \varphi + \frac{1}{2} |\nabla \varphi|^2 dx + \int_{\Omega} f \varphi dx$$

Stokes thm. /
integration by parts

$$= \int_{\partial\Omega} \varphi \nabla u \cdot n dx + \int_{\Omega} \frac{1}{2} |\nabla \varphi|^2 - \varphi \operatorname{div} \nabla u dx + \int_{\Omega} f \varphi dx \stackrel{!}{\geq} 0 \quad \forall \varphi \text{ with } \varphi = 0 \text{ on } \partial\Omega$$

\Rightarrow Poisson equation $\operatorname{div} \nabla u = \Delta u = f$ with Dirichlet bc $u = g$ on $\partial\Omega$

Introduction: PDEs of 2. order

Parabolic PDEs: gradient flows

A physical system first has to equilibrate, to get to the energy minimum.

Parabolic PDEs derive from the physical principle of gradient flows:

At any time point, the state moves into the direction which allows the fastest energy decrease (at same effort/cost).

Let $\tau D(\frac{\varphi}{\tau})$ be the cost of achieving the change $u \rightarrow u + \varphi$ in time τ

$$\Rightarrow u = \lim_{\tau \rightarrow 0} \left(\operatorname{argmin}_{\varphi} \tau D(\frac{\varphi}{\tau}) + E(u + \varphi) \right) / \tau$$

Exm: A change of the charge distribution (displacement/temperature) costs $D(\varphi) = \frac{1}{2} \int_{\Omega} \varphi^2 dx$.

$$\tau D(\frac{\varphi + \tilde{\varphi}}{\tau}) + E(u + \varphi + \tilde{\varphi}) - [\tau D(\frac{\varphi}{\tau}) + E(u + \varphi)] = \int_{\Omega} \frac{\varphi \tilde{\varphi}}{\tau} + \frac{\tilde{\varphi}^2}{2\tau} + \frac{|\nabla \tilde{\varphi}|^2}{2} + \nabla(u + \varphi) \cdot \nabla \tilde{\varphi} + f \tilde{\varphi} dx$$

$$= \int_{\Omega} \tilde{\varphi} \left(\frac{\varphi}{\tau} - \Delta(u + \varphi) + f \right) dx + \int_{\partial\Omega} \tilde{\varphi} \nabla(u + \varphi) \cdot n dx + O(\tilde{\varphi}^2) \geq 0 \quad \forall \tilde{\varphi} \text{ with } \tilde{\varphi} = 0 \text{ on } \partial\Omega$$

$$\Rightarrow \frac{\varphi}{\tau} = \Delta(u + \varphi) - f \Rightarrow \varphi \rightarrow 0 \text{ for } \tau \rightarrow 0$$

heat equation $u - \Delta u = -f$ with Dirichlet bc $u = g$ on $\partial\Omega$, $u = u_0$ for $t = 0$

Introduction: PDEs of 2. order

Hyperbolic PDEs: conservation laws

Hyperbolic PDEs often derive from the physical principle of conservation:

The change of an extensive quantity (e.g. mass, momentum, energy) in a volume V is only possible via transport through ∂V .

Let $q(x,t)$ be the flux of the (mass, energy etc.) density $\rho \Rightarrow \frac{d}{dt} \int_V \rho(x,t) dx = - \int_{\partial V} q(x,t) \cdot n dx$

$$\Rightarrow 0 = \int_V \dot{\rho} + \operatorname{div} q dx \quad \forall V \subset \Omega \quad \Rightarrow \quad \dot{\rho} + \operatorname{div} q = 0$$

Gauss theorem

Exm: A material with mass distribution ρ moves at velocity $v(x,t)$

(e.g. water in clouds is transported by wind) $\Rightarrow q = v \rho$

\Rightarrow transport equation $\dot{\rho} + \operatorname{div}(v \rho) = 0$ with bc $\rho = \rho_0$ for $t = 0$.

Exm: The kinetic & potential energy density within a membrane under tension

with displacement u is $\rho(x,t) = \frac{m}{2} |\dot{u}|^2 + E \frac{|\nabla u|^2}{2}$ ($m =$ membrane mass density,

$E =$ elastic modulus), the energy flux is $q = -E \dot{u} \nabla u \Rightarrow \dot{\rho} = m \dot{u} \ddot{u} + E \nabla u \cdot \nabla \dot{u} = E \operatorname{div}(\dot{u} \nabla u)$

\Rightarrow wave equation $\ddot{u} = \frac{E}{m} \Delta u$ with bc $u = g$ on $\partial \Omega$, $u = u_0$ for $t = 0$

Introduction: PDEs of 2. order

Reduction of hyperbolic PDEs from 2. to 1. order in 2D

$$\ddot{u} - \partial^2 u / \partial x^2 = f(u, \dot{u}, \partial u / \partial x, x, t)$$

$$\Leftrightarrow \begin{cases} \dot{v} + \partial v / \partial x = f(u, \dot{u}, \partial u / \partial x, x, t) \\ \dot{u} - \partial u / \partial x = v \end{cases}$$

= system of coupled transport equations of 1. order

$$\text{" } \partial_t^2 - \partial_x^2 = (\partial_t - \partial_x)(\partial_t + \partial_x) \text{"}$$

If f only depends on x & t , one can first solve

$$\dot{v} + \partial v / \partial x = f(x, t)$$

and then

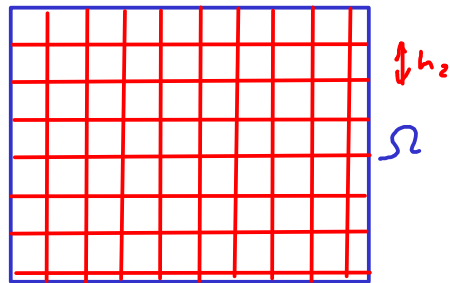
$$\dot{u} - \partial u / \partial x = v$$

=> we will only consider hyperbolic PDEs of 1. order!

Introduction: discretization concepts

Finite Differences

If one has a regular grid on Ω with grid width h_1, \dots, h_d in x_1, \dots, x_d -direction,



then the derivatives of a sufficiently $\rightarrow h_n$ often differentiable function u can be approximated by difference quotients.

Exm: $\frac{\partial u}{\partial x_1} \approx \frac{u(x_1+h_1, x_2) - u(x_1, x_2)}{h_1}$, $\frac{\partial^2 u}{\partial x_2^2} \approx \frac{u(x_1, x_2+h_2) - 2u(x_1, x_2) + u(x_1, x_2-h_2)}{h_2^2}$

We seek an approximation u_{i_1, \dots, i_d} to $u(\underbrace{(i_1 h_1, \dots, i_d h_d)}_{\text{grid points}})$

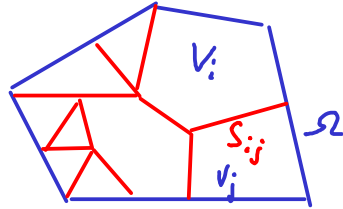
\Rightarrow In the PDE we replace all derivatives by difference quotients and obtain a system of equations in the u_{i_1, \dots, i_d} .

Exm: $i - \partial^2 u / \partial x^2 = 0$, $u(t=0, x) = e^{-x^2} \rightarrow \frac{u_{i+1, j} - u_{i, j}}{h_1} - \frac{u_{i, j+1} - 2u_{i, j} + u_{i, j-1}}{h_2^2} = 0$, $u_{0, j} = e^{-j^2 h_2^2} \forall i > 0, j$

Introduction: discretization concepts

Finite Volumes

On a mesh of volume elements V_i and their sides S_{ij}



one can apply the principle of conservation on each element: If f_{ik} denotes the mean value of f on V_i at time $t_k = k \Delta t$ and $\tilde{q}(f_{ik}, f_{jk}, x_{ij}, t)$ an approximation to the flux through S_{ij} , then

$$\frac{d}{dt} \int_{V_i} f(x, t) dx = - \int_{\partial V_i} q(x, t) \cdot n dx \quad \text{normal from } V_i \text{ to } V_j$$

$$\leadsto f_{i, k+1} - f_{ik} = \frac{-\Delta t}{|V_i|} \sum_j |S_{ij}| \tilde{q}(f_{ik}, f_{jk}, x_{ij}, t) \cdot n_{ij} \quad \forall k, i$$

midpoint of S_{ij}

which is a set of equations for the f_{ik} .

Exm: $\dot{p} + \text{div}(vp) = 0, p(t=0, x) = e^{-x^2}$

$$\leadsto f_{i, k+1} - f_{ik} = \frac{-\Delta t}{|V_i|} \sum_j |S_{ij}| v(x_{ij}, k \Delta t) \cdot n_{ij} \frac{f_{ik} + f_{jk}}{2}, \quad f_{i0} = \int_{V_i} e^{-x^2} dx \quad \forall k > 0, i$$

Introduction: discretization concepts

Galerkin methods/Finite Elements

Let V denote a finite-dimensional subspace of the space F of functions $u: \Omega \rightarrow \mathbb{R}$ with basis ϕ_1, \dots, ϕ_k . The energy minimization principle can also be applied to V instead of to F :

$$u = \operatorname{argmin}_{v \in F} E(v) = \text{solution of the associated elliptic PDE}$$

$$\rightarrow \hat{u} = \sum_{i=1}^k u_i \phi_i = \operatorname{argmin}_{v \in V} E(v) = \text{approximation of the PDE solution}$$

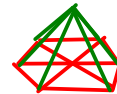
$$\Rightarrow 0 = \frac{\partial}{\partial u_i} E\left(\sum_{i=1}^k u_i \phi_i\right) \quad \forall i = 1, \dots, k$$

\hat{u} is called Galerkin approximation. If the ϕ_i are local (i.e. $\operatorname{supp} \phi_i \cap \operatorname{supp} \phi_j = \emptyset$ for most i, j), one obtains the special case of finite elements.

Exm: $\Delta u = 0$ on Ω , $u(x) = g(x)$ on $\partial\Omega$

triangle mesh  , V spanned by

$$E(u) = \int_{\Omega} \frac{|\nabla u|^2}{2} dx$$



$$\phi_i(x_j) = \begin{cases} 1 & , j=i \\ 0 & , j \neq i \end{cases}$$

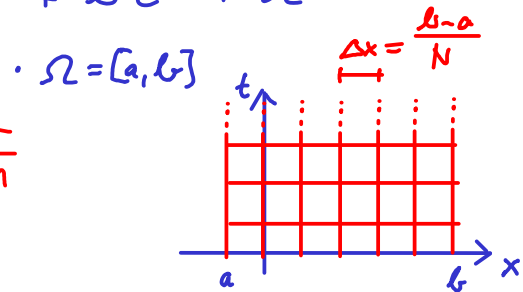
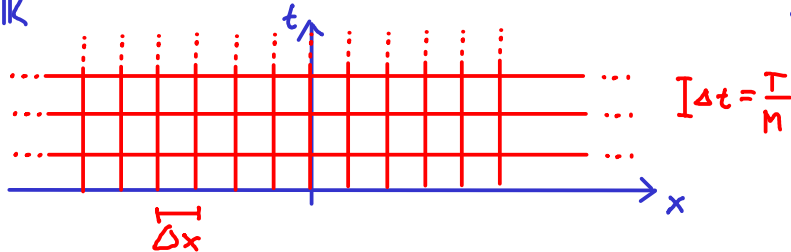
$$\rightarrow \sum_{j=1}^k \int_{\Omega} \nabla \phi_i \cdot \nabla \phi_j dx u_j = 0 \quad \forall i \text{ with } x_i \notin \partial\Omega, \quad u_i = g(x_i) \text{ else}$$

Finite Differences: parabolic PDEs (1D heat equation)

theta-method

$$u_t = u_{xx} \text{ on } [0, T] \times \Omega + bc + ic$$

mesh: $\cdot \Omega = \mathbb{R}$



explicit Euler-method:
$$\frac{U_j^{m+1} - U_j^m}{\Delta t} = \frac{U_{j-1}^m - 2U_j^m + U_{j+1}^m}{\Delta x^2}$$

$$\Leftrightarrow U_j^{m+1} = U_j^m + \frac{\Delta t}{\Delta x^2} (U_{j-1}^m - 2U_j^m + U_{j+1}^m)$$

implicit Euler-method:
$$\frac{U_j^{m+1} - U_j^m}{\Delta t} = \frac{U_{j-1}^{m+1} - 2U_j^{m+1} + U_{j+1}^{m+1}}{\Delta x^2}$$

$$\Leftrightarrow U_j^{m+1} - \frac{\Delta t}{\Delta x^2} (U_{j-1}^{m+1} - 2U_j^{m+1} + U_{j+1}^{m+1}) = U_j^m$$

θ-method:
$$\frac{U_j^{m+1} - U_j^m}{\Delta t} = (1-\theta) \frac{U_{j-1}^m - 2U_j^m + U_{j+1}^m}{\Delta x^2} + \theta \frac{U_{j-1}^{m+1} - 2U_j^{m+1} + U_{j+1}^{m+1}}{\Delta x^2}$$

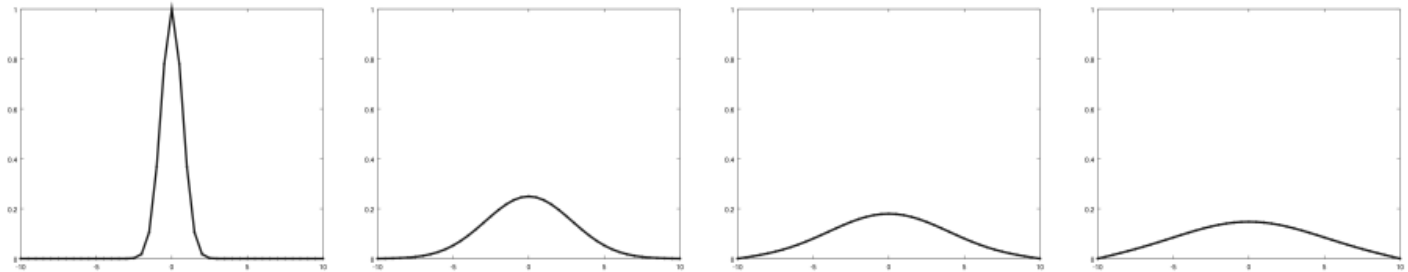
$\theta = 0$: expl. Euler ; $\theta = 1$: impl. Euler ; $\theta = \frac{1}{2}$: Crank-Nicolson

initial condition (ic): $U_j^0 = u(t=0, x_j)$; Dirichlet-bc: $U_0^m = u(t^m, a)$, $U_N^m = u(t^m, b)$

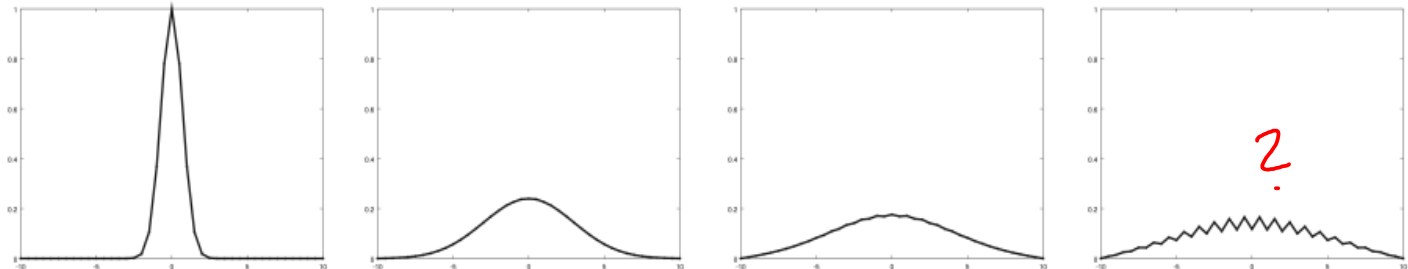
Finite Differences: parabolic PDEs (1D heat equation)

Example simulation $u_t = u_{xx}$ on $\Omega = [-10, 10]$, $u(t=0, x) = e^{-x^2}$, $u(t, \partial\Omega) = u(0, \partial\Omega) = 0$

expl. Euler, $\Delta x = \frac{1}{2}$, $\Delta t = 0,12$



expl. Euler, $\Delta x = \frac{1}{2}$, $\Delta t = 0,13$



Finite Differences: parabolic PDEs (1D heat equation)

Fourier-Transform & Plancherel's/Parseval's Theorem

Def: Fourier transform (FT) of $u: \mathbb{R} \rightarrow \mathbb{C}$: $\hat{u}(\xi) = F[u](\xi) = \int_{\mathbb{R}} u(x) e^{-ix\xi} dx$

Def: L^2 -norm of $u: \Omega \rightarrow \mathbb{C}$: $\|u\|_{L^2} = \left(\int_{\Omega} |u(x)|^2 dx \right)^{1/2}$

$L^2(\Omega) = \{u: \Omega \rightarrow \mathbb{C} \mid u \text{ measurable, } \|u\|_{L^2} < \infty\}$

corresponding L^2 inner product $(u, v)_{L^2} = \int_{\Omega} (u(x), v(x))_{\mathbb{C}} dx = \int_{\Omega} \operatorname{Re}(\overline{u(x)}v(x)) dx$

Thm: Inverse FT (IFT) of $\hat{u}: \mathbb{R} \rightarrow \mathbb{C}$: $u(x) = F^{-1}[\hat{u}](x) = \frac{1}{2\pi} \int_{\mathbb{R}} \hat{u}(\xi) e^{ix\xi} d\xi$

Pf: $\frac{1}{2\pi} \int_{\mathbb{R}} \int_{\mathbb{R}} u(\tilde{x}) e^{-i\tilde{x}\xi} d\tilde{x} e^{ix\xi} d\xi = \frac{1}{2\pi} \int_{\mathbb{R}} \int_{\mathbb{R}} e^{i(x-\tilde{x})\xi} d\xi u(\tilde{x}) d\tilde{x} = \int_{\mathbb{R}} u(\tilde{x}) \delta(\tilde{x}-x) d\tilde{x}$
 \Rightarrow make rigorous via smoothing & approximation argument □

Thm: (Plancherel's/Parseval's Theorem) $\|\hat{u}\|_{L^2} = \sqrt{2\pi} \|u\|_{L^2}$

$\Rightarrow \frac{F}{\sqrt{2\pi}}$ is isometry from $L^2(\mathbb{R})$ to $L^2(\mathbb{R})$

Pf: $\int_{\mathbb{R}} \hat{u}(\xi) \overline{v(\xi)} d\xi = \int_{\mathbb{R}} \int_{\mathbb{R}} u(x) e^{-ix\xi} dx \overline{v(\xi)} d\xi$

$$= \int_{\mathbb{R}} \int_{\mathbb{R}} v(\xi) e^{-ix\xi} d\xi u(x) dx = \int_{\mathbb{R}} u(x) \hat{v}(x) dx$$

• choose $v(\xi) = \overline{\hat{u}(\xi)} = 2\pi F^{-1}[\overline{\hat{u}}](x)$ □

Finite Differences: parabolic PDEs (1D heat equation)

L2 stability

Stability $\hat{=}$ small perturbation in the data (e.g. bc or rhs of PDE) results only in small change of the solution

measurement of perturbation strength possible with different norms
 \Rightarrow stability wrt. different norms

$$\left. \begin{array}{l} u_t - u_{xx} = f, \quad u(t=0) = u_0 \\ \tilde{u}_t - \tilde{u}_{xx} = \tilde{f}, \quad \tilde{u}(t=0) = \tilde{u}_0 \end{array} \right\} \Rightarrow v = \tilde{u} - u \text{ solves } v_t - v_{xx} = \tilde{f} - f, \quad v(t=0) = \tilde{u}_0 - u_0$$

for Neumann bc see homework

Thm: Let u solve $u_t - u_{xx} = f$ on $[0, \infty) \times \Omega$ ($\Omega = \mathbb{R}$ or $\Omega = [a, b]$ with 0-Dirichlet-bc).

$$\|u(t)\|_{L^2} \leq \|u(t=0)\|_{L^2} + \int_0^t \|f(s)\|_{L^2} ds$$

Cauchy-Schwarz

Pf: for $\Omega = [a, b]$, $2\|u\|_{L^2} \frac{d}{dt} \|u\|_{L^2} = 2 \int_a^b u u_t dx = 2 \int_a^b u (\Delta u + f) dx = 2 \int_a^b -|u_x|^2 + u f dx \leq 2 \|u\|_{L^2} \|f\|_{L^2}$

for $\Omega = \mathbb{R}$, $F[u_t - u_{xx} - f](\xi) = \hat{u}_t(\xi, t) - (-\xi^2) \hat{u}(\xi, t) - \hat{f}(\xi, t)$

$$\Rightarrow \forall \xi: \frac{d}{dt} \hat{u} = -\xi^2 \hat{u} + \hat{f} \Rightarrow 2\|\hat{u}\|_{L^2} \frac{d}{dt} \|\hat{u}\|_{L^2} = 2 \int_{\mathbb{R}} \xi^2 |\hat{u}|^2 - \hat{u} \hat{f} d\xi \leq 2\|\hat{u}\|_{L^2} \|\hat{f}\|_{L^2}$$

in both cases $\frac{d}{dt} \|u\|_{L^2} \leq \|f\|_{L^2}$ □

Finite Differences: parabolic PDEs (1D heat equation)

Semidiscrete Fourier-Transform

Def: The semidiscrete FT of a function U defined on the infinite grid $x_j = j\Delta x, j \in \mathbb{Z}$, is

$$\hat{U}: \left[-\frac{\pi}{\Delta x}, \frac{\pi}{\Delta x}\right] \rightarrow \mathbb{C}, \quad \hat{U}(k) = \Delta x \sum_{j=-\infty}^{\infty} U_j e^{-ikx_j}.$$

Def: l_2 -norm of $U: \mathbb{Z} \subset \Delta x \mathbb{Z} \rightarrow \mathbb{C}$: $\|U\|_{l_2} = \left(\Delta x \sum_{x_j \in \mathbb{Z}} |U_j|^2\right)^{\frac{1}{2}}$

Thm: Inverse semidiscrete FT is $U_j = \frac{1}{2\pi} \int_{-\pi/\Delta x}^{\pi/\Delta x} \hat{U}(k) e^{ikx_j} dk$

Thm: (Parseval) $\|\hat{U}\|_{l_2} = \sqrt{2\pi} \|U\|_{l_2}$

l2-Stability

θ -method for $u - \Delta u = f$:

(*)
↓

$$\frac{U_j^{m+1} - U_j^m}{\Delta t} = (1-\theta) \frac{U_{j-1}^m - 2U_j^m + U_{j+1}^m}{\Delta x^2} + \theta \frac{U_{j-1}^{m+1} - 2U_j^{m+1} + U_{j+1}^{m+1}}{\Delta x^2} + (1-\theta)F_j^m + \theta F_j^{m+1}, \quad U_j^0 = U_j^o$$

$$\frac{\tilde{U}_j^{m+1} - \tilde{U}_j^m}{\Delta t} = (1-\theta) \frac{\tilde{U}_{j-1}^m - 2\tilde{U}_j^m + \tilde{U}_{j+1}^m}{\Delta x^2} + \theta \frac{\tilde{U}_{j-1}^{m+1} - 2\tilde{U}_j^{m+1} + \tilde{U}_{j+1}^{m+1}}{\Delta x^2} + (1-\theta)\tilde{F}_j^m + \theta\tilde{F}_j^{m+1}, \quad \tilde{U}_j^0 = \tilde{U}_j^o$$

$\Rightarrow \tilde{u} - u$ solves same equations for data $\tilde{F} - F$ & ic $\tilde{U} - U$.

\Rightarrow for stability consider (as before) norm of the soln. as fcn. of the data Δic .

Thm: Let $U = (U_j^m)_{j \in Z}^{m \geq 0}$ solve (*) (on $Z = \mathbb{Z}$ or $Z = \{0, 1, \dots, N\}$ with 0-Dirichlet-bc).

If $(1-2\theta) \frac{\Delta t}{\Delta x^2} \leq \frac{1}{2}$, $\|U\|_{l^2} \leq \|U^0\|_{l^2} + \Delta t \sum_{k=1}^m \|\theta F^k + (1-\theta) F^{k-1}\|_{l^2}$.

'method is l^2 -stable'. $\leq \Delta t (\sum_{k=1}^{m-1} \|F^k\|_{l^2} + (1-\theta)\|F^0\|_{l^2} + \theta\|F^m\|_{l^2}) = \int_0^{\Delta t} \|F\|_{l^2} dt$ (quadrature for $\int_0^{\Delta t} \|F\|_{l^2} dt$)

Cor: Impl. Euler- & Crank-Nicolson-method are unconditionally l^2 -stable, expl. Euler-method is conditionally l^2 -stable (for $\Delta t \leq \frac{\Delta x^2}{2}$).

Pf : ∴ Abgeleitet

$$\mu = \Delta t / \Delta x^2$$

• $\hat{z} = \mathbb{Z}$: Plug in $U_j^m = \frac{1}{2\pi} \int_{-\pi/\Delta x}^{\pi/\Delta x} e^{ikj\Delta x} \hat{U}^m(k) dk$

$$\Rightarrow \frac{1}{2\pi} \int_{-\pi/\Delta x}^{\pi/\Delta x} e^{ikj\Delta x} \frac{\hat{U}^{m+1} - \hat{U}^m}{\Delta t} dk = \frac{1}{2\pi} \int_{-\pi/\Delta x}^{\pi/\Delta x} (1-\theta) \frac{e^{ik(j-1)\Delta x} - 2e^{ikj\Delta x} + e^{ik(j+1)\Delta x}}{\Delta x^2} \hat{U}^m$$

$$+ \theta \frac{e^{ik(j-1)\Delta x} - 2e^{ikj\Delta x} + e^{ik(j+1)\Delta x}}{\Delta x^2} \hat{U}^{m+1} + e^{ikj\Delta x} [(1-\theta) \hat{F}^m + \theta \hat{F}^{m+1}] dk$$

$$\Rightarrow 0 = \frac{1}{2\pi} \int_{-\pi/\Delta x}^{\pi/\Delta x} e^{ikj\Delta x} \left[\hat{U}^{m+1} (1 - \theta \mu (e^{-ik\Delta x} - 2 + e^{ik\Delta x})) - \hat{U}^m (1 + (1-\theta)\mu (e^{-ik\Delta x} - 2 + e^{ik\Delta x})) \right.$$

$$\left. + \Delta t [(1-\theta) \hat{F}^m + \theta \hat{F}^{m+1}] \right] dk$$

$$\Rightarrow \hat{U}^{m+1}(k) = \frac{1 + (1-\theta)\mu (e^{-ik\Delta x} - 2 + e^{ik\Delta x})}{1 - \theta \mu (e^{-ik\Delta x} - 2 + e^{ik\Delta x})} \hat{U}^m(k) - \frac{\Delta t [(1-\theta) \hat{F}^m(k) + \theta \hat{F}^{m+1}(k)]}{1 - \theta \mu (e^{-ik\Delta x} - 2 + e^{ik\Delta x})}, k \in \left[\frac{-\pi}{\Delta x}, \frac{\pi}{\Delta x} \right]$$

$-4 \sin^2(k\Delta x/2)$

$$\Rightarrow \|\hat{U}^{m+1}\|_{L^2} \leq \underbrace{\max_k \left| \frac{1 - (1-\theta)4\mu \sin^2(\frac{k\Delta x}{2})}{1 + \theta 4\mu \sin^2(\frac{k\Delta x}{2})} \right|}_{\leq 1} \|\hat{U}^m\|_{L^2} + \Delta t [\|\theta \hat{F}^{m+1} + (1-\theta) \hat{F}^m\|_{L^2}]$$

$$\leq 1 \Leftrightarrow (1-2\theta)\mu \leq \frac{1}{2}$$

$$\cdot z = \{0, 1, \dots, N\}; \quad U^{m+1} = (I + \theta \mu L)^{-1} \left[(I - (1-\theta) \mu L) U^m + \theta \Delta t F^{m+1} + (1-\theta) \Delta t F^m \right]$$

$$U^m = \begin{pmatrix} U_1^m \\ \vdots \\ U_N^m \end{pmatrix}, \quad L = \begin{pmatrix} 2 & -1 & & \\ -1 & 2 & \dots & -1 \\ & & \dots & \\ -1 & 2 & \dots & -1 \\ & & & 2 & -1 \end{pmatrix} \in \mathbb{R}^{(N-1) \times (N-1)}$$

By Gerschgorin's Theorem, L is pos. semidefinite with eigenvalues $\in [0, 4]$.

$$\Rightarrow \|U^{m+1}\|_{\ell^2} \leq \underbrace{\|(I + \theta \mu L)^{-1} (I - (1-\theta) \mu L)\|}_{\leq 1} \|U^m\|_{\ell^2} + \underbrace{\|(I + \theta \mu L)^{-1}\|}_{\leq 1} \Delta t \|\theta F^{m+1} + (1-\theta) F^m\|_{\ell^2}$$

$$\leq \max_{\text{eigenvalues } \lambda} \left| \frac{1 - (1-\theta)\mu\lambda}{1 + \theta\mu\lambda} \right| \leq 1 \quad \text{if } (1-2\theta)\mu \leq \frac{1}{2}$$

By induction we deduce in both cases

$$\|U^m\|_{\ell^2} \leq \|U^0\|_{\ell^2} + \Delta t \sum_{k=1}^m \|\theta F^k + (1-\theta) F^{k-1}\|_{\ell^2} \quad \square$$

ℓ^2 -stability for Neumann boundary conditions: see home work

l2-von Neumann-Stability

Conditions for stability are typically strongest for $z = z$, since then all frequencies k occur, while on a finite grid there are only finitely many frequencies/eigenmodes. Formally, we can check l^2 -stability by inserting $u_j^m = [\lambda(k)]^m e^{ijk\Delta x}$ into the discrete method with $F_j^m = 0$ and by then requiring $|\lambda(k)| \leq 1$.

Def: The method is called l^2 -von Neumann-stable, if there exists $C > 0$ independent of Δt such that $\|u^m\|_{l^2} \leq C \left[\|u^0\|_{l^2} + \Delta t \sum_{k=0}^m \|F^k\|_{l^2} \right]$, $m = 1, \dots, M = T/\Delta t$.

Thm: The θ -method is l^2 -von Neumann-stable if $(1-2\theta) \frac{\Delta t}{\Delta x^2} \leq \frac{1}{2} + c\Delta t$.

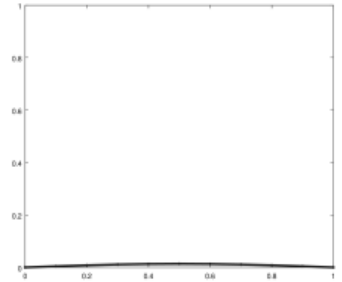
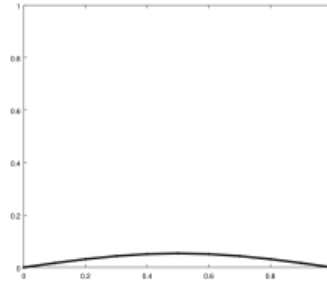
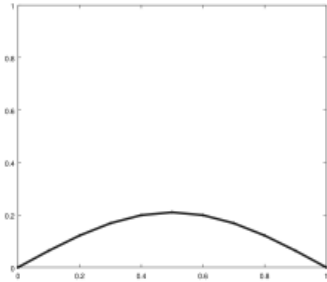
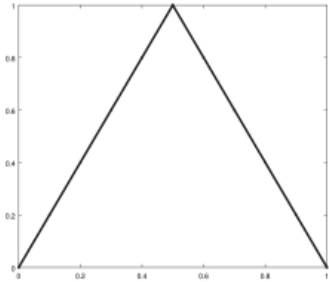
Pf : homework (use $(1+c\Delta t)^m \leq (1 + \frac{cT}{M})^M \leq e^{cT}$) □

Finite Differences: parabolic PDEs (1D heat equation)

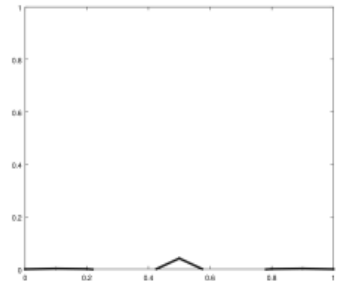
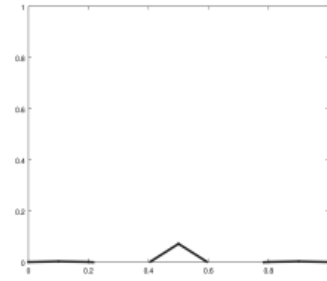
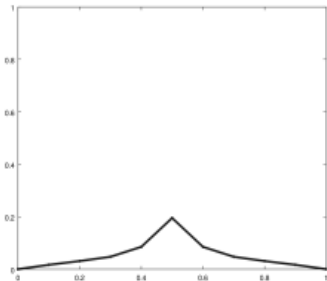
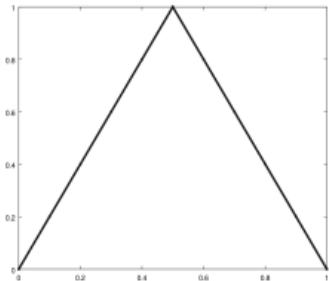
Example simulation

$$u_t = u_{xx} \text{ on } [0, T] \times [0, 1], \quad u(0, x) = 1 - 2|x - \frac{1}{2}|, \quad u(t, 0, 1) = 0$$

$$\text{impl. Euler} \quad \Delta t = \Delta x = \frac{1}{10}$$



$$\text{Crank-Nicolson} \quad \Delta t = \Delta x = \frac{1}{10}$$



Finite Differences: parabolic PDEs (1D heat equation)

L^∞ -stability: maximum principle

$$\Omega = (a, b); \quad \left. \begin{array}{l} u - u_{xx} = f, \quad u|_{t=0} = u_0, \quad u(\cdot, a) = u_a, \quad u(\cdot, b) = u_b \\ \tilde{u} - \tilde{u}_{xx} = \tilde{f}, \quad \tilde{u}|_{t=0} = \tilde{u}_0, \quad \tilde{u}(\cdot, a) = \tilde{u}_a, \quad \tilde{u}(\cdot, b) = \tilde{u}_b \end{array} \right\} \Rightarrow \begin{cases} v = \tilde{u} - u \text{ solves } v - v_{xx} = 0, \\ v|_{t=0} = \tilde{u}_0 - u_0, \quad v(\cdot, a) = \tilde{u}_a - u_a, \quad v(\cdot, b) = \tilde{u}_b - u_b \end{cases}$$

Thm: (Maximum principle) Let $u - \Delta u = f \leq 0$, then $\min_{(t,x) \in [0,T] \times \bar{\Omega}} u(t,x) = \max_{t=0 \vee x \in \partial\Omega} u(t,x)$.

Pf: \because Let $u(t,x)$, $t \in (0,T]$, $x \in \Omega$ be a loc. maximum $\Rightarrow u_x = 0$, $u_t \geq 0$, $u_{xx} \leq 0$

• if $f(t,x) < 0$, we have $0 > f = u_t - u_{xx} \geq 0 \quad \downarrow$

• if $f(t,x) \leq 0$, set $v(t,x) = u(t,x) + \frac{\varepsilon}{2} x^2 \Rightarrow v - v_{xx} = f - \varepsilon < 0$

\Rightarrow by previous argument, v attains maximum on $\underbrace{\{0\} \times \Omega \cup (0,T] \times \partial\Omega}_{=: \mathcal{R}}$

$$\Rightarrow u(t,x) \leq v(t,x) \leq \max_{(t,x) \in \mathcal{R}} v(t,x) \leq \max_{(t,x) \in \mathcal{R}} u(t,x) + \varepsilon \max(|a|, |b|)^2 \xrightarrow{\varepsilon \rightarrow 0} \max_{(t,x) \in \mathcal{R}} u(t,x)$$

• analogous for minima □

Cor: (L^∞ -stability) $\|\tilde{u} - u\|_{L^\infty} = \sup_{x,t} |\tilde{u}(t,x) - u(t,x)| \leq \max(\|\tilde{u}_0 - u_0\|_{L^\infty}, \|\tilde{u}_a - u_a\|_{L^\infty}, \|\tilde{u}_b - u_b\|_{L^\infty})$

$\sup_{x \in \Omega} |\tilde{u}_0(x) - u_0(x)|$ $\sup_{t \in (0,T)} |\tilde{u}_a(t) - u_a(t)|$

Finite Differences: parabolic PDEs (1D heat equation)

Γ^∞ -stability: discrete maximum principle

θ -method for $u - \Delta u = f$: ($\mu = \frac{\Delta t}{\Delta x^2}$)

(*)
↓

$$\frac{U_j^{m+1} - U_j^m}{\Delta t} = (1-\theta) \frac{U_{j-1}^m - 2U_j^m + U_{j+1}^m}{\Delta x^2} + \theta \frac{U_{j-1}^{m+1} - 2U_j^{m+1} + U_{j+1}^{m+1}}{\Delta x^2} + (1-\theta)F_j^m + \theta F_j^{m+1}$$

$$\left\{ \begin{array}{l} U_0^m = A_j^m \\ U_N^m = B_j^m \\ U_j^0 = O_j \end{array} \right.$$

If \tilde{U}_j^m is soln. for data F & i.c. \tilde{O} & b.c. \tilde{A}, \tilde{B} ,

then $\tilde{U} - U$ solves the eq. for data 0, i.c. $\tilde{O} - O$ & b.c. $\tilde{A} - A, \tilde{B} - B$

\Rightarrow for stability wrt. i.c. & b.c. examine, how $\|\tilde{U} - U\|$ depends on $\|\tilde{A} - A\|, \|\tilde{B} - B\|, \|\tilde{O} - O\|$

Thm: (Discrete maximum principle) Let U solve (*) with $F=0$ for $\theta \in [0, 1]$ & $\mu(1-\theta) \leq \frac{1}{2}$. Then

$$\min \{ U_j^m \mid m=0 \vee j=0 \vee j=N \} \leq U_j^m \leq \max \{ U_j^m \mid m=0 \vee j=0 \vee j=N \} \quad \forall j, m$$

$$\begin{aligned} \text{Pf: } (1+2\theta\mu)U_j^{m+1} &= \overbrace{\theta\mu(U_{j+1}^{m+1} + U_{j-1}^{m+1})}^{\geq 0} + (1-\theta)\mu(U_{j+1}^m + U_{j-1}^m) + \overbrace{(1-2(1-\theta)\mu)U_j^m}^{\geq 0} \\ &\leq 2\theta\mu U^* + 2(1-\theta)\mu U^* + (1-2(1-\theta)\mu)U^* \leq (1+2\theta\mu)U^* \\ &\quad \text{where } U^* = \max(U_{j+1}^{m+1}, U_{j-1}^{m+1}, U_{j+1}^m, U_{j-1}^m, U_j^m) \end{aligned}$$

\Rightarrow there cannot be a strict maximum for $m \neq 0$ & $j \notin \{0, N\}$ (analogous for minimal) \square

Cor: (Γ^∞ -stability) $\theta \in [0, 1], \mu(1-\theta) \leq \frac{1}{2}$. $\|u - \tilde{u}\|_{\ell^\infty} = \max_{j, m} |U_j^m - \tilde{U}_j^m| \leq \max(\|\tilde{A} - A\|_{\ell^\infty}, \|\tilde{B} - B\|_{\ell^\infty}, \|\tilde{O} - O\|_{\ell^\infty})$.

Finite Differences: parabolic PDEs (1D heat equation)

l^∞ -stability wrt. data term

Note : l^∞ -stability has stricter conditions on μ than l^2 -stability.

Crank-Nicolson-method : l^2 -stable $\forall \mu$, l^∞ -stable for $\mu \leq 1$.

Thm: Let U solve (X) with $A = B = 0 = 0$ for $\theta \in [0, 1]$, $\mu(1-\theta) \leq \frac{1}{2}$, then $\|U\|_{l^\infty} \leq T \|F\|_{l^\infty}$

$$\|U\|_{l^\infty} \leq \Delta t \sum_{k=1}^m \|(1-\theta)F^{k-1} + \theta F^k\|_{l^\infty} \leq T \|F\|_{l^\infty}.$$

Pf : $(1+2\theta\mu)U_j^{m+1} = \theta\mu(U_{j+1}^{m+1} + U_{j-1}^{m+1}) + (1-\theta)\mu(U_{j+1}^m + U_{j-1}^m) + (1-2(1-\theta)\mu)U_j^m + \Delta t[(1-\theta)F_j^m + \theta F_j^{m+1}]$

$$\leq 2\theta\mu \|U^{m+1}\|_{l^\infty} + 2(1-\theta)\mu \|U^m\|_{l^\infty} + (1-2(1-\theta)\mu) \|U^m\|_{l^\infty} + \Delta t \|(1-\theta)F_j^m + \theta F_j^{m+1}\|_{l^\infty}$$

$$\|G\|_{l^\infty} = \max\{|G_1|, \dots, |G_{N-1}|\}$$

Take maximum over $j \in \{1, \dots, N-1\} \Rightarrow \|U^{m+1}\|_{l^\infty} \leq \|U^m\|_{l^\infty} + \Delta t \|(1-\theta)F^m + \theta F^{m+1}\|_{l^\infty}$

$$\Rightarrow \|U^m\|_{l^\infty} \leq \underbrace{\|U^0\|_{l^\infty}}_{=0} + \Delta t \sum_{k=1}^m \|(1-\theta)F^{k-1} + \theta F^k\|_{l^\infty}, \quad m = 1, \dots, M \quad \square$$

Finite Differences: parabolic PDEs (1D heat equation)

Consistency: truncation error

Def: Let $u_j^m = u(t_m, x_j)$ for a smooth soln. u . The truncation error of the θ -method is

$$T_j^m = \frac{u_j^{m+1} - u_j^m}{\Delta t} - (1-\theta) \frac{u_{j-1}^m - 2u_j^m + u_{j+1}^m}{\Delta x^2} - \theta \frac{u_{j-1}^{m+1} - 2u_j^{m+1} + u_{j+1}^{m+1}}{\Delta x^2} - (1-\theta)f_j^m - \theta f_j^{m+1}$$

How big is the truncation error? Suppose u solves $u_t - u_{xx} = f$ and is sufficiently often differentiable. Then we perform a Taylor expansion:

$$u_j^{m+1} = u_j^m + u_t(t_m, x_j) \Delta t + O(\Delta t^2) \quad (\text{analogous for } f)$$

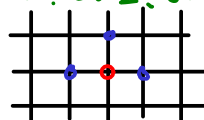
$$u_{j\pm 1}^m = u_j^m \pm u_x(t_m, x_j) \Delta x + \frac{1}{2} u_{xx}(t_m, x_j) \Delta x^2 \pm \frac{1}{6} u_{xxx}(t_m, x_j) \Delta x^3 + O(\Delta x^4)$$

$$\Rightarrow \text{Expl. Euler-method: } T_j^m = u_t + O(\Delta t) - u_{xx} + O(\Delta x^2) - f_j^m = O(\Delta t + \Delta x^2)$$

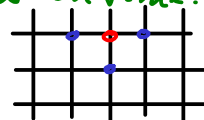
Thm: (Order of consistency) The θ -method satisfies $T_j^m = \begin{cases} O(\Delta t + \Delta x^2), & \theta \neq \frac{1}{2} \\ O(\Delta t^2 + \Delta x^2), & \theta = \frac{1}{2} \end{cases}$,

i.e. it is consistent of 2. order in space and 1. or 2. order in time.

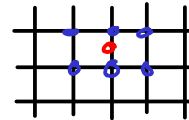
Pf; homework; note: Taylor exp. easiest around



expl. Euler



imp. Euler



CN \square

\Rightarrow Only C-N-method is 2. order consistent in time!

Finite Differences: parabolic PDEs (1D heat equation)

Convergence (consistency + stability = convergence!)

$$\frac{u_j^{m+1} - u_j^m}{\Delta t} = (1-\theta) \frac{u_{j-1}^m - 2u_j^m + u_{j+1}^m}{\Delta x^2} + \theta \frac{u_{j-1}^{m+1} - 2u_j^{m+1} + u_{j+1}^{m+1}}{\Delta x^2} + (1-\theta)f_j^m + \theta f_j^{m+1}$$

$$\frac{u_j^{m+1} - u_j^m}{\Delta t} = (1-\theta) \frac{u_{j-1}^m - 2u_j^m + u_{j+1}^m}{\Delta x^2} + \theta \frac{u_{j-1}^{m+1} - 2u_j^{m+1} + u_{j+1}^{m+1}}{\Delta x^2} + (1-\theta)f_j^m + \theta f_j^{m+1} + \tau_j^m$$

Set $e_j^m = u_j^m - u_j^m \Rightarrow \frac{e_j^{m+1} - e_j^m}{\Delta t} = (1-\theta) \frac{e_{j-1}^m - 2e_j^m + e_{j+1}^m}{\Delta x^2} + \theta \frac{e_{j-1}^{m+1} - 2e_j^{m+1} + e_{j+1}^{m+1}}{\Delta x^2} - \tau_j^m$

$\Rightarrow e$ solves θ -method with 0-i.c. & 0-b.c & data F_j^m so that $(1-\theta)F_j^m + \theta F_j^{m+1} = -\tau_j^m$

Thm: Let u (smooth) solve $u_t - u_{xx} = f$ with i.c. & b.c. and U_j^m solve the associated θ -method with $\theta \in [0, 1]$

$$\cdot \|u^m - u^m\|_{l^2} = \begin{cases} O(\Delta t + \Delta x^2) & , \theta \neq \frac{1}{2} \\ O(\Delta t^2 + \Delta x^2) & , \theta = \frac{1}{2} \end{cases} \quad \text{if } (1-2\theta) \frac{\Delta t}{\Delta x^2} \leq \frac{1}{2} \quad , m=1, \dots, M$$

$$\cdot \|u^m - u^m\|_{l^\infty} = \begin{cases} O(\Delta t + \Delta x^2) & , \theta \neq \frac{1}{2} \\ O(\Delta t^2 + \Delta x^2) & , \theta = \frac{1}{2} \end{cases} \quad \text{if } (1-\theta) \frac{\Delta t}{\Delta x^2} \leq \frac{1}{2} \quad , m=1, \dots, M$$

Pf: \Rightarrow stability $\|e^m\|_{l^2} \leq \Delta t \sum_{k=1}^m \|\theta F^k + (1-\theta)F^{k-1}\|_{l^2} \leq T \max_{k=1, \dots, M} \|\theta F^k + (1-\theta)F^{k-1}\|_{l^2}$

$$\|e^m\|_{l^\infty} \leq T \max_{k=1, \dots, M} \|\theta F^k + (1-\theta)F^{k-1}\|_{l^\infty}$$

consistency $\Rightarrow \|e^m\|_{l^2}, \|e^m\|_{l^\infty} \leq O(\Delta t^{1 \text{ or } 2} + \Delta x^2)$

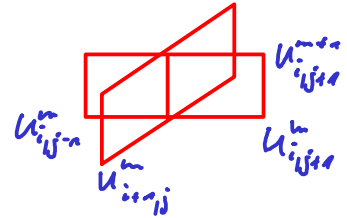
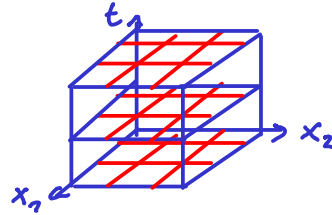
□

Finite Differences: parabolic PDEs (2D heat equation)

Summary 2D domain

$$u - \Delta u = f \text{ on } \Omega = [a_1, b_1] \times [a_2, b_2]$$

with i.c. & b.c.



Thm: 1) The θ -method for $\Omega = [a_1, b_1] \times [a_2, b_2]$ with grid widths $\Delta t, \Delta x_1, \Delta x_2$

and $\mu_1 = \frac{\Delta t}{\Delta x_1^2}, \mu_2 = \frac{\Delta t}{\Delta x_2^2}, \theta \in [0, 1]$, is

- \cdot l^2 -stable for $(1-2\theta)(\mu_1 + \mu_2) \leq \frac{1}{2}$,

- \cdot l^∞ -stable for $(1-\theta)(\mu_1 + \mu_2) \leq \frac{1}{2}$,

2) it is consistent with truncation error

- \cdot $\tau_{ij}^m = O(\Delta t + \Delta x_1^2 + \Delta x_2^2)$ for $\theta \neq \frac{1}{2}$

- \cdot $\tau_{ij}^m = O(\Delta t^2 + \Delta x_1^2 + \Delta x_2^2)$ for $\theta = \frac{1}{2}$

3) it is convergent of same order in the l^2/l^∞ -norm, if it is stable in that norm.

Pf: homework \square

Finite Differences: parabolic PDEs (2D heat equation)

Alternating Direction Implicit (ADI) method

$$u_t - \Delta u = 0$$

Write $\delta_x^2 U_{ij}^m = U_{i-1,j}^m - 2U_{ij}^m + U_{i+1,j}^m$; analogously δ_y^2

Crank-Nicolson: $(1 - \frac{\tau}{2} \mu_x \delta_x^2 - \frac{\tau}{2} \mu_y \delta_y^2) U_{ij}^{m+1} = (1 + \frac{\tau}{2} \mu_x \delta_x^2 + \frac{\tau}{2} \mu_y \delta_y^2) U_{ij}^m$

Def: ADI-method: $(1 - \frac{\tau}{2} \mu_x \delta_x^2) (1 - \frac{\tau}{2} \mu_y \delta_y^2) U_{ij}^{m+1} = (1 + \frac{\tau}{2} \mu_x \delta_x^2) (1 + \frac{\tau}{2} \mu_y \delta_y^2) U_{ij}^m$

$$\Leftrightarrow \begin{cases} (1 - \frac{\tau}{2} \mu_x \delta_x^2) U_{ij}^{m+1/2} = (1 + \frac{\tau}{2} \mu_y \delta_y^2) U_{ij}^m & | (1 + \frac{\tau}{2} \mu_x \delta_x^2) \dots \\ (1 - \frac{\tau}{2} \mu_y \delta_y^2) U_{ij}^{m+1} = (1 + \frac{\tau}{2} \mu_x \delta_x^2) U_{ij}^{m+1/2} & | (1 - \frac{\tau}{2} \mu_y \delta_y^2) \dots \end{cases} (*)$$

ADI differs from CN only in the term $\frac{\tau}{4} \mu_x \mu_y \delta_x^2 \delta_y^2$. It has the advantage that the resulting linear systems of equations (with matrices $1 - \frac{\tau}{2} \mu_x \delta_x^2$ and $1 - \frac{\tau}{2} \mu_y \delta_y^2$) are tridiagonal and can therefore be readily inverted by the Thomas-algorithm (the matrices of CN are much more complicated)!

Thm: The ADI method is L^2 -stable and for $\mu_x, \mu_y \leq 1$ also L^∞ -stable. It is consistent & convergent of 2. order in space & time.

Pf: homework; use (*) to derive stability.

□

Finite Differences: parabolic PDEs (2D heat equation)

Perspective: nonlinear PDEs

For nonlinear PDEs the stability of the method is often difficult to show.

Ex: $u_t = \Delta u + 2u(1-u^2) + f$ can be discretized as

$$U_{ij}^{m+1} - U_{ij}^m = \mu_1 \delta_n^2 U_{ij}^{m+1} + \mu_2 \delta_2^2 U_{ij}^{m+1} - 2\Delta t (U_{ij}^{m+1})^3 + 2\Delta t U_{ij}^m + \Delta t F_{ij}^m$$

U does not satisfy the maximum principle (u does not either), but one has:

Thm: $\|U^m\|_{\ell^\infty} \leq \max(1, \|U^0\|_{\ell^\infty}, \max_{x_{ij} \in \partial\Omega, n \leq m} |U_{ij}^n|) + T \|F\|_{\ell^\infty}$

Pf: Let $|U_{ij}^{m+1}| > 1$ for some $x_{ij} \in \text{int}\Omega$,

$$\begin{aligned} |U_{ij}^{m+1}|(1 + 2\mu_1 + 2\mu_2 + 2\Delta t) &\leq |U_{ij}^{m+1}|(1 + 2\mu_1 + 2\mu_2) + 2\Delta t (U_{ij}^{m+1})^3 \\ &= |U_{ij}^m + \mu_1(U_{i-1,j}^m + U_{i+1,j}^m) + \mu_2(U_{i,j-1}^m + U_{i,j+1}^m) + 2\Delta t U_{ij}^m + \Delta t F_{ij}^m| \\ &\leq (1 + 2\Delta t) \|U^m\|_{\ell^\infty} + \Delta t \|F\|_{\ell^\infty} + 2(\mu_1 + \mu_2) \|U^{m+1}\|_{\ell^\infty} \end{aligned}$$

$$\Rightarrow \|U^{m+1}\|_{\ell^\infty} \leq \max(1, \max_{x_{ij} \in \partial\Omega} |U_{ij}^{m+1}|) \quad \text{or} \quad \|U^{m+1}\|_{\ell^\infty} \leq \|U^m\|_{\ell^\infty} + \Delta t \|F\|_{\ell^\infty}$$

$$\Rightarrow \|U^{m+1}\|_{\ell^\infty} \leq \max(1, \max_{x_{ij} \in \partial\Omega} |U_{ij}^{m+1}|, \|U^m\|_{\ell^\infty}) + \Delta t \|F\|_{\ell^\infty}$$

\Rightarrow result follows by induction □

Finite Differences: parabolic PDEs (2D heat equation)

Perspective: nonlinear PDEs

Such boundedness results allow to reduce everything to the linear case.

Thm: (stability) Let \tilde{u} solve the discrete system for \tilde{F} , but same i.c./b.c. Then there exist

$C > 0$ s.t. for Δt small enough (both depends on $\|F\|_{\ell^\infty}, \|\tilde{F}\|_{\ell^\infty}, ic/bc$): $\|\tilde{u} - u\|_{\ell^\infty} \leq C \|\tilde{F} - F\|_{\ell^\infty}$

PF: $V_{ij}^{m+1} - V_{ij}^m = \mu_1 \delta_1^2 V_{ij}^{m+1} + \mu_2 \delta_2^2 V_{ij}^{m+1} - 2\Delta t \left[\underbrace{\left((u_{ij}^{m+1})^3 - (u_{ij}^m)^3 \right)}_{=: V_{ij}^{m+1} B_{ij}^{m+1}} \right] + 2\Delta t V_{ij}^m + \Delta t (\tilde{F}_{ij}^m - F_{ij}^m)$

$V = \tilde{u} - u$

Now standard argument for linear case:

$$V_{ij}^{m+1} (1 + 2\mu_1 + 2\mu_2 + 2\Delta t B_{ij}^{m+1}) = V_{ij}^m + \mu_1 (V_{i-1,j}^m + V_{i+1,j}^m) + \mu_2 (V_{i,j-1}^m + V_{i,j+1}^m) + 2\Delta t V_{ij}^m + \Delta t (\tilde{F}_{ij}^m - F_{ij}^m)$$

$$\Rightarrow \|V^{m+1}\|_{\ell^\infty} (1 - 2\Delta t \|B\|_{\ell^\infty}) \leq \|V^m\|_{\ell^\infty} (1 + 2\Delta t) + \Delta t \|\tilde{F} - F\|_{\ell^\infty}$$

$\|B\|_{\ell^\infty} \leq (\|\tilde{u}\|_{\ell^\infty} + \|u\|_{\ell^\infty})^2$ depends only on F, \tilde{F}, ic, bc . If $\Delta t \leq \min\{\frac{1}{2}, \frac{1}{8\|B\|_{\ell^\infty}}\}$, then

$$\frac{1+2\Delta t}{1-2\|B\|_{\ell^\infty}\Delta t} \leq \underbrace{1+2\Delta t+4\|B\|_{\ell^\infty}\Delta t}_{=: q}, \quad \frac{1}{1-2\|B\|_{\ell^\infty}\Delta t} \leq \frac{4}{3} \quad \text{and thus}$$

$$\|V^{m+1}\|_{\ell^\infty} \leq \|V^m\|_{\ell^\infty} (1 + 2\Delta t + 4\|B\|_{\ell^\infty}\Delta t) + 2\Delta t \|\tilde{F} - F\|_{\ell^\infty}$$

$$\leq \|V^0\|_{\ell^\infty} q^{m+1} + 2\Delta t \|\tilde{F} - F\|_{\ell^\infty} \sum_{i=0}^{m-1} q^i = \frac{q^m - 1}{q - 1} \leq \frac{q^m}{\Delta t} \leq \frac{(1 + (2+4\|B\|_{\ell^\infty})\frac{\Delta t}{M})^M}{\Delta t} \leq \frac{e^{(2+4\|B\|_{\ell^\infty})T}}{\Delta t} \quad \square$$

Convergence now follows from consistency and stability as before.

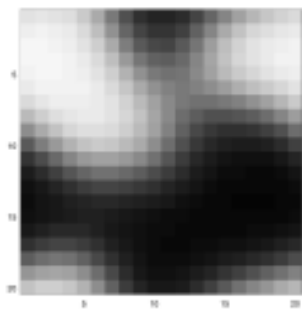
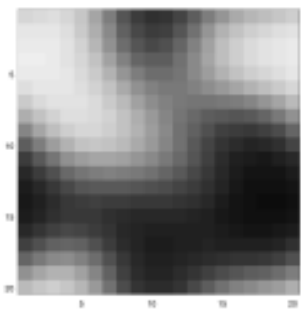
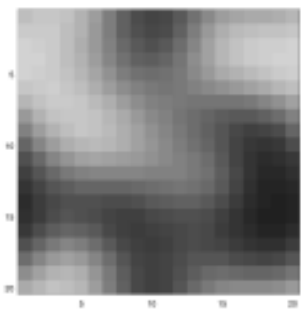
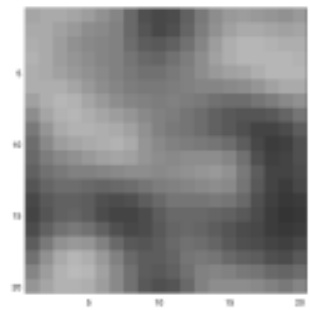
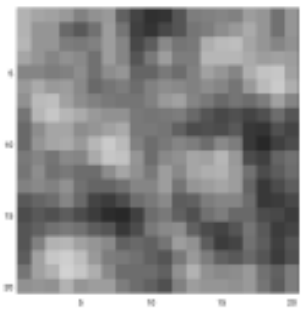
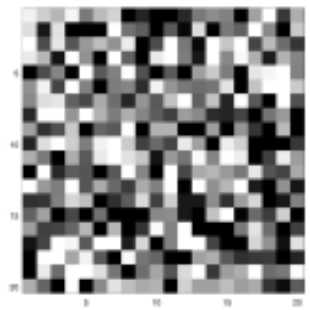
Finite Differences: parabolic PDEs (2D heat equation)

Example nonlinear PDE

$u_t = \Delta u + 2u(1-u^2)$ is gradient flow of

$$E(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 + (u^2 - 1)^2 dx$$

$\Rightarrow u = 1$ (white) or -1 (black) are assumed in the long-run.



Finite differences: hyperbolic PDEs (1D transport equation)

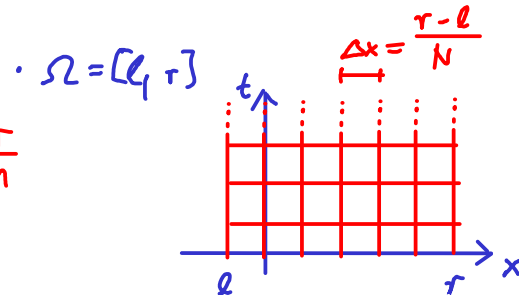
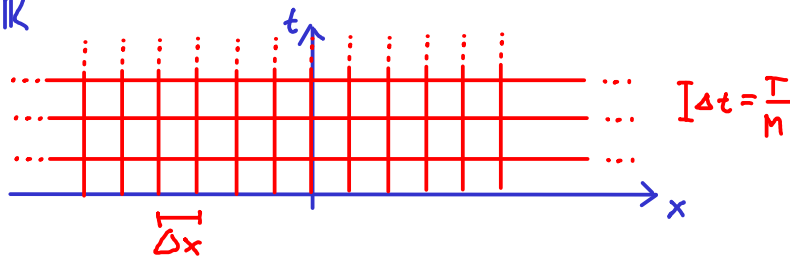
Upwind-method

$$u_t + a(t,x)u_x + c(t,x)u = f(t,x) \text{ on } [0,T] \times \Omega, \quad a \geq 0$$

$$\text{ic : } u(t=0, x) = u_0(x), \quad \text{bc : } u(t, l) = u_l(t)$$

what is interpretation?

grid : $\cdot \Omega = \mathbb{R}$



upwind-method

$$: \frac{U_j^{m+1} - U_j^m}{\Delta t} + A_j^m \frac{U_j^m - U_{j-1}^m}{\Delta x} + C_j^m U_j^m = F_j^m$$

$$\Leftrightarrow U_j^{m+1} = U_j^m - \frac{\Delta t}{\Delta x} A_j^m (U_j^m - U_{j-1}^m) - \Delta t C_j^m U_j^m + \Delta t F_j^m$$

$$\text{ic : } U_j^0 = u(t=0, x_j) \quad ; \quad \text{bc : } U_0^m = u(t^m, l)$$

$$\text{If } a \leq 0 \text{ choose : } \frac{U_j^{m+1} - U_j^m}{\Delta t} + A_j^m \frac{U_{j+1}^m - U_j^m}{\Delta x} + C_j^m U_j^m = F_j^m$$

$$\text{ic : } U_j^0 = u(t=0, x_j) \quad ; \quad \text{bc : } U_N^m = u(t^m, r)$$

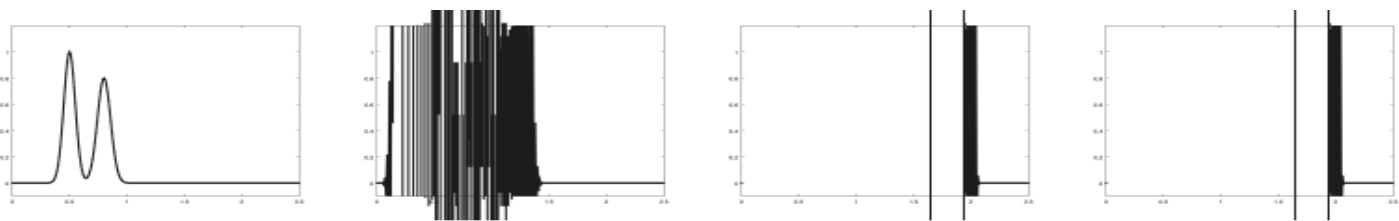
Finite differences: hyperbolic PDEs (1D transport equation)

Example simulation

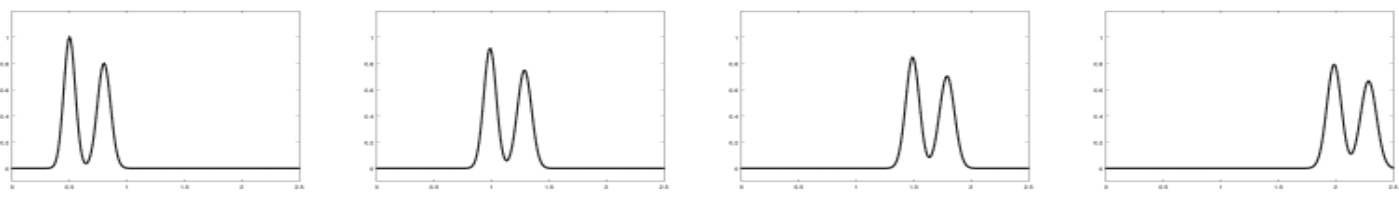
$$u_t + u_x = 0, \quad u(0, x) = e^{-200(x-\frac{1}{2})^2} + 0.8 e^{-150(x-0.8)^2} =: u_0(x)$$

$$\Rightarrow u(t, x) = u_0(x-t)$$

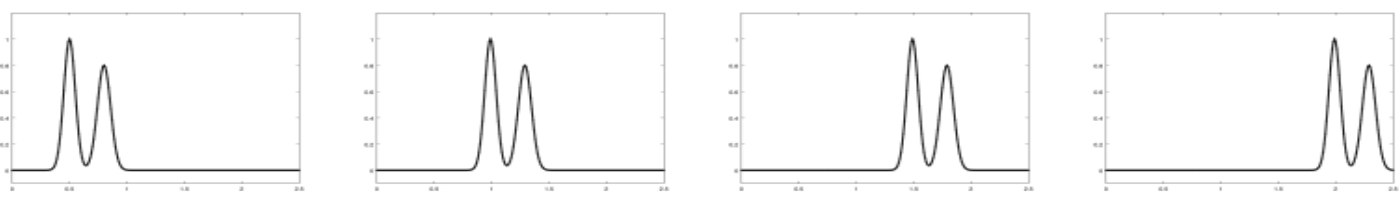
forward
diffmas



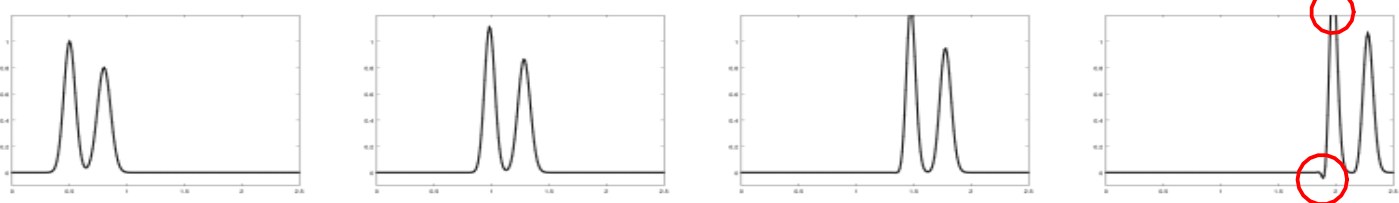
$\frac{\Delta t}{\Delta x} = 0.9$



$\frac{\Delta t}{\Delta x} = 1$



$\frac{\Delta t}{\Delta x} = 1.1$



Finite differences: hyperbolic PDEs (1D transport equation)

Lebesgue-spaces

Def. Let $\Omega \subset \mathbb{R}^n$ open, $p \in [1, \infty]$. $\|f\|_{L^p} = \begin{cases} \left(\int_{\Omega} |f|^p dx \right)^{\frac{1}{p}} & p \neq \infty \\ \operatorname{ess\,sup}_{x \in \Omega} |f(x)| = \inf \tilde{p} = f \text{ a.e. } \sup_{x \in \Omega} |\tilde{f}(x)| & p = \infty \end{cases}$

$$L^p(\Omega) = \{f: \Omega \rightarrow \mathbb{R} \text{ measurable} \mid \|f\|_{L^p} < \infty\}$$

Thm. (Hölder-inequality) For $p, q \in [1, \infty]$, $\frac{1}{p} + \frac{1}{q} = 1$, $f \in L^p(\Omega)$, $g \in L^q(\Omega)$ we have

$$\int_{\Omega} |fg| dx \leq \|f\|_{L^p} \|g\|_{L^q}$$

Pf. \because For $a, b \geq 0$, $\lambda \in (0, 1)$ one has

$$a^\lambda b^{(1-\lambda)} \leq \lambda a + (1-\lambda)b \quad (\text{geometric} \leq \text{arithmetic mean})$$

since $f(t) = t^\lambda - \lambda t + \lambda - 1$ satisfies $f(t) \leq f(1) = 0 \Rightarrow t^\lambda \leq \lambda t + 1 - \lambda$; now pick $t = \frac{a}{b}$

$$\cdot \text{ choose } a = \left(\frac{|f(x)|}{\|f\|_{L^p}} \right)^p, \quad b = \left(\frac{|g(x)|}{\|g\|_{L^q}} \right)^q, \quad \lambda = \frac{1}{p}, \quad 1-\lambda = \frac{1}{q}$$

$$\Rightarrow \frac{|f(x)g(x)|}{\|f\|_{L^p} \|g\|_{L^q}} \leq \frac{1}{p} \left(\frac{|f(x)|}{\|f\|_{L^p}} \right)^p + \frac{1}{q} \left(\frac{|g(x)|}{\|g\|_{L^q}} \right)^q; \text{ now integrate on } \Omega. \quad \square$$

Finite differences: hyperbolic PDEs (1D transport equation)

Lp-stability

$$u_t + a(t,x)u_x + c(t,x)u = f(t,x) \text{ on } [0,T] \times \Omega, \quad p \in [1, \infty)$$

Thm: Let $u_0(x), f(x) = 0$ for $|x|$ big enough, $\Omega = \mathbb{R}$, sol. $u \geq 0$, then

$$\|u(t, \cdot)\|_{L^p} \leq \|u_0\|_{L^p} \exp\left(\int_0^t \|c(s, \cdot)\|_{L^\infty} + \|a_x(s, \cdot)\|_{L^0} ds\right) + \int_0^t \|f(\tau, \cdot)\|_{L^p} \exp\left(\int_\tau^t \|c(s, \cdot)\|_{L^\infty} + \|a_x(s, \cdot)\|_{L^0} ds\right) d\tau$$

$$\mathcal{P} : p \|u\|_{L^p}^{p-1} \frac{d}{dt} \|u\|_{L^p} = \frac{d}{dt} \|u\|_{L^p}^p = p \int_{\Omega} u^{p-1} (f - a u_x - c u) dx$$

$$\leq p \left[\|u\|_{L^p}^p \|c\|_{L^\infty} + \|u\|_{L^p}^{p-1} \|f\|_{L^p} \right] - \underbrace{\int_{\Omega} a (u^p)_x dx}_{= - \int_{\Omega} a_x u^p dx} \leq p \left[\|u\|_{L^p}^p \|c\|_{L^\infty} + \|u\|_{L^p}^{p-1} \|f\|_{L^p} + \|a_x\|_{L^\infty} \|u\|_{L^p}^p \right]$$

$$\Rightarrow \frac{d}{dt} \|u\|_{L^p} \leq (\|c\|_{L^\infty} + \|a_x\|_{L^\infty}) \|u\|_{L^p} + \|f\|_{L^p}; \text{ now use Gronwall} \quad \square$$

Thm: For a, c constant, $\Omega = \mathbb{R}$, one has $u(t, x) = e^{-ct} u_0(x - at) + \int_0^t e^{-c(t-s)} f(s, x - a(t-s)) ds$

$$\Rightarrow \|u(t, \cdot)\|_{L^p} \leq e^{-ct} \|u_0\|_{L^p} + \int_0^t e^{-c(t-s)} \|f(s, \cdot)\|_{L^p} ds \quad \forall p \in [1, \infty]$$

Finite differences: hyperbolic PDEs (1D transport equation)

l2-stability

$$u_t + a u_x + c u = f, \quad a, c \text{ constant}, a \geq 0, \Omega = \mathbb{R}$$

$$\frac{u_j^{m+1} - u_j^m}{\Delta t} + a \frac{u_j^m - u_{j-1}^m}{\Delta x} + c u_j^m = F_j^m$$

Thm: (l^2 -stability) Let u_j^m be the soln. of the upwind method, $\mu = a \frac{\Delta t}{\Delta x}$. One has

$$\|u^m\|_{l^2} \leq \begin{cases} e^{-c t^m} \|u^0\|_{l^2} + \Delta t \sum_{n=0}^{m-1} \|F^n\|_{l^2} e^{-c(t^m - t^n)} & \text{if } \mu \leq 1 - c \Delta t \text{, } l^2\text{-stable''} \\ e^{(c+2k)t^m} \|u^0\|_{l^2} + \Delta t \sum_{n=0}^{m-1} \|F^n\|_{l^2} e^{(c+2k)(t^m - t^n)} & \text{if } \mu \leq 1 + k \Delta t \text{, } l^2\text{-von Neumann-stable''} \end{cases}$$

Pf: \therefore Set $u_j^m = \frac{1}{2\pi} \int_{-\pi/\Delta x}^{\pi/\Delta x} \hat{u}^m(k) e^{ikj\Delta x} dk$

$$\Rightarrow \frac{\hat{u}^{m+1}(k) - \hat{u}^m(k)}{\Delta t} + a \frac{\hat{u}^m(k) - e^{-ik\Delta x} \hat{u}^m(k)}{\Delta x} + c \hat{u}^m(k) = \hat{F}^m(k) \quad \forall k \in [-\pi/\Delta x, \pi/\Delta x]$$

$$\Rightarrow \hat{u}^{m+1}(k) = (1 - \mu(1 - e^{-ik\Delta x}) - c\Delta t) \hat{u}^m(k) + \Delta t \hat{F}^m(k) =: \lambda(k) \hat{u}^m(k) + \Delta t \hat{F}^m(k)$$

Let $\Lambda = \max_k |\lambda(k)| \Rightarrow \|u^{m+1}\|_{l^2} = \frac{1}{\sqrt{2\pi}} \|\hat{u}^{m+1}\|_{l^2} \leq \Lambda \frac{1}{\sqrt{2\pi}} \|\hat{u}^m\|_{l^2} + \Delta t \frac{1}{\sqrt{2\pi}} \|\hat{F}^m\|_{l^2}$

for $a < 0$ this would be $\approx 1 + c|\mu|!!$ $= \Lambda \|u^m\|_{l^2} + \Delta t \|F^m\|_{l^2} \leq \Lambda^{m+1} \|u^0\|_{l^2} + \sum_{n=0}^m \Delta t \Lambda^{m-n} \|F^n\|_{l^2}$

$|\lambda(k)|^2 = (1 - c\Delta t)^2 + 4\mu \left(\sin^2 \frac{k\Delta x}{2}\right) (\mu - (1 - c\Delta t)) \Rightarrow \Lambda \leq \begin{cases} 1 - c\Delta t & \leq e^{-c\Delta t} & \text{if } \mu \leq 1 - c\Delta t \\ 1 + (c+2k)\Delta t & \leq e^{(c+2k)\Delta t} & \text{if } \mu \leq 1 + k\Delta t \end{cases}$

Finite differences: hyperbolic PDEs (1D transport equation)

l^∞ -stability

$$u_t + a u_x + c u = f, \quad a \geq 0, \quad \Omega = \mathbb{R}$$

$$\frac{u_j^{m+1} - u_j^m}{\Delta t} + A_j^m \frac{u_j^m - u_{j-1}^m}{\Delta x} + C_j^m u_j^m = F_j^m$$

Thm: (l^∞ -stability) Let u_j^m be the soln. of the upwind method, $m \geq 0, j \in \mathbb{Z}, \mu = \|A\|_{l^\infty} \frac{\Delta t}{\Delta x}$.

$$\text{If } \mu \leq 1, \quad \|u^m\|_{l^\infty} \leq e^{t^m \|C\|_{l^\infty}} \|u^0\|_{l^\infty} + \Delta t \sum_{k=0}^{m-1} e^{(t^{m-1} - t^k) \|C\|_{l^\infty}} \|F^k\|_{l^\infty}.$$

$$\text{Pf: } u_j^{m+1} = u_j^m \left(1 - A_j^m \frac{\Delta t}{\Delta x}\right) + A_j^m \frac{\Delta t}{\Delta x} u_{j-1}^m - \Delta t C_j^m u_j^m + \Delta t F_j^m$$

$$\Rightarrow |u_j^{m+1}| \leq \|u^m\|_{l^\infty} \left(1 - A_j^m \frac{\Delta t}{\Delta x}\right) + A_j^m \frac{\Delta t}{\Delta x} \|u^m\|_{l^\infty} + \Delta t \|C\|_{l^\infty} \|u^m\|_{l^\infty} + \Delta t \|F^m\|_{l^\infty}$$

$$\Rightarrow \|u^{m+1}\|_{l^\infty} \leq (1 + \|C\|_{l^\infty} \Delta t) \|u^m\|_{l^\infty} + \Delta t \|F^m\|_{l^\infty}$$

\Rightarrow now induction in m

□

Thm: For $C = \text{constant}$ & $\mu + C \Delta t \leq 1$ even $\|u^m\|_{l^\infty} \leq e^{-c t^m} \|u^0\|_{l^\infty} + \Delta t \sum_{k=0}^{m-1} e^{-c(t^{m-1} - t^k)} \|F^k\|_{l^\infty}$

Pf: homework □

How about $\Omega = [l, r]$ and Dirichlet b.c. in l ?

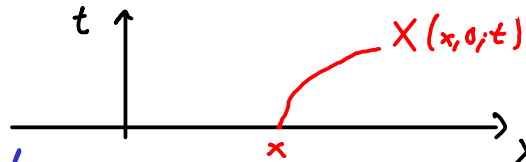
Finite differences: hyperbolic PDEs (1D transport equation)

Domain of Dependence

$$u_t(t,x) + a(t,x)u_x(t,x) + c(t,x)u(t,x) = f(t,x)$$

Def: The curve defined by $\frac{d}{dt} X(x,s;t) = a(t, X(x,s;t))$, $X(x,s;s) = x$ is called characteristic curve of the PDE.

Ex: $u_t + tu_x = 0 \Rightarrow \frac{d}{dt} X(x,s;t) = t \Rightarrow X(x,s;t) = x + \frac{t^2 - s^2}{2}$



We have $\frac{d}{dt} u(t, X(x,s;t)) = u_t(t, X(x,s;t)) + u_x(t, X(x,s;t)) \frac{d}{dt} X(x,s;t) = -c(t, X(x,s;t))u(t, X(x,s;t)) + f(t, X(x,s;t))$

With $\tilde{g}(t) := g(t, X(x,s;t))$ we obtain $\tilde{u}' = -\tilde{c}\tilde{u} + \tilde{f} \Rightarrow$ Along char. curve one has an ODE
 $\Rightarrow \tilde{u}(t) = \tilde{u}(s)e^{-\int_s^t \tilde{c}(\tau) d\tau} + \int_s^t \tilde{f}(\tau) e^{-\int_\tau^t \tilde{c}(\tau) d\tau} d\tau$ only depends on c, f, u along $(t, X(x,s;t))!$

Ex: $u_t + tu_x = 0 \Rightarrow u(t, x + \frac{t^2 - s^2}{2}) = u(s, x) \Rightarrow u(t, x) = u(0, x - \frac{t^2}{2})$

Def: The analytical domain of dependence (dod) of (t,x) is $\{X(x,s;t), t \mid 0 \leq t \leq s\}$.

Courant-Friedrichs-Lewy-condition: analytical & numerical dod 

$\Rightarrow \alpha \frac{\Delta t}{\Delta x} \leq 1$

Finite differences: hyperbolic PDEs (1D transport equation)

Convergence (consistency + stability)

Def: Let u solve $u_t + a u_x + c u = f$. The truncation error of the upwind method is

$$\tau_j^m = \frac{u_j^{n+1} - u_j^m}{\Delta t} + a_j^m \frac{u_j^m - u_{j-1}^m}{\Delta x} + c_j^m u_j^m - f_j^m \quad (\text{with } g_j^m = g(t^m, x_j))$$

Thm: (consistency) The upwind method is first order consistent, i.e. a smooth soln u satisfies $\tau_j^m = O(\Delta t + \Delta x)$.

Pf : homework (via Taylor expansion) \square

Just like for the heat eq the error solves $e_j^m = u_j^m - u_j^m$: $\frac{e_j^{n+1} - e_j^m}{\Delta t} + a_j^m \frac{e_j^m - e_{j-1}^m}{\Delta x} + c_j^m e_j^m = \tau_j^m$
we only prove ℓ^2 -convergence for $a, c = \text{const}$

Thm: (Convergence) Let $U_j^m, j \in \mathbb{Z}, m = 0, \dots, M = \frac{T}{\Delta t}$ solve the upwind-method for $u_t + a u_x + c u = f$ on $[0, T] \times \Omega, \Omega = \mathbb{R}, a \geq 0$ + i.c., u smooth, $\mu \leq 1$. Then $\|e^m\|_{\ell^2}, \|e^m\|_{\ell^\infty} = O(\Delta t + \Delta x) \forall m$

Pf : Stability $\Rightarrow \|e^m\| \leq \Delta t \sum_{k=0}^{m-1} e^{\text{const} \cdot T} \|T^m\| \leq T e^{\text{const} \cdot T} O(\Delta t + \Delta x)$. \square
only on that boundary where characteristics enter

Thm: Now let $\Omega = [l, r]$ & use b.c. $u(t, l) = g(t)$. $\|e^m\|_{\ell^\infty} = O(\Delta t + \Delta x), m = 1, \dots, M$.

Pf : homework (first derive ℓ^∞ -stability for 0-b.c.) \square

Finite differences: hyperbolic PDEs (1D transport equation)

Methods of higher order consistency

$$u_t + au_x = 0$$

Def: method of central differences

$$\frac{U_j^{m+n} - U_j^m}{\Delta t} + A_j \frac{U_{j+n}^m - U_{j-n}^m}{2\Delta x} = 0$$

Lax-Friedrichs-method

$$\frac{U_j^{m+n} - \frac{1}{2}(U_{j-n}^m + U_{j+n}^m)}{\Delta t} + A_j \frac{U_{j+n}^m - U_{j-n}^m}{2\Delta x} = 0$$

Lax-Wendroff-method

$$\frac{U_j^{m+n} - U_j^m}{\Delta t} + a \frac{U_{j+n}^m - U_{j-n}^m}{2\Delta x} = \frac{1}{2} \Delta t a^2 \frac{U_{j+n}^m - 2U_j^m + U_{j-n}^m}{\Delta x^2}$$

sign of a plays no role any more!

for $a = \text{const.}$

What is the Lax-Wendroff-method based on?

$$u_t = -au_x \Rightarrow u_{tt} = -a_t u_x - au_{xt} = -a_t u_x - a(-au_x)_x$$

$$\text{Taylor expansion} \Rightarrow U_j^{m+n} = U_j^m + \Delta t [u_t]_j^m + \frac{\Delta t^2}{2} [u_{tt}]_j^m + O(\Delta t^3)$$

$$= -[au_x]_j^m = [a_t u_x - a(au_x)_x]_j^m$$

$$\Rightarrow \frac{U_j^{m+n} - U_j^m}{\Delta t} + [(a + \frac{1}{2} \Delta t a_t) u_x]_j^m = \frac{1}{2} \Delta t [a(au_x)_x]_j^m + O(\Delta t^2)$$

Thm: (Consistency) The truncation error of the three methods for a smooth soln u is $O(\Delta t + \Delta x^2)$, resp. $O(\Delta t^2 + \Delta x^2)$ for LW.

Pf: homework \square

Finite differences: hyperbolic PDEs (1D transport equation)

l2-stability

$$u_t + au_x = 0 \text{ on } \mathbb{R} \text{ with } a = \text{const.}$$

If we insert $u_j^m = \frac{1}{2\pi} \int_{-\pi/\Delta x}^{\pi/\Delta x} \hat{u}^m(k) e^{ikx_j} dk$ into the methods, we obtain

centr. diff:
$$\frac{\hat{u}^{m+1}(k) - \hat{u}^m(k)}{\Delta t} + a \frac{e^{ik\Delta x} \hat{u}^m(k) - e^{-ik\Delta x} \hat{u}^m(k)}{2\Delta x} = 0$$

Lax-Friedrichs:
$$\frac{\hat{u}^{m+1}(k) - \frac{1}{2}(e^{ik\Delta x} \hat{u}^m(k) + e^{-ik\Delta x} \hat{u}^m(k))}{\Delta t} + a \frac{e^{ik\Delta x} \hat{u}^m(k) - e^{-ik\Delta x} \hat{u}^m(k)}{2\Delta x} = 0$$

Lax-Wendroff:
$$\frac{\hat{u}^{m+1}(k) - \hat{u}^m(k)}{\Delta t} + a \frac{e^{ik\Delta x} \hat{u}^m(k) - e^{-ik\Delta x} \hat{u}^m(k)}{2\Delta x} = \frac{1}{2} \Delta t a^2 \frac{\hat{u}^m(k) (e^{ik\Delta x} - 2 + e^{-ik\Delta x})}{\Delta x^2}$$

$\Rightarrow \hat{u}^{m+1}(k) = \lambda(k) \hat{u}^m(k)$ for all $k \in [-\frac{\pi}{\Delta x}, \frac{\pi}{\Delta x}]$ with

$$\lambda(k) = \begin{cases} 1 - \frac{\mu}{2}(e^{ik\Delta x} - e^{-ik\Delta x}) = 1 - \mu i \sin k\Delta x & , \text{ CD} \\ \frac{1}{2}[e^{ik\Delta x} + e^{-ik\Delta x} - \mu(e^{ik\Delta x} - e^{-ik\Delta x})] = \cos k\Delta x - \mu i \sin k\Delta x & , \text{ LF} \\ 1 - \frac{\mu}{2}(e^{ik\Delta x} - e^{-ik\Delta x}) + \frac{\mu^2}{2}(e^{ik\Delta x} - 2 + e^{-ik\Delta x}) = 1 - \mu i \sin k\Delta x - 2\mu^2 \sin^2 \frac{k\Delta x}{2} & , \text{ LW} \end{cases}$$

Thm: (l^2 -stability) LF & LW are l^2 -stable for $|\mu| \leq 1$, CD is never l^2 -stable.

Pf : LF: $|\lambda(k)|^2 = \cos^2 k\Delta x + \mu^2 \sin^2 k\Delta x \leq 1$; LW: $|\lambda(k)|^2 = (1 - 2\mu^2 \sin^2 \frac{k\Delta x}{2})^2 + \mu^2 \sin^2 k\Delta x = 1 - 4\mu^2(1 - \mu^2) \sin^4 \frac{k\Delta x}{2} \leq 1$

CD: $|\lambda(k)|^2 = 1 + \mu^2 \sin^2 k\Delta x > 1 \Rightarrow |\hat{u}^m(k)| \geq (1 + c)^{\frac{m}{2}} |\hat{u}^0(k)|$ for $\forall c \leq \mu^2 \sin^2(k\Delta x)$

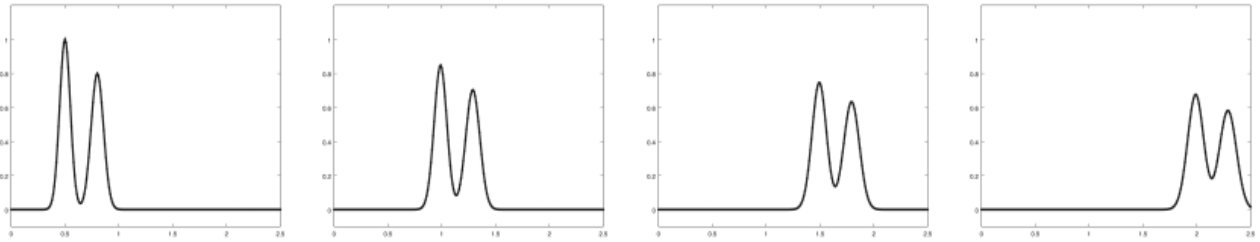
$\Rightarrow \Delta t \xrightarrow{\mu \text{ const.}} 0$ implies $\sqrt{2\pi} \|u^m\|_{l^2} = \|\hat{u}^m\|_{l^2} \geq \|\hat{u}^m\|_{\left[\frac{\pi}{4\Delta x}, \frac{3\pi}{4\Delta x}\right]}_{l^2} \geq (1 + \frac{c}{2})^{\frac{m}{2}} \|\hat{u}^0\|_{\left[\frac{\pi}{4\Delta x}, \frac{3\pi}{4\Delta x}\right]}_{l^2} \rightarrow \infty$

Finite differences: hyperbolic PDEs (1D transport equation)

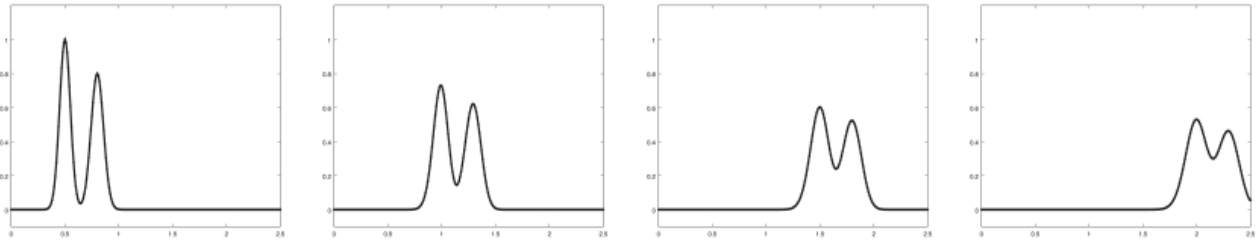
Example simulation

$$u_t + u_x = 0, \quad \mu = \frac{\Delta t}{\Delta x} = 0.7$$

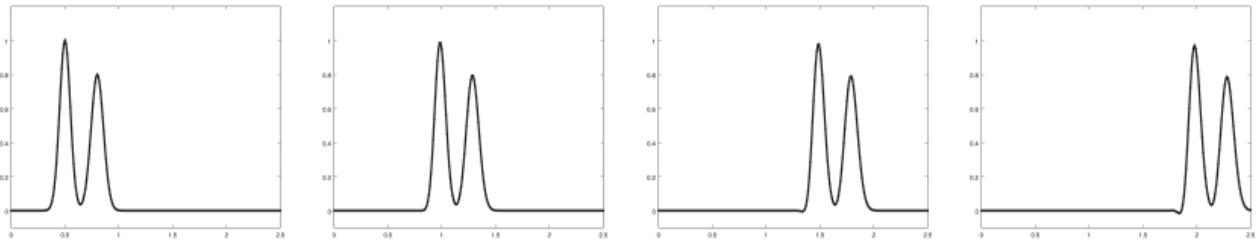
upwind



LF



LW



All methods are convergent in Q^2 (follows from stability + consistency), but ...

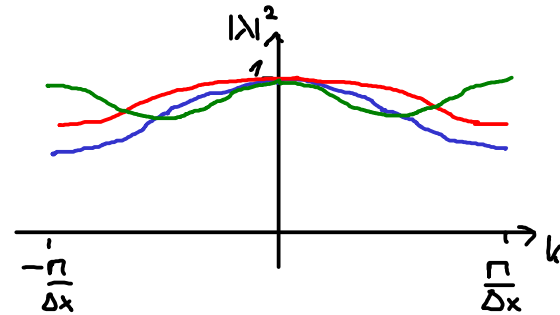
Finite differences: hyperbolic PDEs (1D transport equation)

Numerical dissipation

$$u_t + u_x = 0 \text{ has soln. } u(t, x) = u(0, x-t) \Rightarrow \|u\|_{L^2} = \text{const.}$$

$$\sqrt{2\pi} \|u^m\|_{L^2} = \|\hat{u}^m\|_{L^2} \text{ with } |\hat{u}^m(k)| \leq |\lambda(k)|^m |\hat{u}^0(k)|$$

$$\text{for } |\lambda(k)|^2 = \begin{cases} 1 - 4\mu(1-\mu) \sin^2 \frac{k\Delta x}{2} & \text{upwind} \\ 1 - (1-\mu^2) \sin^2 k\Delta x & \text{LF} \\ 1 - 4\mu^2(1-\mu^2) \sin^4 \frac{k\Delta x}{2} & \text{LW} \end{cases}$$



\Rightarrow frequency components at $k=0$ stay unchanged due to $|\lambda(0)|=1$, but frequency components further away from $k=0$ are damped \Rightarrow "numerical dissipation"

Def: A scheme is called dissipative of order $r \in \mathbb{N}$ if there is $\delta > 0$ such that

$$|\lambda(k)|^2 \leq 1 - \delta \left| \sin \frac{k\Delta x}{2} \right|^{2r} \text{ near } k=0. \quad (\text{upwind \& LF: } r=1, \text{ LW: } r=2)$$

Finite differences: hyperbolic PDEs (1D transport equation)

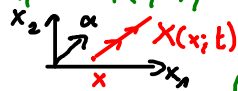
Perspective: 2D and nonlinearity

Conservation principle: $u_t + \operatorname{div}[q(x,u)] = 0 \Rightarrow u_t + q_u \cdot u_x + \operatorname{tr} q_x = 0, \Omega \subset \mathbb{R}^2$

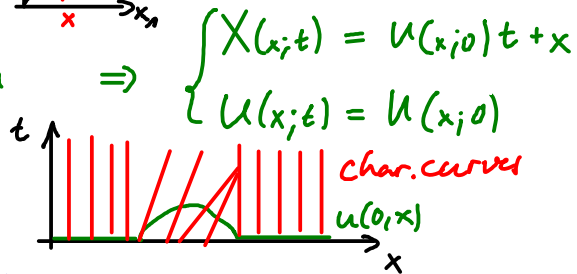
\Rightarrow general form: $u_t(t,x) + a(t,x,u(t,x)) \cdot \nabla u(t,x) = f(t,x,u(t,x))$

Solu. satisfies $u(t, X(x;t)) = U(x;t)$
 for characteristic equations $\begin{cases} \dot{X}(x;t) = a(t, X(x;t), U(x;t)), & X(x;0) = x \\ \dot{U}(x;t) = f(t, X(x;t), U(x;t)), & U(x;0) = u_0(x) \end{cases}$
ODE, i.e. concept of d.o.d. also applies here

Ex: $u_t + a \cdot \nabla u = 0, a \in \mathbb{R}^2 \Rightarrow X(x;t) = at + x, u(t, X(x;t)) = U(x;t) = U(x;0), \text{d.h. } u(t,y) = u(0, y-at)$
char. curve



Ex: (Burgers equation) $u_t + (\frac{u^2}{2})_x = 0, \text{ i.e. } a(t,x,u) = u$
 $\Rightarrow u$ is constant along char. curves



The upwind method now reads

$$\frac{U_{ij}^{m+n} - U_{ij}^m}{\Delta t} + \frac{1}{\Delta x} \left[a(t^m, x_{ij}, U_{ij}^m) \cdot \begin{pmatrix} U_{ij}^m - U_{i-1,j}^m \\ U_{ij}^m - U_{i,j-1}^m \end{pmatrix} + a(t^m, x_{ij}, U_{ij}^m) \cdot \begin{pmatrix} U_{i+1,j}^m - U_{ij}^m \\ U_{ij}^{m+n} - U_{ij}^m \end{pmatrix} \right] = f(t^m, x_{ij}, U_{ij}^m)$$

+ i.c. + b.c. on that part of $\partial\Omega$ where char. curves enter

If $|a(t,x,u) \cdot (\frac{\Delta t}{\Delta x})| \leq 1 \forall t,x,u$, then l^2 - l^∞ stability follows as before.

Finite Elements: elliptic PDEs (2D Poisson equation)

Elliptic PDE theory: introduction to Euler-Lagrange equations

Poisson equation with homogeneous Dirichlet boundary conditions:

$$\begin{cases} \Delta u = f & \text{on } \Omega \subset \mathbb{R}^d \\ u = 0 & \text{on } \partial\Omega \end{cases}$$

We had interpreted this as necessary condition for the fact that u minimizes

$$E(u) = \int_{\Omega} \frac{1}{2} |\nabla u(x)|^2 - f(x) u(x) dx \quad (*)$$

among all functions $\Omega \rightarrow \mathbb{R}$ with 0-b.c. In detail:

Def: (Gateaux derivative) Let X be a vector space, $E: X \rightarrow \mathbb{R}$. The Gateaux derivative of E in $u \in X$ in direction $v \in X$ is $\partial_u E(u)(v) = \left. \frac{d}{dt} E(u+tv) \right|_{t=0}$.

Thm: (Necessary optimality condition) If u minimizes the functional E and if $\partial_u E(u)(v)$ exists, then $\partial_u E(u)(v) = 0$. *why?*

Ex: For (*) we have $\partial_u E(u)(v) = \int_{\Omega} \nabla u \cdot \nabla v - f v dx$. If u is sufficiently often differentiable, using integration by parts we obtain $\partial_u E(u)(v) = \int_{\Omega} v (-\Delta u - f) dx$. If u is a minimum, then $\partial_u E(u)(v) = 0$ for "all" $v \Rightarrow -\Delta u - f = 0$.

Actually, only the so-called weak form of the PDE holds

$$0 = \partial_u E(u)(v) = \int_{\Omega} \nabla u \cdot \nabla v - f v dx \quad \forall v: \Omega \rightarrow \mathbb{R}, v|_{\partial\Omega} = 0.$$

How often must u be differentiable to interpret this as $-\Delta u = f$?

$\Omega \subset \mathbb{R}^d$

Elliptic PDE theory: weak derivative

Def.: $C^m(\bar{\Omega})$ is the space of m -times continuously differentiable functions $f: \bar{\Omega} \rightarrow \mathbb{R}$ with

$$\|f\|_{C^m} := \sup_{x \in \Omega, |\alpha| \leq m} |D^\alpha f(x)| < \infty$$

$C_0^\infty(\Omega)$ is the space of infinitely often continuously differentiable functions $f: \Omega \rightarrow \mathbb{R}$ with compact support $\text{supp}(f) := \{x \in \Omega \mid f(x) \neq 0\} \subset \Omega$.

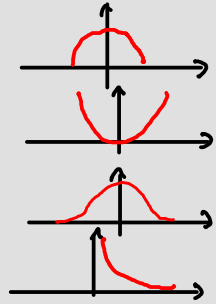
$L^1_{loc}(\Omega) = \{f: \Omega \rightarrow \mathbb{R} \mid f \in L^1(\omega) \forall \bar{\omega} \subset \Omega \text{ kompakt}\} = \text{locally integrable functions}$

Ex.: $x \mapsto \sqrt{1-x^2} \in C^0([-1,1]) \setminus C^1([-1,1])$

$x \mapsto |x|^3 \in C^2([-1,1]) \setminus C^3([-1,1])$

$x \mapsto \begin{cases} \exp(-\frac{1}{1-x^2}) & x \in (-1,1) \\ 0 & \text{else} \end{cases} \in C_0^\infty(\mathbb{R})$

$x \mapsto \frac{1}{x} \in L^1_{loc}((0,\infty))$



Def.: Let $u \in L^1_{loc}(\Omega)$, α a multiindex. $w \in L^1_{loc}$ is called weak α -derivative of u , if $\int_{\Omega} w(x)v(x) dx = (-1)^{|\alpha|} \int_{\Omega} u(x) D^\alpha v(x) dx \quad \forall v \in C_0^\infty(\Omega)$.

Ex.: $u \in C^m(\bar{\Omega}), v \in C_0^\infty(\Omega) \Rightarrow \int_{\Omega} v \partial_{x_i} u dx = - \int_{\Omega} u \partial_{x_i} v dx \Rightarrow \partial_{x_i} u$ is weak derivative with respect to x_i .

analogously, $D^\alpha u$ is weak α -derivative for $|\alpha| \leq m$

$x \mapsto |x|$ has weak derivative $x \mapsto \begin{cases} -1 & x < 0 \\ 1 & \text{else} \end{cases}$ ← check!

Elliptic PDE theory: Sobolev spaces

Def: A metric space X is called complete, if every Cauchy sequence in X converges.

A complete normed space X with inner product $\langle \cdot, \cdot \rangle_X$ is called Hilbert space, if

$$\|x\|_X = \sqrt{\langle x, x \rangle} \quad \forall x \in X$$

Let $m \in \mathbb{N}$, $p \in [1, \infty]$. The (m, p) -Sobolev space is

$$W^{m,p}(\Omega) = \{f \in L^p(\Omega) \mid f \text{ has weak } \alpha\text{-derivative } D^\alpha f \in L^p(\Omega) \forall |\alpha| \leq m\}$$

with norm
$$\|f\|_{W^{m,p}} = \left(\sum_{|\alpha| \leq m} \|D^\alpha f\|_{L^p}^p \right)^{\frac{1}{p}}$$

Let $m \in \mathbb{N}$. The m -Hilbert space is $H^m(\Omega) = W^{m,2}(\Omega)$

with inner product
$$(f, g)_{H^m} = \sum_{|\alpha| \leq m} \int_{\Omega} D^\alpha f D^\alpha g \, dx.$$

Thm: $W^{m,p}(\Omega)$ is complete.

Ex: $x \mapsto \frac{1}{x^2} \in W^{m,p}((1, \infty)) \quad \forall p \in [1, \infty], m \geq 0$

$x \mapsto \sqrt{|x|} \in (W^{1,p}((-1, 1)) \cap L^q((-1, 1))) \setminus W^{2,r}((-1, 1)) \quad \forall p \in [1, 2), q, r \in [1, \infty]$

$x \mapsto \begin{cases} -1 & x \leq 0 \\ 1 & \text{else} \end{cases} \notin W^{1,p}((-1, 1)) \quad \forall p$

Thm: (trace theorem) Let Ω have Lipschitz boundary. $T: W^{1,p}(\Omega) \rightarrow L^p(\partial\Omega)$, $u \mapsto u|_{\partial\Omega}$ is linear & bounded.

Finite Elements: elliptic PDEs (2D Poisson equation)

Elliptic PDE theory: linear elliptic boundary value problems

Apparently, the minimization of $E(u) = \int_{\Omega} \frac{|\nabla u|^2}{2} - \int_{\Omega} f u \, dx$ already makes sense on $H^1(\Omega) \supset C^2(\bar{\Omega})$.

Def: $u \in H_0^1(\Omega) = \{v \in H^1(\Omega) \mid v=0 \text{ on } \partial\Omega\}$ is called weak solution of $-\Delta u = f$ in Ω , $u=0$ on $\partial\Omega$, if $0 = \partial_u E(u)(v) = \int_{\Omega} \nabla u \cdot \nabla v - f v \, dx \quad \forall v \in H_0^1(\Omega)$.

To obtain more general elliptic PDEs and boundary conditions, we now consider

$$E(u) = \int_{\Omega} \frac{1}{2} \nabla u(x)^T A(x) \nabla u(x) + \frac{q(x)}{2} u(x)^2 - f(x) u(x) \, dx + \int_{\Gamma_1} \frac{\rho(x)}{2} u(x)^2 - g(x) u(x) \, dx$$

for $A(x) \in \mathbb{R}^{d \times d}$ continuously differentiable symmetric positive definite, $q \in L^\infty(\Omega)$, $q \geq 0$, $\Gamma_1 \subset \partial\Omega$ with positive measure, $\rho \in L^\infty(\Gamma_1)$, $\rho \geq 0$, $f \in L^2(\Omega)$, $g \in L^2(\Gamma_1)$, to be minimized for $u \in H^1(\Omega)$ (why does this make sense?) with $u = h \in L^2(\Gamma_2)$ on $\Gamma_2 = \partial\Omega \setminus \Gamma_1$.

\Rightarrow For all $v \in H^1(\Omega)$ with $v|_{\Gamma_2} = 0$ we have

$$0 = \partial_u E(u)(v) = \int_{\Omega} \nabla v^T A \nabla u + q u v - f v \, dx + \int_{\Gamma_1} \rho u v - g v \, dx \quad \text{"weak form"}$$

if u smooth $\nearrow \int_{\Omega} v [-\operatorname{div}(A \nabla u) + q u - f] \, dx + \int_{\Gamma_1} (n^T A \nabla u + \rho u - g) v \, dx$

$\Rightarrow -\operatorname{div}(A \nabla u) + q u = f$ in Ω with $\begin{cases} u = h & \text{on } \Gamma_2 & \text{Dirichlet bc} \\ n^T A \nabla u = g & \text{on } \Gamma_1 \setminus \{\rho=0\} & \text{Neumann bc} \\ n^T A \nabla u + \rho u = g & \text{on } \Gamma_1 & \text{Robin bc} \end{cases}$

\Rightarrow exactly one bc everywhere on $\partial\Omega$

Elliptic PDE theory: weak solution

The weak form of the PDE is obtained from the energy E as $0 = \partial_u E(u)(v)$ for all allowed v ; from the PDE itself it is obtained by multiplication of the PDE with v & integration by parts.

Ex: $\Delta u = f$ in Ω , $u = 0$ on $\partial\Omega \Rightarrow$ allowed v are $v \in H_0^1(\Omega)$

$$\rightarrow \int_{\Omega} v \Delta u \, dx = \int_{\Omega} v f \, dx \Rightarrow \int_{\Omega} \nabla u \cdot \nabla v \, dx = \int_{\Omega} v f \, dx \quad \forall v \in H_0^1(\Omega) \text{ is weak form}$$

Def: The weak form of
$$\begin{cases} -\operatorname{div}(A(x) \nabla u(x)) + b(x) \cdot \nabla u(x) + q(x) u(x) = f(x) & \text{in } \Omega \\ u(x) = h(x) & \text{on } \Gamma_2 = \partial\Omega \setminus \Gamma_1 \\ (A(x) \nabla u(x)) \cdot n + p(x) u(x) = g(x) & \text{on } \Gamma_1 \end{cases}$$

$$\text{is } 0 = \int_{\Omega} \nabla v^T A \nabla u + v b \cdot \nabla u + q u v - f v \, dx + \int_{\Gamma_1} p u v - g v \, dx \quad \forall v \in H_{\Gamma_2}^1(\Omega) = \{v \in H^1(\Omega) \mid v|_{\Gamma_2} = 0\}$$

& $u = h$ on Γ_2

$u \in H^1(\Omega)$ is called a weak solution of the PDE if it satisfies its weak form.

Finite Elements: elliptic PDEs (2D Poisson equation)

Elliptic PDE theory: associated bilinear form

By setting $u(x) = \tilde{u}(x) + h(x) \quad \forall x \in \Omega$ one can reduce the bc $u|_{\Gamma_2} = h$ to $\tilde{u}|_{\Gamma_2} = 0$.

Ex: $-\Delta u = 0$ in Ω , $u = h$ on $\partial\Omega \Rightarrow$ Find $u \in H^1(\Omega)$ with $\int_{\Omega} \nabla u \cdot \nabla v \, dx = 0 \quad \forall v \in H_0^1(\Omega)$ & $u|_{\Gamma_2} = h$

\Rightarrow Find $\tilde{u} \in H_0^1(\Omega)$ with $\int_{\Omega} \nabla \tilde{u} \cdot \nabla v + \nabla h \cdot \nabla v \, dx = 0 \quad \forall v \in H_0^1(\Omega)$

Therefore we only consider 0-boundary values. Hence, let u be a weak soln, i.e.,

$u \in H_{\Gamma_2}^1$ with $0 = \int_{\Omega} \nabla v^T A \nabla u + v b \cdot \nabla u + q u v - f v \, dx + \int_{\Gamma_1} p u v - g v \, dx \quad \forall v \in H_{\Gamma_2}^1(\Omega)$.

Set $a(u, v) = \int_{\Omega} \nabla v^T A \nabla u + v b \cdot \nabla u + q u v \, dx + \int_{\Gamma_1} p u v \, dx$ "associated bilinear form"

$$l(v) = \int_{\Omega} f v \, dx + \int_{\Gamma_1} g v \, dx$$

a is a bilinear form, i.e. linear in u & v ($a(u + \lambda \tilde{u}, v) = a(u, v) + \lambda a(\tilde{u}, v)$, analogous in v)

l is a linear form ($l(v + \lambda \tilde{v}) = l(v) + \lambda l(\tilde{v})$)

u is weak soln. of PDE

$$\Leftrightarrow u \in H_{\Gamma_2}^1 \quad \& \quad a(u, v) = l(v) \quad \forall v \in H_{\Gamma_2}^1$$

Finite Elements: elliptic PDEs (2D Poisson equation)

Elliptic PDE theory: existence, uniqueness, equivalence to minim.

Def: Let X be a Hilbert space, $a: X \times X \rightarrow \mathbb{R}$ a bilinear form, $l: X \rightarrow \mathbb{R}$ linear.

We require the following conditions: $\exists c, C > 0$ such that for all $u, v \in X$ we have

coercivity/ellipticity: $a(v, v) \geq c \|v\|_X^2$

Boundedness: $|a(u, v)| \leq C \|u\|_X \|v\|_X$

Boundedness: $|l(v)| \leq C \|v\|_X$

Thm: (Lax-Milgram-lemma) Let X be a Hilbert space, a a coercive bounded bilinear form, l linear and bounded. Then there exists a unique solution $u \in X$ of

$$a(u, v) = l(v) \quad \forall v \in X.$$

Thm: Let in addition a be symmetric, i.e., $a(u, v) = a(v, u) \quad \forall u, v \in X$. Then

$$a(u, v) = l(v) \quad \forall v \in X \quad \Leftrightarrow \quad u = \operatorname{argmin}_{\tilde{u} \in X} \frac{1}{2} a(\tilde{u}, \tilde{u}) - l(\tilde{u})$$

Pf : " \Rightarrow " Let $v \in X \Rightarrow \frac{1}{2} a(v, v) - l(v) - (\frac{1}{2} a(u, u) - l(u)) = \frac{1}{2} a(v-u, v-u) + a(v-u, u) - l(v-u)$
 $= \frac{1}{2} a(v-u, v-u) \geq \frac{c}{2} \|v-u\|_X^2$

" \Leftarrow " Set $E(u) = \frac{1}{2} a(u, u) - l(u)$, then $0 = \partial_u E(u)(v) = a(u, v) - l(v) \quad \forall v \in X \quad \square$

Elliptic PDE theory: weak solution to Poisson-problem

$\Omega \subset \mathbb{R}^d$ bounded, Lipschitz

Thm: $a(v, u) = \int_{\Omega} \nabla u \cdot \nabla v \, dx$ is a coercive, bounded bilinear form on $H_0^1(\Omega)$,
 $\ell(v) = \int_{\Omega} f v \, dx$ with $f \in L^2(\Omega)$ is a bounded linear form on $H_0^1(\Omega)$.

Cor: There exists a unique weak solution u of $-\Delta u = f$ in Ω , $u = 0$ on $\partial\Omega$.

Pf (Thm): $\cdot | \ell(v) | \leq \int_{\Omega} | f v | \, dx \leq \| f \|_{L^2} \| v \|_{L^2} \leq \| f \|_{L^2} \sqrt{\| v \|_{L^2}^2 + \| \nabla v \|_{L^2}^2} = \| f \|_{L^2} \| v \|_{H^1}$

$\cdot | a(u, v) | \leq \sum_{i=1}^d \int_{\Omega} | \partial_{x_i} u \partial_{x_i} v | \, dx \leq \sum_{i=1}^d \| \partial_{x_i} u \|_{L^2} \| \partial_{x_i} v \|_{L^2}$
 $\leq \left(\| u \|_{L^2}, \| \partial_{x_1} u \|_{L^2}, \dots, \| \partial_{x_d} u \|_{L^2} \right)^T \cdot \left(\| v \|_{L^2}, \| \partial_{x_1} v \|_{L^2}, \dots, \| \partial_{x_d} v \|_{L^2} \right)^T$
 $\leq \sqrt{\| u \|_{L^2}^2 + \| \partial_{x_1} u \|_{L^2}^2 + \dots + \| \partial_{x_d} u \|_{L^2}^2} \cdot \sqrt{\| v \|_{L^2}^2 + \| \partial_{x_1} v \|_{L^2}^2 + \dots + \| \partial_{x_d} v \|_{L^2}^2} = \| u \|_{H^1} \| v \|_{H^1}$

$\cdot a(v, v) = \int_{\Omega} | \nabla v |^2 \, dx = \frac{1}{2} \| \nabla v \|_{L^2}^2 + \frac{1}{2} \| \nabla v \|_{L^2}^2 \geq \frac{1}{2} \| \nabla v \|_{L^2}^2 + c \| v \|_{L^2}^2 \geq \min\left(\frac{1}{2}, c\right) \| v \|_{H^1}^2$
 $= \| \partial_{x_1} v \|^2 + \dots + \| \partial_{x_d} v \|^2$ Poincaré-inequality!

\cdot bilinearity & linearity obvious (check!) □

Finite Elements: elliptic PDEs (2D Poisson equation)

Elliptic PDE theory: Poincaré inequality

Thm: (Poincaré - (Friedrichs-) inequality) Let $\Omega \subset \mathbb{R}^d$ be bounded, connected with Lipschitz bdy.

Let $\Gamma \subset \partial\Omega$ be simply connected with $|\Gamma| > 0$. There exists $C > 0$ such that

$$\|v\|_{L^2} \leq C \|\nabla v\|_{L^2} \quad \text{for all } v \in H^1_\Gamma = \{v \in H^1(\Omega) \mid v=0 \text{ on } \Gamma\}.$$

Pf: in 1D, wlog. $\Omega = [a, b]$, $v(a) = 0$:

$$\begin{aligned} \|v\|_{L^2}^2 &= \int_a^b v(x)^2 dx = \int_a^b \left(\int_a^x v'(y) dy \right)^2 dx \leq \int_a^b \left(\int_a^b |v'(y)| dy \right)^2 dx = (b-a) \left(\int_a^b |v'(x)| dx \right)^2 \\ &\leq (b-a) \left(\|1\|_{L^2} \|v'\|_{L^2} \right)^2 = (b-a)^2 \|v'\|_{L^2}^2 \end{aligned}$$

in 2D with $\Omega = [a, b] \times [c, d]$, $v(a, x_2) = 0$:

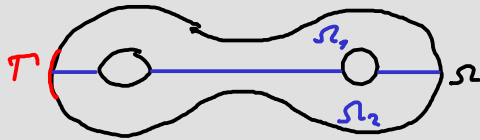
$$\begin{aligned} \|v\|_{L^2}^2 &= \int_c^d \int_a^b v^2 dx = \int_c^d \|v(\cdot, x_2)\|_{L^2}^2 dx_2 \leq (b-a)^2 \int_c^d \|\partial_{x_1} v(\cdot, x_2)\|_{L^2}^2 dx_2 \\ &\leq (b-a)^2 \int_c^d \int_a^b (\partial_{x_1} v)^2 + (\partial_{x_2} v)^2 dx = (b-a)^2 \|\nabla v\|_{L^2}^2 \end{aligned}$$

in 2D, Ω simply connected: $\exists \phi: [a, b] \times [c, d] \rightarrow \Omega$ Lipschitz with ϕ^{-1} Lipschitz, $\phi^{-1}(\Gamma) = \{a\} \times [c, d]$; let L be the Lipschitz constant of ϕ, ϕ^{-1} .

$$\begin{aligned} \|v\|_{L^2}^2 &= \int_\Omega v^2 dx = \int_a^b \int_c^d (v \circ \phi)^2 |\det D\phi| dx \leq 2L^2 \|v \circ \phi\|_{L^2}^2 \leq 2L^2 (b-a)^2 \|\nabla(v \circ \phi)\|_{L^2}^2 = C \|\nabla \phi(\nabla v) \circ \phi\|_{L^2}^2 \\ &\leq C L^2 \|(\nabla v) \circ \phi\|_{L^2}^2 = C L^2 \int_a^b \int_c^d (\nabla v)^2 \phi dx = C L^2 \int_\Omega \nabla v^2 |\det D\phi^{-1}| dx \leq 2CL^4 \|\nabla v\|_{L^2}^2 \end{aligned}$$

Elliptic PDE theory: Poincaré inequality II

- analogous proof in d dimensions for simply connected domain Ω Lipschitz-boundary, $v=0$ on $T \subset \partial\Omega$ simply connected with $|T| > 0$.
- For arbitrary connected Ω with Lipschitz boundary, $v=0$ on $T \subset \partial\Omega$ connected with $|T| > 0$ decompose Ω into Lipschitz domains $\Omega_1, \dots, \Omega_N$ with $v=0$ on part of the boundary



and use $\|v\|_{L^2}^2 = \|v|_{\Omega_1}\|_{L^2}^2 + \dots + \|v|_{\Omega_N}\|_{L^2}^2$,
 $\|\nabla v\|_{L^2}^2 \geq \|\nabla v|_{\Omega_1}\|_{L^2}^2 + \dots + \|\nabla v|_{\Omega_N}\|_{L^2}^2$ □

Thm: The following, non-improvable inequalities hold for all $v \in H_0^1(\Omega)$.

- 1) $\|v\|_{L^2}^2 \leq \left(\frac{b-a}{\pi}\right)^2 \|\nabla v\|_{L^2}^2$ if $\Omega = [a, b]$
- 2) $\|v\|_{L^2}^2 \leq \frac{1}{2\pi^2} \|\nabla v\|_{L^2}^2$ if $\Omega = [0, 1]^2$

Pf : 1) homework: with semidiscrete Fourier-transform

2) reduce to case 1) □

Finite Elements: elliptic PDEs (2D Poisson equation)

Elliptic PDE theory: weak solution to linear elliptic bvp

Thm: Let Ω be Lipschitz, $\Gamma_1, \Gamma_2 = \partial\Omega \setminus \Gamma_1$ with $|\Gamma_1|, |\Gamma_2| > 0$. $A \in L^\infty(\Omega; \mathbb{R}^{d \times d})$ with $\xi^T A \xi \geq c |\xi|^2 \forall \xi \in \mathbb{R}^d$

also other conditions possible, see homework

$B \in L^\infty(\Omega; \mathbb{R}^d)$ with $\|b(x)\| \leq C \leq q(x) \forall x \in \Omega$

$q \in L^\infty(\Omega)$ with $q \geq 0$

$p \in L^0(\Gamma_1)$ with $p \geq 0$

$f \in L^2(\Omega), g \in L^2(\Gamma_1)$

Then $a(u, v) = \int_{\Omega} \nabla v^T A \nabla u + v b \cdot \nabla u + q u v \, dx + \int_{\Gamma_1} p u v \, dx$ is a coercive, bounded bilinear form on $H_{\Gamma_2}^1(\Omega)$, and $l(v) = \int_{\Omega} f v \, dx + \int_{\Gamma_1} g v \, dx$ is bounded & linear.

Cor: $\exists!$ weak solution of
$$\begin{cases} -\operatorname{div}(A(x) \nabla u(x)) + b(x) \cdot \nabla u(x) + q(x) u(x) = f(x) & \text{in } \Omega \\ u(x) = 0 & \text{on } \Gamma_2 \\ (A(x) \nabla u(x)) \cdot n + p(x) u(x) = g(x) & \text{on } \Gamma_1 \end{cases}$$
 trace theorem

Pf (Thm): $|l(v)| \leq \int_{\Omega} |f v| \, dx + \int_{\Gamma_1} |g v| \, dx \leq \|f\|_{L^2} \|v\|_{L^2} + \|g\|_{L^2(\Gamma_1)} \|v\|_{L^2(\Gamma_1)} \leq (\|f\|_{L^2} + C \|g\|_{L^2(\Gamma_1)}) \|v\|_{H^1}$

$|a(u, v)| \leq \|A\|_{L^\infty} \|\nabla v\|_{L^2} \|\nabla u\|_{L^2} + \|b\|_{L^\infty} \|v\|_{L^2} \|\nabla u\|_{L^2} + \|q\|_{L^\infty} \|u\|_{L^2} \|v\|_{L^2} + \|p\|_{L^\infty} \|u\|_{L^2(\Gamma_1)} \|v\|_{L^2(\Gamma_1)} \leq \operatorname{const} \cdot \|v\|_{H^1} \|u\|_{H^1}$

$a(v, v) \geq \int_{\Omega} c |\nabla v|^2 - \|b\| \|v\| |\nabla v| + q v^2 \geq \int_{\Omega} (c - \varepsilon) |\nabla v|^2 + \tilde{c} \left(|\nabla v| - \frac{\|b\|}{2\tilde{c}} |v| \right)^2 \, dx \geq (c - \varepsilon) \|\nabla v\|_{L^2}^2 \geq \operatorname{const} \|v\|_{H^1}^2 \square$

Ritz-Galerkin-method: existence, uniqueness, minimization

Thm: Let $a: X \times X \rightarrow \mathbb{R}$ be a coercive, bounded bilinear form on the Hilbert space X , $l: X \rightarrow \mathbb{R}$ linear and bounded. Let $X_h \subset X$ be a closed subspace. Then there exists a unique solution $u_h \in X_h$ of $a(u_h, v_h) = l(v_h) \quad \forall v_h \in X_h$.

Def: u_h is called Ritz-Galerkin-approximation of the solution $u \in X$ to $a(u, v) = l(v) \quad \forall v \in X$.

Pf: a is coercive & bounded bilinear form on X_h , l is bounded linear form on X_h (check!) \square
 \Rightarrow Lax-Milgram yields unique u_h . \square

Thm: If additionally a is symmetric, then $u_h = \operatorname{arg\,min}_{u \in X_h} \frac{1}{2} a(u, u) - l(u)$.

Pf: analogous to proof in X . \square

The Ritz-Galerkin-approximation u_h thus possesses the same properties as u !

If $\dim X_h < \infty$, then there is a basis $\phi_1, \dots, \phi_N \in X_h$. Each $u_h \in X_h$ can be written as $u_h = \sum_{i=1}^N U_i \phi_i$.

$$\Rightarrow a(u_h, v_h) = (U_1, \dots, U_N) A \begin{pmatrix} v_1 \\ \vdots \\ v_N \end{pmatrix}, \quad l(v_h) = (B_1, \dots, B_N) \begin{pmatrix} v_1 \\ \vdots \\ v_N \end{pmatrix}$$

with $A_{ij} = a(\phi_i, \phi_j)$, $B_i = l(\phi_i)$

\Rightarrow Ritz-Galerkin-approximation solves $A^T \begin{pmatrix} U_1 \\ \vdots \\ U_N \end{pmatrix} = \begin{pmatrix} B_1 \\ \vdots \\ B_N \end{pmatrix}$

Ritz-Galerkin-method: basics of error estimates

$a: X \times X \rightarrow \mathbb{R}$ with $a(v,v) \geq c \|v\|_X^2$ & $|a(u,v)| \leq C \|u\|_X \|v\|_X$, $\ell: X \rightarrow \mathbb{R}$ with $|\ell(v)| \leq \bar{C} \|v\|_X$

$X_h \subset X$, $a(u,v) = \ell(v) \quad \forall v \in X$, $a(u_h, v_h) = \ell(v_h) \quad \forall v_h \in X_h$

Thm: (Galerkin-orthogonality) $a(u_h - u, v_h) = 0 \quad \forall v_h \in X_h$

Pf: Subtract $a(u, v_h) = \ell(v_h) \quad \forall v_h \in X_h$ from $a(u_h, v_h) = \ell(v_h) \quad \forall v_h \in X_h$ \square

Def: If a is symmetric, then $\langle u, v \rangle_a = a(u, v)$ is an inner product on X , the energy inner product. $\|u\|_a = \sqrt{\langle u, u \rangle_a}$ is called energy norm.

Galerkin-orthogonality says that the approximation error $u_h - u$ is orthogonal to X_h wrt. $\langle \cdot, \cdot \rangle_a$.

Cor.: (Céa-lemma) $\|u_h - u\|_X \leq \frac{C}{c} \min_{v_h \in X_h} \|v_h - u\|_X$ „almost-best approximation“

(best approximation) If a is symmetric, then $\|u_h - u\|_a \leq \min_{v_h \in X_h} \|v_h - u\|_a$

Pf: $c \|u_h - u\|_X^2 \leq a(u_h - u, u_h - u) = a(u_h - u, v_h - u) \leq C \|u_h - u\|_X \|v_h - u\|_X \quad \forall v_h \in X_h$

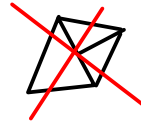
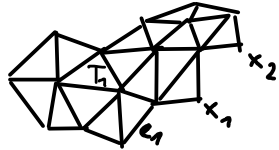
\cdot a symmetric $\Rightarrow \|u_h - u\|_a^2 = a(u_h - u, u_h - u) = a(u_h - u, v_h - u) \leq \|u_h - u\|_a \|v_h - u\|_a \quad \forall v_h \in X_h$ \square
↑ Galerkin-orth.
↑ Cauchy-Schwarz

u_h is (almost) the best possible approximation to u . It only remains to estimate $\min_{v_h \in X_h} \|v_h - u\|_X$.

Finite Elements: elliptic PDEs (2D Poisson equation)

Linear FE: ansatz

Def: Let $\Omega \subset \mathbb{R}^2$. A triangulation of Ω is a triple (V, E, T) with vertex set $V = (x_1, \dots, x_N) \in (\mathbb{R}^2)^N$, edge set $E = (e_1, \dots, e_L)$, triangle set $T = (T_1, \dots, T_M)$, where every edge / every triangle is the convex hull of two / three vertices, every intersection of two triangles is either a common side / corner or empty, $\bar{\Omega} = \bigcup_{i=1}^M T_i$, each edge is a triangle side, and $T_i \neq \emptyset \forall i$.

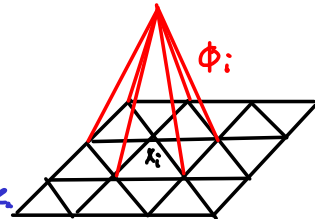


$$h_T = \text{diam } T, \quad h = \max_{T \in T} h_T = \text{grid width}$$

If one chooses $X_h \subset H^1(\Omega)$ as the set of piecewise affine functions, $X_h = \{v : C(\Omega) \mid v|_T \in P_1 \forall T \in T\}$ ($P_1 = \{\text{polynomials of degree} \leq 1\}$), a basis is given by the "FE hat functions" $\phi_1, \dots, \phi_N \in X_h$, which are defined by $\phi_i(x_j) = \begin{cases} 1 & \text{if } i=j \\ 0 & \text{else} \end{cases}$.

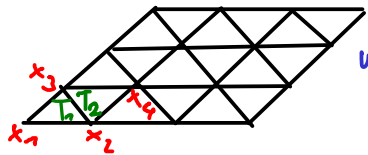
The elements of the chosen basis are also called FE form functions.

Since $\text{supp } \phi_i \cap \text{supp } \phi_j = \emptyset$ for most FE form functions, $A_{ij} = a(\phi_i, \phi_j)$ is sparse.



Finite Elements: elliptic PDEs (2D Poisson equation)

Linear FE: matrix assembly



vertex array =
$$\begin{bmatrix} -x_1 \\ -x_2 \\ \vdots \\ -x_N \end{bmatrix} = \begin{bmatrix} 0 & 0 \\ 1 & 0 \\ 0.5 & 0.5 \\ 1.5 & 0.5 \\ \vdots \end{bmatrix}$$

connectivity array =
$$\begin{bmatrix} 1 & 2 & 3 \\ 3 & 2 & 4 \\ \vdots & \vdots & \vdots \\ 7 & 8 & 9 \\ \vdots & \vdots & \vdots \end{bmatrix} \begin{matrix} \leftarrow T_1 \\ \\ \\ \leftarrow T_i \\ \end{matrix}$$

The computation of matrix $A = (\alpha(\phi_i, \phi_j))_{ij}$ is called assembly.

Since a is of the form $a(u, v) = \int_{\Omega} Q(u, v) dx = \sum_{k=1}^M \int_{T_k} Q(u, v) dx$ for a bilinear operator Q , one iterates over the triangles.
 here without 2Ω -contribution; this works analogously...

Let $T_e = \text{co}(x_{i_1}, x_{i_2}, x_{i_3})$, then only $\phi_{i_1}, \phi_{i_2}, \phi_{i_3}$ are nonzero on T_e .

$\Rightarrow \int_{T_e} Q(\phi_m, \phi_n) dx = 0$ for $m \notin \{i_1, i_2, i_3\}$ or $n \notin \{i_1, i_2, i_3\}$

\Rightarrow one computes the local system matrix $A^e = \left(\int_{T_e} Q(\phi_{i_m}, \phi_{i_n}) dx \right)_{m, n=1, 2, 3}$ and adds it into A at position (i_m, i_n) , $m, n=1, 2, 3$, i.e. one computes

$$A = \sum_{e=1}^M L_e^T A^e L_e \text{ for } L_e = \begin{pmatrix} 0 & \dots & 1 & 0 & \dots & 0 & 0 & \dots & 0 \\ 0 & \dots & 0 & 0 & \dots & 1 & 0 & \dots & 0 \\ 0 & \dots & 0 & 0 & \dots & 0 & 1 & \dots & 0 \end{pmatrix}$$

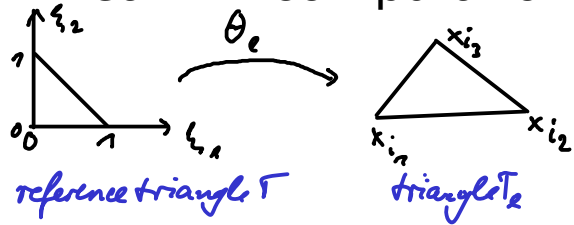
$\uparrow \quad \quad \uparrow \quad \quad \uparrow$
 $i_1 \quad \quad i_2 \quad \quad i_3$

Ex:
$$A = \begin{bmatrix} \dots & \dots & \dots & \dots & \dots \\ \dots & *K & \dots & *K & \dots & *K & \dots \\ \dots & *K & \dots & *K & \dots & *K & \dots \\ \dots & *K & \dots & *K & \dots & *K & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots \end{bmatrix}$$

$$A^e = \begin{pmatrix} x & x & x \\ x & x & x \\ x & x & x \end{pmatrix}$$

Finite Elements: elliptic PDEs (2D Poisson equation)

Linear FE: computation of local system matrix



$$\Theta_e(\xi) = x_{i_1} + \overbrace{(x_{i_2} - x_{i_1} \mid x_{i_3} - x_{i_1})}^{D\Theta_e} \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix} \text{ maps } T \text{ onto } T_e.$$

$$\text{triangle area} = |T_e| = |(x_{i_2} - x_{i_1}) \times (x_{i_3} - x_{i_1})| / 2 = \det D\Theta_e / 2$$

Set $\Psi_1(\xi) = 1 - \xi_1 - \xi_2$, $\Psi_2(\xi) = \xi_1$, $\Psi_3(\xi) = \xi_2 \Rightarrow$ on T_e one has $\phi_{i_m} \circ \Theta_e = \Psi_m / \nabla \phi_{i_m} \circ \Theta_e = D\Theta_e^{-T} \nabla \Psi_m$

Now A^e can be computed using the Ψ_i and Θ_e !

Ex: $\cdot a(u, v) = \int_{\Omega} u v dx \Rightarrow A_{mn}^e = \int_{T_e} \phi_{i_m} \phi_{i_n} dx = \int_T \Psi_m \Psi_n |\det D\Theta_e| d\xi = 2 |T_e| \int_T \Psi_m \Psi_n d\xi$

$$\Rightarrow A^e = \frac{|T_e|}{12} \begin{pmatrix} 2 & 1 & 1 \\ 1 & 2 & 1 \\ 1 & 1 & 2 \end{pmatrix}$$

$\cdot a(u, v) = \int_{\Omega} \nabla u \cdot \nabla v dx \Rightarrow A_{mn}^e = \int_{T_e} \nabla \phi_{i_m} \cdot \nabla \phi_{i_n} dx = \int_T \nabla \Psi_m^T D\Theta_e^{-T} D\Theta_e^{-T} \nabla \Psi_n |\det D\Theta_e| d\xi$

$$= |T_e| \nabla \Psi_m^T D\Theta_e^{-T} D\Theta_e^{-T} \nabla \Psi_n$$

$$D\Theta_e^{-T} = \frac{1}{2|T_e|} \begin{pmatrix} (x_{i_3} - x_{i_1})_2 & -(x_{i_3} - x_{i_1})_1 \\ -(x_{i_2} - x_{i_1})_2 & (x_{i_2} - x_{i_1})_1 \end{pmatrix} \Rightarrow D\Theta_e^{-T} D\Theta_e^{-T} = \frac{1}{4|T_e|^2} \underbrace{\begin{pmatrix} |x_{i_3} - x_{i_1}|^2 & -(x_{i_3} - x_{i_1}) \cdot (x_{i_2} - x_{i_1}) \\ -(x_{i_3} - x_{i_1}) \cdot (x_{i_2} - x_{i_1}) & |x_{i_2} - x_{i_1}|^2 \end{pmatrix}}_R$$

$$\Rightarrow A^e = \frac{1}{4|T_e|} \begin{pmatrix} -1 & -1 \\ 1 & 0 \\ 0 & 1 \end{pmatrix} R \begin{pmatrix} -1 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} = \frac{1}{4|T_e|} \begin{pmatrix} -x_{i_2} - x_{i_3} \\ -x_{i_2} - x_{i_1} \\ -x_{i_1} - x_{i_2} \end{pmatrix} \begin{pmatrix} x_{i_2} - x_{i_3} & | & x_{i_3} - x_{i_1} & | & x_{i_1} - x_{i_2} \end{pmatrix}$$

Finite Elements: elliptic PDEs (2D Poisson equation)

Linear FE: mass & stiffness matrix for regular grid

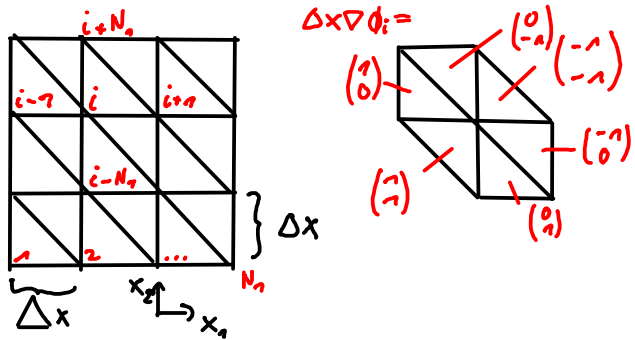
Def: Let Φ_1, \dots, Φ_N be a basis of the FE-space X_h .

mass matrix $M_h = \left(\int_{\Omega} \Phi_i \Phi_j dx \right)_{ij}$, stiffness matrix $L_h = \left(\int_{\Omega} \nabla \Phi_i \cdot \nabla \Phi_j dx \right)_{ij}$

Assembly of $B_i = \ell(\Phi_i)$ also works via iteration over all triangles (homework);

if $\ell(v) = \int_{\Omega} f v dx$ for $f \in X_h$ one can also calculate $B = M_h \begin{pmatrix} f(x_1) \\ \vdots \\ f(x_N) \end{pmatrix}$.

On a regular grid we can often calculate the system matrix by hand:



Ex: $\int_{\Omega} |\nabla \Phi_i|^2 dx = \frac{1}{2} \left[\begin{pmatrix} 0 \\ -1 \end{pmatrix}^2 + \begin{pmatrix} 1 \\ 0 \end{pmatrix}^2 + \begin{pmatrix} 1 \\ 1 \end{pmatrix}^2 + \begin{pmatrix} 0 \\ 1 \end{pmatrix}^2 + \begin{pmatrix} -1 \\ 0 \end{pmatrix}^2 + \begin{pmatrix} -1 \\ -1 \end{pmatrix}^2 \right] = 4$

$\int_{\Omega} \nabla \Phi_i \cdot \nabla \Phi_{i-1} dx = \frac{1}{2} \left[\begin{pmatrix} -1 \\ -1 \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \begin{pmatrix} -1 \\ 0 \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 1 \end{pmatrix} \right] = -1$

$\int_{\Omega} \nabla \Phi_i \cdot \nabla \Phi_{i-1} dx = \int_{\Omega} \nabla \Phi_i \cdot \nabla \Phi_{i+N_x} dx = \int_{\Omega} \nabla \Phi_i \cdot \nabla \Phi_{i-N_x} dx = -1$

$\Rightarrow L_h = \begin{pmatrix} \dots & -1 & \dots & -4 & \dots & -1 & \dots \\ \dots & & & & & & \dots \\ \dots & & & & & & \dots \end{pmatrix}$

$\Rightarrow L_h \begin{pmatrix} u_1 \\ \vdots \\ u_N \end{pmatrix} = -u_{i-N_x} - u_{i+N_x} + 4u_i - u_{i-1} - u_{i+1}$

finite differences for Δ : $\frac{-u_{i-N_x} - u_{i+N_x} + 4u_i - u_{i-1} - u_{i+1}}{\Delta x^2}$

Ex: $\int_{\Omega} \Phi_i^2 dx = \frac{\Delta x^2}{2}$, $\int_{\Omega} \Phi_i \Phi_{i \pm 1} dx = \int_{\Omega} \Phi_i \Phi_{i \pm N_x} dx = \int_{\Omega} \Phi_i \Phi_{i \pm (N_x-1)} dx = \frac{\Delta x^2}{12}$ (homework) $\Rightarrow M_h = \frac{\Delta x^2}{12} \begin{pmatrix} \dots & 1 & \dots & 6 & \dots & 1 & \dots \\ \dots & & & & & & \dots \end{pmatrix}$

Finite Elements: elliptic PDEs (2D Poisson equation)

Linear FE: interpolation error (on reference triangle)

Def: $I_h u = \sum_{i=1}^N u(x_i) \Phi_i \in X_h$ is called FE-interpolant of u

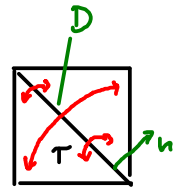
Thm: There exists $C > 0$ such that for all $u \in H^2(\tau)$ we have $\|u - I_h u\|_{L^2} \leq C |u|_{H^2} = C \|D^2 u\|_{L^2}$,
 $\|\nabla u - \nabla I_h u\|_{L^2} \leq C |u|_{H^2}$.
reference triangle $\|A\|_{L^2}^2 = \int \sum_{i,j} |A_{ij}|^2 dx$

Pf: Let $v \in H^1([0,1]^2)$ with $\int_0^1 v(s,0) ds = 0$.

$$\begin{aligned} \|v\|_{L^2}^2 &= \int_0^1 \int_0^1 (v(x_1, x_2) - \int_0^1 v(s,0) ds)^2 dx_1 dx_2 = \int_0^1 \int_0^1 \left| \int_0^1 v(x_1, x_2) - v(s, x_2) ds + \int_0^1 v(s, x_2) - v(s, 0) ds \right|^2 dx_1 dx_2 \\ &= \int_0^1 \int_0^1 \left| \int_0^1 \int_0^1 \frac{\partial v}{\partial x_1}(t, x_2) dt ds + \int_0^1 \int_0^1 \frac{\partial v}{\partial x_2}(s, t) dt ds \right|^2 dx_1 dx_2 \\ &\leq \int_0^1 \int_0^1 \left(\int_0^1 \int_0^1 \left| \frac{\partial v}{\partial x_1}(t, x_2) \right| dt ds + \int_0^1 \int_0^1 \left| \frac{\partial v}{\partial x_2}(s, t) \right| dt ds \right)^2 dx_1 dx_2 \\ &\leq 2 \int_0^1 \int_0^1 \left(\int_0^1 \int_0^1 \left| \frac{\partial v}{\partial x_1}(t, x_2) \right| dt ds \right)^2 + \left(\int_0^1 \int_0^1 \left| \frac{\partial v}{\partial x_2}(s, t) \right| dt ds \right)^2 dx_1 dx_2 \\ &\leq 2 \int_0^1 \int_0^1 \int_0^1 \int_0^1 \left| \frac{\partial v}{\partial x_1}(t, x_2) \right|^2 dt ds dx_1 dx_2 + 2 \int_0^1 \int_0^1 \int_0^1 \int_0^1 \left| \frac{\partial v}{\partial x_2}(s, t) \right|^2 ds dt dx_1 dx_2 \\ &= 2 \left\| \frac{\partial v}{\partial x_1} \right\|_{L^2}^2 + 2 \left\| \frac{\partial v}{\partial x_2} \right\|_{L^2}^2 = 2 \|\nabla v\|_{L^2}^2 \end{aligned}$$

Analogous for v with $\int_0^1 v(0, s) ds = 0$.

Linear FE: interpolation error - cont.



Let $v \in H^1(T)$, $\hat{v}: [0,1]^2 \rightarrow \mathbb{R}$, $\hat{v} = v$ on T , $\hat{v}(x_1, x_2) = v(1-x_2, 1-x_1)$ on $[0,1]^2 \setminus T$
 then $\hat{v} \in H^1([0,1]^2)$ with $\|\hat{v}\|_{L^2}^2 = 2\|v\|_{L^2}^2$ (clear), $\|\nabla \hat{v}\|_{L^2}^2 = 2\|\nabla v\|_{L^2}^2$ (since $\nabla \hat{v} = \nabla v$ on T and

$\nabla \hat{v}(x_1, x_2) = \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix} \nabla v(1-x_2, 1-x_1)$ on $[0,1]^2 \setminus T$: Let $\varphi \in C_0^\infty([0,1]^2; \mathbb{R}^2)$, then

$$\int_0^1 \int_0^1 \hat{v} \operatorname{div} \varphi \, dx = \int_T v \operatorname{div} \varphi \, dx + \int_{T^c} \hat{v} \operatorname{div} \varphi \, dx = \int_D v \varphi \cdot n \, dx - \int_T \nabla v \cdot \varphi \, dx + \int_D \hat{v} \varphi \cdot (-n) \, dx - \int_{T^c} \nabla \hat{v} \cdot \varphi \, dx = - \int_0^1 \int_0^1 \nabla \hat{v} \cdot \varphi \, dx$$

Let $v \in H^1(T)$ with $\int_0^1 v(s, 0) \, ds = 0$ or $\int_0^1 v(0, s) \, ds = 0$, then $\|v\|_{L^2}^2 = \frac{1}{2} \|\hat{v}\|_{L^2}^2 \leq \|\nabla \hat{v}\|_{L^2}^2 = 2\|\nabla v\|_{L^2}^2$
 \Rightarrow with $v = \nabla u - \nabla I_h u$ we have $\|\nabla u - \nabla I_h u\|_{L^2}^2 \leq 2\|\mathcal{D}^2 u - \mathcal{D}^2 I_h u\|_{L^2}^2 = 2\|\mathcal{D}^2 u\|_{L^2}^2$

Let $v \in H^2(T)$ with $v(0,0) = v(1,0) = 0$.

$$\begin{aligned} \|v\|_{L^2}^2 &= \int_0^1 \|v(x_n, \cdot)\|_{L^2}^2 \, dx_n \leq 2 \int_0^1 v(x_n, 0)^2 + \|v(x_n, \cdot) - v(x_n, 0)\|_{L^2}^2 \, dx_n \stackrel{\text{Poincaré}}{\leq} 2\|v(\cdot, 0)\|_{L^2}^2 + 2 \int_0^1 \left\| \frac{\partial v}{\partial x_2}(x_n, \cdot) \right\|_{L^2}^2 \, dx_n \\ &\leq 2 \left\| \frac{\partial v}{\partial x_n}(\cdot, 0) \right\|_{L^2}^2 + 2 \left\| \frac{\partial v}{\partial x_2} \right\|_{L^2}^2 \leq 2C \left\| \frac{\partial v}{\partial x_n} \right\|_{H^1}^2 + 2\|\nabla v\|_{L^2}^2 \leq 2(C+1)\|\nabla v\|_{L^2}^2 + 2C\|\mathcal{D}^2 v\|_{L^2}^2 \\ &\stackrel{\text{Poincaré from before}}{\leq} (6C+4)\|\mathcal{D}^2 v\|_{L^2}^2 \end{aligned}$$

\Rightarrow with $v = u - I_h u$ we have $\|u - I_h u\|_{L^2}^2 \leq (6C+4)\|\mathcal{D}^2 u - \mathcal{D}^2 I_h u\|_{L^2}^2 = (6C+4)\|\mathcal{D}^2 u\|_{L^2}^2 \square$

Finite Elements: elliptic PDEs (2D Poisson equation)

Linear FE: quasiuniform triangulation

Def: A triangulation $(V, \mathcal{E}, \mathcal{T})$ is called $\frac{1}{\rho}$ -quasiuniform, if for all $T_e \in \mathcal{T}$ the affine transformation $\theta_e: T \rightarrow T_e$ of the reference triangle T onto T_e satisfies $\text{cond } D\theta_e \leq \rho$.

Condition number e.g. wrt. 2-norm

Here's a geometric interpretation:

Thm: $r_T \geq 2h_T/\rho \quad \forall T \in \mathcal{T} \iff \frac{1}{\rho}$ -quasiuniform $\iff r_T \geq h_T/8\rho \quad \forall T \in \mathcal{T}$



Pf: Let θ be the reference transformation. Incircle $r_T = \frac{2|T|}{a+b+c}$, $2|T| = |\det D\theta| = b_1 b_2$, $\|D\theta\|_2 = b_2 \begin{cases} \geq h_T/2 \\ \leq 2h_T \end{cases}$
 $b_1 \leq a+b+c \leq b_1 + b_1 + \sqrt{2} b_1 \leq 4b_1$ 1) $\frac{\|D\theta\|_2}{\rho} \leq \frac{2h_T}{\rho} \leq r_T \leq b_2 = \|D\theta^{-1}\|_2^{-1}$ 2) $r_T \geq \frac{b_2}{4} \geq \frac{b_1}{4\rho} \geq \frac{h_T}{8\rho}$ \square

Thm: Let T_e be a triangle of a $\frac{1}{\rho}$ -quasiuniform triangulation, $e \subset T_e$ an edge. There exists $C \equiv C(\rho)$ such that

- trace theorem: $\|v\|_{L^2(e)} \leq C(h_{T_e}^{-\frac{1}{2}} \|v\|_{L^2(T_e)} + h_{T_e}^{\frac{1}{2}} \|\nabla v\|_{L^2(T_e)}) \quad \forall v \in H^1(T_e)$
- Poincaré-inequality: $\|v\|_{L^2(T_e)} \leq C h_{T_e} \|\nabla v\|_{L^2(T_e)} \quad \forall v \in H_0^1(T_e)$
- Poincaré-inequality II: $\|v - \bar{v}\|_{L^2(T_e)} \leq C h_{T_e} \|\nabla v\|_{L^2(T_e)}$, $\bar{v} = \frac{1}{|T_e|} \int_{T_e} v dx \quad \forall v \in H^1(T_e)$
- interp. error: $\|v - I_h v\|_{L^2(T_e)} \leq C h_{T_e}^2 |v|_{H^2(T_e)}$, $\|\nabla v - \nabla I_h v\|_{L^2(T_e)} \leq C h_{T_e} |v|_{H^2(T_e)} \quad \forall v \in H^2(T_e)$

Pf homework (show & use $h_{T_e}/\sqrt{2} \leq \|D\theta_e\|_2 \leq \sqrt{2} h_{T_e}$, $\|D\theta_e^{-1}\|_2 \leq \rho / \|D\theta_e\|_2 \leq \sqrt{2} \rho / h_{T_e}$, $|T_e| = \frac{|\det D\theta_e|}{2} = \frac{\|D\theta_e\|_2}{2\|D\theta_e^{-1}\|_2}$)

Finite Elements: elliptic PDEs (2D Poisson equation)

Linear FE: a priori error estimates

$a: X \times X \rightarrow \mathbb{R}$ with $a(v,v) \geq c \|v\|_X^2$ & $|a(u,v)| \leq C \|u\|_X \|v\|_X$, $\ell: X \rightarrow \mathbb{R}$ with $|\ell(v)| \leq \bar{C} \|v\|_X$

$X_h \subset X = H_0^1(\Omega)$, $a(u,v) = \ell(v) \quad \forall v \in X$, $a(u_h, v_h) = \ell(v_h) \quad \forall v_h \in X_h$

We already know $\|u - u_h\|_{H^1}^2 \leq C \min_{v_h \in X_h} \|u - v_h\|_{H^1}^2$ — how big is this?

Thm: On a \hat{p} -quasiuniform triangulation of grid with h there exists $\hat{C} > 0$, depending only on \hat{p} ,

with $\|u - I_h u\|_{L^2} \leq \hat{C} h^2 |u|_{H^2}$, $\|\nabla u - \nabla I_h u\|_{L^2} \leq \hat{C} h |u|_{H^2}$.

Pf: $\|u - I_h u\|_{L^2}^2 = \sum_{\tau \in \mathcal{T}_h} \|u - I_h u\|_{L^2(\tau)}^2 \leq \sum_{\tau \in \mathcal{T}_h} \hat{C} h^4 |u|_{H^2(\tau)}^2 = \hat{C} h^4 |u|_{H^2}^2$

• analogous for $\|\nabla u - \nabla I_h u\|_{L^2}^2$ □

Cor: $\|u - u_h\|_{H^1}^2 \leq C \|u - I_h u\|_{H^1}^2 \leq C \hat{C}^2 h^2 (1+h^2) |u|_{H^2}^2 \leq \tilde{C} h^2 |u|_{H^2}^2$

Finite Elements: elliptic PDEs (2D Poisson equation)

Linear FE: Aubin-Nitsche duality trick

$a: X \times X \rightarrow \mathbb{R}$ coercive, bounded, bilinear, $l: X \rightarrow \mathbb{R}$ bounded, linear, $X_h \subset X$ subspace H_T^1

$u \in X: a(u, v) = l(v) \quad \forall v \in X$, $u_h \in X_h: a(u_h, v_h) = l(v_h) \quad \forall v_h \in X_h$

We know $\|u_h - u\|_{H^1} \leq Ch \|u\|_{H^2} \Rightarrow \|u_h - u\|_{L^2} \leq Ch \|u\|_{H^1}$; however, one often actually has $\|u_h - u\|_{L^2} \leq Ch^2 \|u\|_{H^2}$

Thm: (Aubin-Nitsche-Lemma) Let H be a Hilbert space with continuous embedding $X \hookrightarrow H$, $w \in X$

with $a(v, w) = (u_h - u, v)_H \quad \forall v \in X$. There exists $C > 0$ depending on a with

adjoint bilinear form $\|u_h - u\|_H^2 \leq C \|u_h - u\|_X \inf_{w_h \in X_h} \|w - w_h\|_X$

Pf: $\|u_h - u\|_H^2 = (u_h - u, u_h - u)_H = a(u_h - u, w - w_h) = a(u_h - u, w - w_h) \leq C \|u_h - u\|_X \|w - w_h\|_X \quad \forall w_h \in X_h$
 \ is this well-defined?

Cor: Let $X = H_T^1(\Omega)$. Assume for a that the solution w_r of $a(v, w_r) = (r, v)_{L^2} \quad \forall v \in H_T^1$ satisfies

$|w_r|_{H^2} \leq C \|r\|_{L^2} \quad \forall r \in L^2(\Omega)$. Then $\|u_h - u\|_{L^2} \leq \tilde{C} h^2 \|u\|_{H^2}$

Pf: Take $H = L^2(\Omega)$, then $\|u_h - u\|_{L^2}^2 \leq C \|u_h - u\|_{H^1} \inf_{w_h \in X_h} \|w - w_h\|_{H^1} \leq \hat{C} h \|u\|_{H^2} h \|w\|_{H^2} \leq \hat{C} h^2 \|u\|_{H^2} \|u_h - u\|_{L^2}$

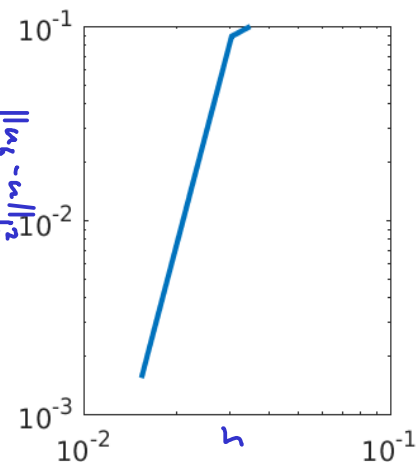
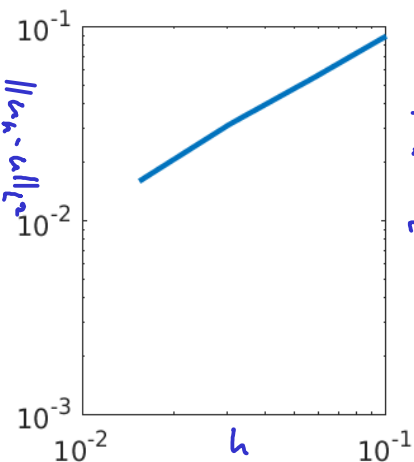
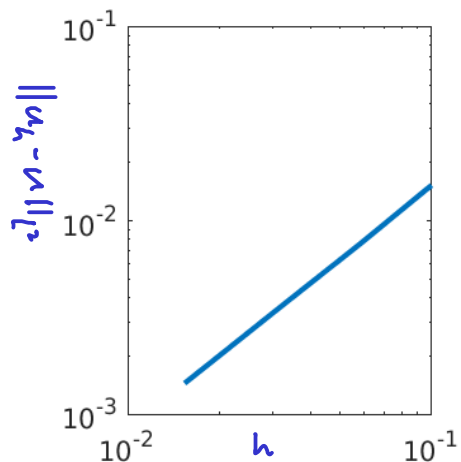
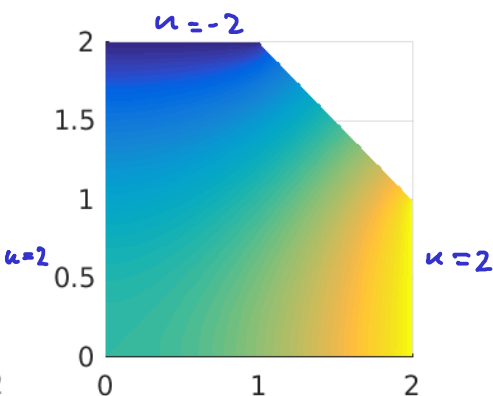
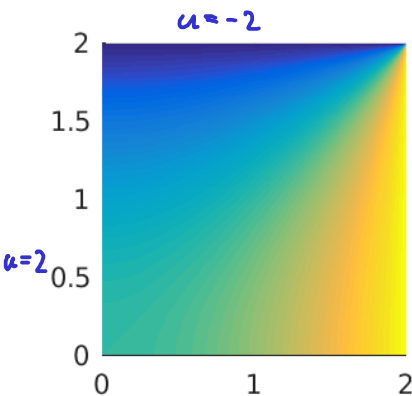
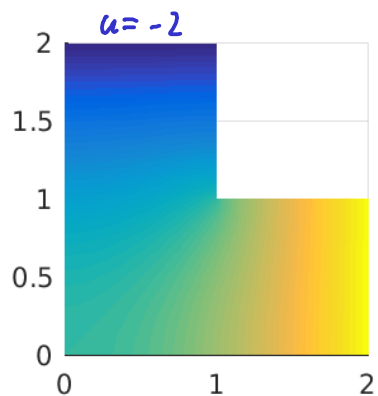
Note: The estimates require that $u \in H^2$ & $w_r \in H^2$ with $|w_r|_{H^2} \leq C \|r\|_{L^2}$.

There are situations in which $u, w_r \notin H^2$, e.g. for $\Omega = \square \Rightarrow \|u_h - u\| = O(h^{\frac{1}{2}})$ smaller

Finite Elements: elliptic PDEs (2D Poisson equation)

Linear FE: example simulation

$$-\Delta u = 0 \text{ on } \Omega$$



Finite Elements: elliptic PDEs (2D Poisson equation)

Linear FE: regularity (Friedrichs' theorem)

We showed

$$\|u_h - u\|_{H^1} \leq Ch |u|_{H^2} \quad \& \quad \|u_h - u\|_{L^2} \leq Ch^2 |u|_{H^2} \quad \text{if} \quad |w_r|_{H^2} \leq C \|r\|_{L^2}$$

Do $u, w_r \in H^2$ and $|w_r|_{H^2} \leq C \|r\|_{L^2}$ hold, and how big is $|u|_{H^2}$?

Ex: $\Omega = [0,1]^2$, $a(u,v) = \int_{\Omega} \nabla u \cdot \nabla v + uv \, dx$, $\ell(v) = \int_{\Omega} f v \, dx$,

$$\left. \begin{array}{l} \{ u \in X = H_0^1(\Omega), a(u,v) = \ell(v) \forall v \in X \\ \{ w_r \in X, a(v, w_r) = (r, v)_{L^2} \forall v \in X \} \end{array} \right\} \Rightarrow \left\{ \begin{array}{l} |u|_{H^2} \leq C \|f\|_{L^2} \\ |w_r|_{H^2} \leq C \|r\|_{L^2} \end{array} \right.$$

Pf: (proof for u , analogous for w_r) $\|u\|_{L^2}^2 \leq \|u\|_{H^1}^2 \leq a(u,u) = \ell(u) \leq \|f\|_{L^2} \|u\|_{L^2}$

$$\begin{aligned} |u|_{H^2}^2 &= \int_{\Omega} \left| \frac{\partial^2 u}{\partial x_1^2} \right|^2 + 2 \left| \frac{\partial^2 u}{\partial x_1 \partial x_2} \right|^2 + \left| \frac{\partial^2 u}{\partial x_2^2} \right|^2 dx = \int_{\Omega} \left| \frac{\partial^2 u}{\partial x_1^2} \right|^2 - 2 \frac{\partial}{\partial x_2} \frac{\partial^2 u}{\partial x_1 \partial x_2} \frac{\partial u}{\partial x_1} + \left| \frac{\partial^2 u}{\partial x_2^2} \right|^2 dx \\ &= \int_{\Omega} \left| \frac{\partial^2 u}{\partial x_1^2} \right|^2 + 2 \frac{\partial u}{\partial x_2} \frac{\partial^2 u}{\partial x_1^2} + \left| \frac{\partial^2 u}{\partial x_2^2} \right|^2 dx \\ &= \|\Delta u\|_{L^2}^2 = \|f - u\|_{L^2}^2 \leq 2 \|f\|_{L^2}^2 + 2 \|u\|_{L^2}^2 \leq 2(1+1) \|f\|_{L^2}^2 \quad \square \end{aligned}$$

Thm: (e.g. Gilbarg & Trudinger Thm. 8.12) Let Ω bounded, $\partial\Omega$ C^2 -regular,

$$a(u,v) = \int_{\Omega} \nabla u^T A \nabla v + u b \cdot \nabla v - c \cdot \nabla u v - d u v \, dx, \quad \ell(v) = \int_{\Omega} f v \, dx,$$

A, b Lipschitz continuous, $c, d \in L^\infty(\Omega)$, $f \in L^2(\Omega)$, a coercive, $h \in H^2(\Omega)$.

The weak solution $u \in H^1(\Omega)$ of $a(u,v) = \ell(v) \forall v \in H^1(\Omega)$, $u = h$ on $\partial\Omega$,

$$\text{satisfies} \quad \|u\|_{H^2} \leq C (\|u\|_{L^2} + \|f\|_{L^2} + \|h\|_{H^2}) \quad (C \text{ depending on } \Omega, A, b, c, d)$$

Similar statements hold for piecewise smooth boundary such that no two neighbouring pieces have Neumann boundary conditions (see e.g. Bacuta, Mazzucato, Nistor, Zikatanov: Interface and mixed boundary value problems on n-dimensional polyhedral domains)

Finite Elements: elliptic PDEs (2D Poisson equation)

FEM: general FE spaces

Def: A finite element is a triple $(K, \mathcal{P}, \mathcal{N}) =$

(element domain, space of form functions, set of nodal variables) with

- $K \subset \mathbb{R}^n$ open, bounded, simply connected with piecewise smooth boundary
- \mathcal{P} a k -dimensional space of functions $f: K \rightarrow \mathbb{R}$
- $\mathcal{N} = \{N_1, \dots, N_k\}$ a set of linearly independent linear functionals on \mathcal{P}

Def: Let $(K, \mathcal{P}, \mathcal{N})$ be a finite element, and let $\{\psi_1, \dots, \psi_k\}$ be a basis of \mathcal{P} dual to \mathcal{N} , i.e. $N_i(\psi_j) = \delta_{ij}$. Such a basis is called nodal basis for \mathcal{P} .

Thm: $\psi, \tilde{\psi} \in \mathcal{P}$ with $N(\psi) = N(\tilde{\psi}) \forall N \in \mathcal{N} \Rightarrow \psi = \tilde{\psi}$

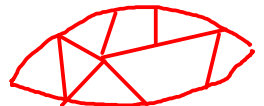
Pf: $N(\psi - \tilde{\psi}) = N(\psi) - N(\tilde{\psi}) = 0 \forall N \Rightarrow$ it suffices to show that $N(\phi) = 0 \forall N \Rightarrow \phi = 0$.

Let $\phi = \sum_{i=1}^k \alpha_i \psi_i$, then $0 = N_j(\phi) = \sum_{i=1}^k \alpha_i N_j(\psi_i) = \alpha_j \quad \forall j. \quad \square$

Def: A subdivision of a domain $\Omega \subset \mathbb{R}^n$ is a finite set $\{K_1, \dots, K_M\}$ of open sets with

$$(1) K_i \cap K_j = \emptyset \text{ for } i \neq j$$

$$(2) \bar{K}_1 \cup \dots \cup \bar{K}_M = \bar{\Omega}$$



Def: Let $\{K_1, \dots, K_M\}$ be a subdivision of Ω and $(K_i, \mathcal{P}^i, \mathcal{N}^i)$ be a finite element $\forall i = 1, \dots, M$. $X_h = \{v: \Omega \rightarrow \mathbb{R} \mid v|_{K_i} \in \mathcal{P}^i, N(\nu|_{K_i}) = N(\nu) \forall N \in \mathcal{N}^i \forall i = 1, \dots, M\}$ is called finite element space.

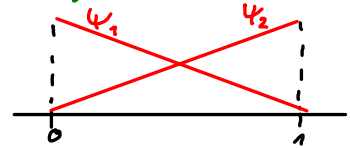
Finite Elements: elliptic PDEs (2D Poisson equation)

FEM: FE examples

Ex: (1D Lagrange element) $K=(0,1)$, $\mathcal{P}=\mathcal{P}_1=\{\text{polynomials of degree 1}\}$, $\mathcal{N}=\{N_1, N_2\}$

$$N_1(v) = v(0), \quad N_2(v) = v(1)$$

$$\text{nodal basis: } \psi_1(x) = 1-x, \quad \psi_2(x) = x$$

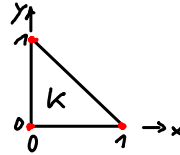


$\dim \mathcal{P} = ?$

Ex: (2D Lagrange element) $K=T:=\{(x,y) \in [0,1]^2 \mid y \leq 1-x\}$, $\mathcal{P}=\mathcal{P}_1$, $\mathcal{N}=\{N_1, N_2, N_3\}$

$$N_1(v) = v(0,0), \quad N_2(v) = v(1,0), \quad N_3(v) = v(0,1)$$

$$\text{nodal basis: } \psi_1(x,y) = 1-x-y, \quad \psi_2(x,y) = x, \quad \psi_3(x,y) = y$$

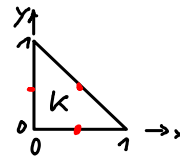


$\dim \mathcal{P} = ?$

Ex: (2D Courser-Raviart element) $K=T$, $\mathcal{P}=\mathcal{P}_1$, $\mathcal{N}=\{N_1, \dots, N_3\}$

$$N_1(v) = v\left(\frac{1}{2}, \frac{1}{2}\right), \quad N_2(v) = v\left(0, \frac{1}{2}\right), \quad N_3(v) = v\left(\frac{1}{2}, 0\right)$$

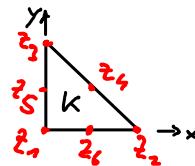
$$\text{nodal basis: } \psi_1(x,y) = 2(x+y)-1, \quad \psi_2(x,y) = 1-2x, \quad \psi_3(x,y) = 1-2y$$



Ex: (2D Lagrange element) $K=T$, $\mathcal{P}=\mathcal{P}_2$, $\mathcal{N}=\{N_1, \dots, N_6\}$; $N_i(v) = v(z_i)$

$$\psi_1(x,y) = (1-2x-2y)(1-x-y), \quad \psi_2(x,y) = -x(1-2x), \quad \psi_3(x,y) = -y(1-2y),$$

$$\psi_4(x,y) = 4xy, \quad \psi_5(x,y) = 4y(1-x-y), \quad \psi_6(x,y) = 4x(1-x-y)$$

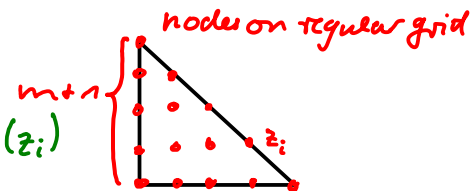


Finite Elements: elliptic PDEs (2D Poisson equation)

FEM: FE examples

Ex: (2D Lagrange element of order m)

$$K = T, \mathcal{P} = \mathcal{P}_m, \mathcal{N} = \{N_1, \dots, N_k\}, k = 1 + 2 + \dots + (m+1), N_i(v) = v(z_i)$$



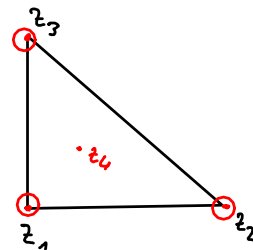
Ex: (2D Hermite element) $K = T, \mathcal{P} = \mathcal{P}_3$ $\dim \mathcal{P} = ?$

$$N_1(v) = v(z_1), N_2(v) = v(z_2), N_3(v) = v(z_3), N_4(v) = v(z_4)$$

$$N_5(v) = \frac{\partial v}{\partial x}(z_1), N_6(v) = \frac{\partial v}{\partial x}(z_2), N_7(v) = \frac{\partial v}{\partial x}(z_3)$$

$$N_8(v) = \frac{\partial v}{\partial y}(z_1), N_9(v) = \frac{\partial v}{\partial y}(z_2), N_{10}(v) = \frac{\partial v}{\partial y}(z_3)$$

nodal basis: determine via polynomial interpolation



Ex: (2D Hermite element of order m)

$$K = T, \mathcal{P} = \mathcal{P}_m, \mathcal{N} = \{N_1, \dots, N_k\}, k = \frac{(m+1)(m+2)}{2}$$

m	1	2	3	4	5	6	7	8
k	3	6	10	15	21	28	36	45

$N_i(v)$ = function values & directional derivatives up to order l in z_1, z_2, z_3

there are $l+1$ independent directional derivatives of order l

$$\Rightarrow 3(1+2+\dots+(l+1)) = 3 \frac{(l+1)(l+2)}{2} =: L \text{ form functions} \rightarrow \text{pick } l = \left\lfloor \frac{-3 + \sqrt{1+8k/3}}{2} \right\rfloor$$

& complement with $k-L$ form functions in z_4

Finite Elements: elliptic PDEs (2D Poisson equation)

FEM: interpolant

Def.: Let $(K, \mathcal{P}, \mathcal{N})$ be a finite element with nodal basis $\{\psi_1, \dots, \psi_n\}$. The local interpolant of a function $v: K \rightarrow \mathbb{R}$ is $I_K v = \sum_{i=1}^n N_i(v) \psi_i$.

Let $\{K_1, \dots, K_M\}$ be a subdivision of Ω and $(K_i, \mathcal{P}_i, \mathcal{N}_i)$ a finite element $\forall i=1, \dots, M$.

The global interpolant $I_h v$ of $v: \Omega \rightarrow \mathbb{R}$ is defined by $I_h v|_{K_i} = I_{K_i} v \quad \forall i$.

The associated finite element space can be represented as $X_h = \{I_h v \mid v \in C^\infty(\bar{\Omega})\}$.

Def.: If $I_h v \in C^r(\Omega) \quad \forall v \in C^\infty(\bar{\Omega})$, then X_h is called a C^r finite element space.

Thm.: For a triangulation with Lagrange or Hermite elements, the associated FE space is a C^0 FE space.

Pf.: It suffices to show that on an edge $e = \bar{K}_i \cap \bar{K}_j$ between two elements we have

$$I_{K_i} v|_e = I_{K_j} v|_e \quad \forall v \in C^\infty(\bar{\Omega}).$$

For a Lagrange element of order m , $I_{K_i} v|_e, I_{K_j} v|_e$ are 1D polynomials of degree m

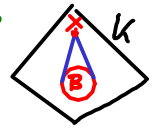
\Rightarrow the conditions $I_{K_i} v(z_r) = v(z_r) = I_{K_j} v(z_r)$ for the $m+1$ points z_r on e imply $I_{K_i} v|_e = I_{K_j} v|_e$.

Hermite element: homework

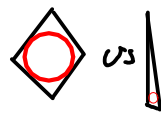
□

FEM: Bramble-Hilbert lemma

Def: $K \subset \mathbb{R}^n$ is called star-shaped wrt. a ball $B \subset K$ if the convex hull of $\{x\} \cup B$ satisfies $\text{co}(\{x\} \cup B) \subset K$ for all $x \in K$.



A subdivision $\{K_1, \dots, K_M\}$ of Ω is called $\frac{1}{\rho}$ -quasiuniform or regular if $\forall i = 1, \dots, M: \frac{\rho}{\rho} \leq \frac{\tau_{K_i}}{h_{K_i}}$, where τ_{K_i} = diameter of the largest ball in K_i .



Stephen, not David

Thm: (Bramble-Hilbert lemma) Let $\{K_1, \dots, K_M\}$ be a $\frac{1}{\rho}$ -quasiuniform subdivision of the bounded, open domain Ω , and let $(K_i, \mathcal{P}^i, \mathcal{N}^i)$ be finite elements with

- (1) K_i is star-shaped wrt. a ball,
- (2) $\mathcal{P}_{m-1} \subset \mathcal{P}^i \subset C^m(\bar{K}_i)$,
- (3) $\mathcal{N}^i \subset$ bounded linear forms on $C^l(\bar{K}_i)$,

this form of the lemma actually is due to Dupont & Scott

then for $m > l + \frac{n}{2}$ and $0 \leq j \leq m$ we have: $|u - I_{K_i} u|_{H^j(K_i)} \leq C(m, n, \rho) h_{K_i}^{m-j} |u|_{H^m(K_i)}$
 $\Rightarrow |u - I_h u|_{H^j(\Omega)} \leq C(m, n, \rho) h^{m-j} |u|_{H^m(\Omega)}$ if $I_h u \in H^j(\Omega)$

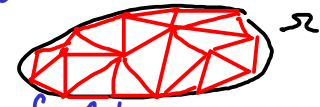
Cor: $\|u_h - u\|_{H^r} \leq C h^{m-r} |u|_{H^m}$ for our elliptic problems & C^0 FE approximations

Finite Elements: elliptic PDEs (2D Poisson equation)

Variational crimes: deviation from exact methods

So far we only considered FE approximations in which u_h satisfies the same weak equation as u (i.e. $a(u, v) = \ell(v) \forall v \in X$), only on a subspace X_h . In practice, however, one often deviates from that, e.g.

• Ω has a curvilinear boundary $\partial\Omega$, but is approximated by triangulations



• the system matrix $A_{mn} = \int_{\Omega} Q(\phi_m, \phi_n) dx$ or right-hand side $B_m = \int_{\Omega} f \phi_m dx$ are only computed approximately, e.g. via numerical quadrature

• the finite element space X_h is not a subspace of the solution space X (e.g. Crouzeix-Raviart-FE-space $\notin H^1(\Omega)$)

Such deviations are called "variational crimes".

Ex: $\ell(v) = \int_{\Omega} f(x) v(x) dx$ is replaced by midpoint quadrature on each triangle

T_e of the triangulation \mathcal{T} , i.e. by

$$\ell_h(v) = \sum_{e=1}^n |T_e| f(\bar{x}_e) v(\bar{x}_e) \quad \text{with } \bar{x}_e \text{ the midpoint of } T_e.$$

Finite Elements: elliptic PDEs (2D Poisson equation)

Variational crimes: Strang's first lemma

$u \in X$ solves $a(u, v) = l(v) \forall v \in X$; a bilinear, coercive bounded; l linear, bounded

$u_h \in X_h$ solves $a_h(u_h, v_h) = l_h(v_h) \forall v_h \in X_h$; a_h bilinear, uniformly coercive; l_h linear

$X_h \subset X$ i.e. $a_h(v_h, v_h) \geq C \|v_h\|_X^2$ indep. of h

By the way: a_h & l_h only need to be defined on X_h ! E.g., $a_h(u_h, v_h)$ could evaluate the gradient of u_h at a quadrature point, which would not be well-defined for an arbitrary function u from $X = H^1(\Omega)$!

Thm: (Strang's first lemma) There exists $K \geq 0$ independent of h such that

$$\|u_h - u\|_X \leq K \left[\inf_{v_h \in X_h} \left(\|u - v_h\|_X + \sup_{w_h \in X_h} \frac{|a(v_h, w_h) - a_h(v_h, w_h)|}{\|w_h\|_X} \right) + \sup_{w_h \in X_h} \frac{|l(w_h) - l_h(w_h)|}{\|w_h\|_X} \right]$$

Pf: Let $v_h \in X_h$ and $w_h = u_h - v_h$.

$$\begin{aligned} c \|u_h - v_h\|_X^2 &\leq a_h(u_h - v_h, w_h) = a(u - v_h, w_h) + [a(v_h, w_h) - a_h(v_h, w_h)] + [a_h(u_h, w_h) - a(u, w_h)] \\ &\leq C \|u - v_h\|_X \|w_h\|_X + [a(v_h, w_h) - a_h(v_h, w_h)] + [l_h(w_h) - l(w_h)] \end{aligned}$$

$$\Rightarrow c \|u_h - v_h\|_X \leq C \|u - v_h\|_X + \frac{|a(v_h, w_h) - a_h(v_h, w_h)|}{\|w_h\|_X} + \frac{|l(w_h) - l_h(w_h)|}{\|w_h\|_X}$$

Now use $\|u_h - u\|_X \leq \|u - v_h\|_X + \|u_h - v_h\|_X$

□

Variational crimes: Strang's first lemma - application

Ex: Solve $-\Delta u + u = f$ on $\Omega = [0,1]^2$, $u=0$ on $\partial\Omega$

$\Rightarrow X = H_0^1(\Omega)$, $a(u,v) = \int_{\Omega} \nabla u \cdot \nabla v + uv \, dx$, $l(v) = \int_{\Omega} f v \, dx$ satisfy LM conditions

$\frac{1}{\rho}$ -quasiuniform triangulation $(\mathcal{T}, \varepsilon, \mathcal{J})$ with grid width h ; $X_h = 1$ -order-Lagrange FE-space

$a_h = a$, $l_h(v_h) = \sum_{T \in \mathcal{T}} |T| \sum_{x \in \mathcal{N} \cap T} \frac{f(x)v(x)}{3}$ "trapezium rule"

$u_h \in X_h$ solves $a_h(u_h, v_h) = l_h(v_h) \quad \forall v_h \in X_h$; how big is $\|u_h - u\|_{H^1}$?

There is $K > 0$ with $\|u_h - u\|_{H^1} \leq K(h \|u\|_{H^2} + h^2 \|f\|_{C^2}) \leq \tilde{K}(h \|f\|_{L^2} + h^2 \|f\|_{C^2})$

Pf: Strang: $\|u_h - u\|_{H^1} \leq K \left[\inf_{v_h \in X_h} \|u - v_h\|_{H^1} + \sup_{w_h \in X_h} \frac{|l(w_h) - l_h(w_h)|}{\|w_h\|_{H^1}} \right]$

Bramble-Hilbert: $\inf_{v_h \in X_h} \|u - v_h\|_{H^1} \leq \|u - I_h u\|_{H^1} \leq C(\rho) h \|u\|_{H^2}$

$$\begin{aligned} |l(w_h) - l_h(w_h)| &= \left| \int_{\Omega} f w_h \, dx - \int_{\Omega} I_h(f w_h) \, dx \right| \leq \sum_{T \in \mathcal{T}} \int_T |f w_h - I_h(f w_h)| \, dx \\ &\leq \sum_{T \in \mathcal{T}} \|f w_h - I_h(f w_h)\|_{L^2(T)} \sqrt{|T|} \leq \sum_{T \in \mathcal{T}} \sqrt{|T|} C(\rho) h^2 |f w_h|_{H^2(T)} \\ &\leq \left(\sum_{T \in \mathcal{T}} |T| \right)^{\frac{1}{2}} C(\rho) h^2 \left(\sum_{T \in \mathcal{T}} |f w_h|_{H^2(T)}^2 \right)^{\frac{1}{2}} = C(\rho) \sqrt{|\Omega|} h^2 \left(\sum_{T \in \mathcal{T}} |f w_h|_{H^2(T)}^2 \right)^{\frac{1}{2}} \end{aligned}$$

$$|f w_h|_{H^2(T)}^2 = \int_T (w_h D^2 f + \nabla f \otimes \nabla v_h + \nabla v_h \otimes \nabla f)^2 \, dx \leq \rho \left(\max_{x \in \Omega} |D^2 f(x)|^2 + \max_{x \in \Omega} |\nabla f(x)|^2 \right) \|w_h\|_{H^1(T)}^2 \quad \square$$

Finite Elements: elliptic PDEs (2D Poisson equation)

Variational crimes: Strang's second lemma

$u \in X$ solves $a(u, v) = \ell(v) \forall v \in X$; a bilinear, coercive, bounded; ℓ linear, bounded

$X_h \not\subset X$, i.e. in particular $\|\cdot\|_X$ is not defined on X_h !

\Rightarrow use (grid-dependent) norm $\|\cdot\|_h$ on $X + X_h$ u_h well-defined if X_h Hilbert-space with norm $\|\cdot\|_h$

$u_h \in X_h$ solves $a_h(u_h, v_h) = \ell_h(v_h) \forall v_h \in X_h$; a_h bilinear; ℓ_h linear

$$a_h(v_h, v_h) \geq c \|v_h\|_h^2 \quad \forall v_h \in X_h; \quad |a_h(u, v_h)| \leq C \|u\|_h \|v_h\|_h \quad \forall u \in X + X_h, v_h \in X_h$$

Thm: (Strang's second lemma; Berger-Scott-Strang-lemma) $\exists K > 0$ independent of h such that

$$\|u_h - u\|_h \leq K \left[\underbrace{\inf_{v_h \in X_h} \|u - v_h\|_h}_{\text{"approximation error"}} + \underbrace{\sup_{w_h \in X_h} \frac{|a_h(u, w_h) - \ell_h(w_h)|}{\|w_h\|_h}}_{\text{"consistency error"}} \right]$$

Pf: Let $v_h \in X_h, w_h = u_h - v_h$. $c \|u_h - u_h\|_h^2 \leq a_h(u_h - v_h, w_h) = a_h(u - v_h, w_h) + [\ell_h(w_h) - a_h(u, w_h)]$

$$\leq C \|u_h - v_h\|_h \|w_h\|_h + |\ell_h(w_h) - a_h(u, w_h)|$$

$$\Rightarrow \|u_h - v_h\|_h \leq \frac{C}{c} \|u - v_h\|_h + \frac{1}{c} \frac{|\ell_h(w_h) - a_h(u, w_h)|}{\|w_h\|_h}; \quad \text{now use } \|u - u_h\|_h \leq \|u - v_h\|_h + \|u_h - v_h\|_h \quad \square$$

Ex: Examples for the norm $\|\cdot\|_h$ are given by

$$\|v\|_{m,h} = \sqrt{\sum_{T \in \mathcal{T}_h} \|v\|_{H^m(T)}^2} \quad \left(\neq \|v\|_{H^m(\Omega)} \right)$$

Finite Elements: elliptic PDEs (2D Poisson equation)

Variational crimes: Strang's second lemma - application

$\Omega = [0, 1]^2$, $X = H_0^1(\Omega)$, $f \in L^2(\Omega)$, $a(u, v) = \int_{\Omega} \nabla u \cdot \nabla v + uv \, dx$, $l(v) = \int_{\Omega} f v \, dx$ for $u, v \in X$

Crouzeix-Raviart-elements on quasiform triangulation \mathcal{T}_h i.e.

$X_h = \{v \in L^2(\Omega) \mid v|_{\tau} \in \mathcal{P}_1 \, \forall \tau \in \mathcal{T}_h \text{ \& } v \text{ is continuous at midpoints of triangle sides \& } v = 0 \text{ at midpoints of triangle sides on } \partial\Omega\}$

$a_h(u, v) = \sum_{\tau \in \mathcal{T}_h} \int_{\tau} \nabla u \cdot \nabla v + uv \, dx$, $l_h(v) = \int_{\Omega} f v \, dx$ for $u, v \in X_h \neq X!$

$u \in X$ solves $a(u, v) = l(v) \, \forall v \in X$, $u_h \in X_h$ solves $a_h(u_h, v_h) = l_h(v_h) \, \forall v_h \in X_h$

we know: $u \in H^2(\Omega)$ with $\|u\|_{H^2} \leq C \|f\|_{L^2}$

Thm: $\|u_h - u\|_{1,h} \leq ch \|f\|_{L^2}$ for a $c > 0$. *int. by parts & $u \in H^2$*

Pf: $\because a_h(u, v_h) - l_h(v_h) = \sum_{\tau_i} \int_{\tau_i} \nabla u \cdot \nabla v_h + u v_h - f v_h \, dx = \sum_{\tau_i} \int_{\partial\tau_i} \nabla u \cdot n v_h \, dx - \int_{\tau_i} v_h (\Delta u - u + f) \, dx$
 $= \sum_{\tau_i} \sum_{\text{edge } e \in \partial\tau_i} \int_e \nabla u(x) \cdot n (v_h(x) - v_h(x_e)) \, dx$ *with $x_e = \text{midpoint of } e$*

since $\nabla u|_{\tau_i} \in H^1(\tau_i)$, for $e = \tau_1 \cap \tau_2$ the trace theorem implies $\nabla u|_{\tau_1} \cdot n_{\tau_1} = -\nabla u|_{\tau_2} \cdot n_{\tau_2}$ on e

$\int_e \omega_h - \omega_h(x_e) \, dx = 0$
 $= \sum_{\tau_i} \sum_{e \in \partial\tau_i} \int_e (\nabla u - \nabla I_h u) \cdot n (v_h - v_h(x_e)) \, dx \leq \sum_{\tau_i} \sum_{e \in \partial\tau_i} \|\nabla(u - I_h u)\|_{L^2(e)} \|v_h - v_h(x_e)\|_{L^2(e)}$

Variational crimes: Strang's second lemma - application (cont.)

$$\begin{aligned}
 \cdot \|\nabla(u - I_h u)\|_{L^2(e)} &\leq \check{C} \left(h^{-\frac{1}{2}} \|\nabla(u - I_h u)\|_{L^2(T_i)} + h^{\frac{1}{2}} \|D^2(u - I_h u)\|_{L^2(T_i)} \right) && \text{trace theorem} \\
 &\leq \bar{C} h^{\frac{1}{2}} \|D^2(u - I_h u)\|_{L^2(T_i)} && \text{Bramble-Hilbert} \\
 &= \bar{C} h^{\frac{1}{2}} |u|_{H^2(T_i)}
 \end{aligned}$$

$$\begin{aligned}
 \cdot \|v_h - \bar{v}_h(x_e)\|_{L^2(e)} &\leq \|v_h - \bar{v}_h(T_i)\|_{L^2(e)} && \bar{v}_h(T_i) = \frac{1}{|T_i|} \int_{T_i} v_h dx \\
 &\leq \check{C} \left(h^{-\frac{1}{2}} \|v_h - \bar{v}_h(T_i)\|_{L^2(T_i)} + h^{\frac{1}{2}} \|\nabla(v_h - \bar{v}_h(T_i))\|_{L^2(T_i)} \right) && \text{trace theorem} \\
 &\leq C h^{\frac{1}{2}} |v_h|_{H^1(T_i)} && \text{Poincaré-ineq.}
 \end{aligned}$$

$$\Rightarrow |a_h(u, v_h) - \bar{a}_h(v_h)| \leq \sum_{T_i} C h |u|_{H^2(T_i)} |v_h|_{H^1(T_i)} \leq C h |u|_{H^2(\Omega)} \|v_h\|_{1,h}$$

Cauchy-Schwarz

$$\cdot \inf_{v_h \in X_h} \|u - v_h\|_{1,h}^2 \leq \|u - I_h u\|_{1,h}^2 = \sum_{T_i} \|u - I_h u\|_{H^1(T_i)}^2 \leq \sum_{T_i} C h^2 |u|_{H^2(T_i)}^2 = C h^2 |u|_{H^2(\Omega)}^2$$

$$\cdot \text{conditions of Strang's second lemma satisfied} \Rightarrow \|u_h - u\|_h \leq \tilde{K} \left(h |u|_{H^2} + h |u|_{H^2} \right) \square$$

Finite Elements: elliptic PDEs (2D Poisson equation)

Variational crimes: duality (extension of Aubin-Nitsche)

- $u \in X$ solves $a(u, v) = \ell(v) \quad \forall v \in X$; a bilinear, coercive, bounded; ℓ linear, bounded
- Aubin-Nitsche-cond.: H Hilbert space with continuous embedding $X \hookrightarrow H$
- Strang's 2. Lemma-cond.: $a_h(v_h, v_h) \geq c \|v_h\|_h^2 \quad \forall v_h \in X_h$; $|a_h(u, v)| \leq C \|u\|_h \|v\|_h \quad \forall u, v \in X + X_h$
 a_h linear & bounded on X_h , $\|\cdot\|_h$ norm on Hilbert space X_h
- $u_h \in X_h$ solves $a_h(u_h, v_h) = \ell_h(v_h) \quad \forall v_h \in X_h$
- additionally: a_h defined on $X + X_h$, $a_h = a$ on X , $X_h \subset H$, $\|\cdot\|_h$ defined on $X + X_h$

Thm: $w_r \in X$ solves $a(v, w_r) = (r, v)_H \quad \forall v \in X$, $w_h \in X_h$ solves $a_h(v_h, w_h) = (r, v_h)_H \quad \forall v_h \in X_h$. is not used in proof →

$$\|u_h - u\|_H \leq \sup_{r \in H} \frac{1}{\|r\|_H} \left(C \|u_h - u\|_h \|w_r - w_h\|_h + [(u_h - u, r)_H - a_h(u_h - u, w_r)] + [\ell_h(w_h) - \ell(w_r) - a_h(u, w_h - w_r)] \right)$$

Pf: Set $r = u_h - u$. $(u_h - u, u_h - u)_H = a_h(u_h - u, w_r - w_h) + \text{rest}$ like Aubin-Nitsche

$$\begin{aligned} \text{rest} &= (u_h - u, r)_H - a_h(u_h - u, w_r - w_h) = [(u_h - u, r)_H - a_h(u_h - u, w_r)] + a_h(u_h - u, w_h) \\ &= [(u_h - u, r)_H - a_h(u_h - u, w_r)] + [\ell_h(w_h) - \ell(w_r) - a_h(u, w_h - w_r)] \end{aligned}$$

$$\Rightarrow \|u_h - u\|_H \leq \frac{1}{\|r\|_H} \left(C \|u_h - u\|_h \|w_r - w_h\|_h + [(u_h - u, r)_H - a_h(u_h - u, w_r)] + [\ell_h(w_h) - \ell(w_r) - a_h(u, w_h - w_r)] \right) \quad \square$$

Variational crimes: duality - application

Solve $-\Delta u = f$ on $[0,1]^2$ with 0-Dirichlet bc with Crouzeix-Raviart elements \Rightarrow setting as before

Thm: $\|u_h - u\|_{L^2(\Omega)} \leq Ch^2 \|f\|_{L^2(\Omega)}$ for a $C > 0$.

Pf: $\because H = L^2(\Omega), \|\cdot\|_H = \|\cdot\|_{L^2}$. We already know $\|u_h - u\|_{1,h} \leq Ch \|f\|_{L^2}, \|w_h - w_r\|_{1,h} \leq Ch \|r\|_{L^2}$.

Furthermore, we have shown two slides ago

$$|(v_h, g)_{L^2} - a_h(v_h, w_g)| = |(v_h, g)_{L^2} - a_h(w_g, v_h)| \leq Ch |w_g|_{H^2} \|v_h\|_{1,h} \quad (*)$$

Note: In previous calculation, test function was $v_h \in X_h$, but also $v_h \in X + X_h$ is possible.

To this end replace $v_h(x_e)$ with $\bar{v}_e = \frac{1}{|e|} \int_e v_h(x) dx$ (is the same on X_h).

$$\Rightarrow \begin{cases} (u_h - u, r)_{L^2} - a_h(u_h - u, w_r) \stackrel{(*)}{\leq} Ch |w_r|_{H^2} \|u_h - u\|_{1,h} \leq \bar{C} h^2 \|r\|_{L^2} \|f\|_{L^2} \\ a_h(w_h) - b(w_r) - a_h(u, w_h - w_r) \stackrel{(*)}{\leq} Ch |u|_{H^2} \|w_h - w_r\|_{1,h} \leq \bar{C} h^2 \|f\|_{L^2} \|r\|_{L^2} \end{cases}$$

conditions from previous slide satisfied

$$\Rightarrow \|u_h - u\|_{L^2} \leq \sup_{r \in L^2(\Omega)} \frac{1}{\|r\|_{L^2}} \left(Ch^2 \|f\|_{L^2} \|r\|_{L^2} \right)$$

□

Duality: dual spaces

Def: Let X be a normed vector space. The dual space X^* is the space of all bounded linear functionals on X . For $f \in X^*$, $x \in X$ we also write $f(x) = \langle f, x \rangle_{X^*, X}$. A norm on X^* is defined by $\|f\|_{X^*} = \sup_{x \in X, x \neq 0} \frac{f(x)}{\|x\|_X}$.

Thm: By construction, $\langle f, x \rangle_{X^*, X} \leq \|f\|_{X^*} \|x\|_X$.

Ex: $L^p(\Omega)^* = L^q(\Omega)$ for $p \in [1, \infty)$ & $\frac{1}{p} + \frac{1}{q} = 1$. ($L^q(\Omega) \subset L^p(\Omega)^*$ by Hölder's inequality.)

Def: Let $\Omega \subset \mathbb{R}^n$ open. The space of test functions is $C_c^\infty(\Omega) = \{\phi \in C^\infty(\Omega) \mid \text{supp } \phi \text{ compact}\}$.

$\phi_j \xrightarrow{j \rightarrow \infty} \phi$ in $C_c^\infty(\Omega)$, if there exists a compact $K \subset \Omega$ with $\text{supp } \phi_j \subset K \forall j$ and $\|\phi_j - \phi\|_{C^k} \xrightarrow{j \rightarrow \infty} 0 \forall k$.

Ex: $\phi(x) = e^{-(1-|x|^2)^{-1}}$ for $|x| < 1$ & $\phi(x) = 0$ else

• The space of distributions is $\mathcal{D}'(\Omega) = \{f: C_c^\infty(\Omega) \rightarrow \mathbb{R} \mid f \text{ linear \& continuous}\}$ (topological dual space to $C_c^\infty(\Omega)$). Any $f \in L^1(\Omega)$ can be interpreted as a distribution by setting $f(\phi) := \int_\Omega f(x) \phi(x) dx$. The notation $\int_\Omega f \phi dx$ for $f(\phi)$ is even used if $f \notin L^1(\Omega)$. Ex: $\int_\Omega \delta \phi dx = \phi(0)$

• For $f \in \mathcal{D}'(\Omega)$ and multiindex α the distributional derivative $D^\alpha f \in \mathcal{D}'(\Omega)$ is defined as $D^\alpha f(\phi) = (-1)^{|\alpha|} f(D^\alpha \phi)$ (if $f \in C^{|\alpha|} \subset L^1 \subset \mathcal{D}'$, this equals the classical derivative!) Ex: $H^1 = \mathcal{D}'$

Finite Elements: elliptic PDEs (2D Poisson equation)

Proof possible with Lax-Milgram!

Duality: H^{-1}

Thm: $C_c^\infty(\Omega)$ is dense in $H_0^1(\Omega)$, i.e. $\forall u \in H_0^1(\Omega), \varepsilon > 0 \exists \phi \in C_c^\infty(\Omega) : \|u - \phi\|_{H^1} < \varepsilon$.

Thm: (Riesz representation theorem) A Hilbert space X is isometrically isomorphic to X^* ,
i.e. for every $f \in X^*$ there exists exactly one $g \in X$ with $\|f\|_{X^*} = \|g\|_X$ and $f(v) = (g, v)_X \forall v \in X$.

Ex: Every $f \in L^2(\Omega)^*$ can be written as $f(v) = \int_\Omega g v \, dx$ for a $g \in L^2(\Omega)$.

Def: $H^{-1}(\Omega) := H_0^1(\Omega)^*$, and we write $\langle f, u \rangle_{H^{-1}, H_0^1} = \int_\Omega f u \, dx$.

Thm: $H^{-1}(\Omega) = \{f \in \mathcal{D}'(\Omega) \mid \exists g \in H_0^1(\Omega) : \langle f, u \rangle_{\mathcal{D}', C_c^\infty} = (g, u)_{H^1} \forall u \in H_0^1(\Omega)\} =: F$.

Pf: Clearly, $H^{-1} \subset \mathcal{D}'$, and by Riesz, $\forall f \in H^{-1}(\Omega) \subset \mathcal{D}'(\Omega) \exists g \in H_0^1(\Omega) : \langle f, v \rangle_{\mathcal{D}', C_c^\infty} = \langle f, v \rangle_{H^{-1}, H_0^1} = (g, v)_{H^1} \forall v \in C_c^\infty(\Omega) \subset H_0^1(\Omega)$

\supset : For $f \in F, v \in H_0^1(\Omega)$ one can define $f(v) = \lim_{i \rightarrow \infty} f(v_i)$ for a sequence $v_i \in C_c^\infty(\Omega)$ with $v_i \xrightarrow{i \rightarrow \infty} v$ in $H_0^1(\Omega)$.

The limit exists since $f(v_i)$ is Cauchy due to $|f(v_i) - f(v_j)| = |(g, v_i - v_j)| \leq \|g\|_{H^1} \|v_i - v_j\|_{H^1} \xrightarrow{i, j \rightarrow \infty} 0$, and

it does not depend on v_i (since $v_i, \tilde{v}_i \rightarrow v \Rightarrow |f(v_i) - f(\tilde{v}_i)| \leq \|g\|_{H^1} \|v_i - \tilde{v}_i\|_{H^1} \rightarrow 0$);

also, this f is clearly linear with $|f(v)| = \lim_{i \rightarrow \infty} |f(v_i)| = \lim_{i \rightarrow \infty} |(g, v_i)_{H^1}| \leq \limsup_i \|g\|_{H^1} \|v_i\|_{H^1} = \|g\|_{H^1} \|v\|_{H^1}$. \square

Ex: $\cdot L^2(\Omega) \subset H^{-1}(\Omega)$ with $\langle f, u \rangle_{H^{-1}, H_0^1} = \int_\Omega f u \, dx$ for $f \in L^2(\Omega)$

$\cdot \delta \in H^{-1}((-1, 1))$: $f_1 : (x \mapsto \frac{1-|x|}{2}) \in H_0^1((-1, 1)) \subset L^2((-1, 1)) \subset H^{-1}((-1, 1))$

$\cdot f_2 = \delta + f_1 \in H^{-1}((-1, 1))$, since $\forall u \in C_c^\infty((-1, 1)) : f(u) = u(0) + \int_{-1}^1 \frac{1-|x|}{2} u(x) \, dx = (-f_1, u)_{H^1}$

Finite Elements: elliptic PDEs (2D Poisson equation)

Duality: elliptic differential operators

Thm: $H^{-1}(\Omega)^* = H_0^1(\Omega)$ with $\langle u, f \rangle_{H^{-1}(\Omega)^*, H^{-1}(\Omega)} = \langle f, u \rangle_{H^{-1}, H_0^1}$.

Pf: $H_0^1(\Omega) \subset H^{-1}(\Omega)^*$ (one always has $X \subset (X^*)^*$)

Let now $z \in H^{-1}(\Omega)^*$ and $R: H^{-1}(\Omega) \rightarrow H_0^1(\Omega)$ be the isomorphism from Riesz' theorem.

$u_z = R(z \circ R^{-1}) \in H_0^1(\Omega)$ satisfies $z(f) = (z \circ R^{-1})(Rf) = \langle z \circ R^{-1}, Rf \rangle_{H^{-1}, H_0^1} = \langle Rf, R(z \circ R^{-1}) \rangle_{H^{-1}} = \langle f, u_z \rangle = f(u_z) \forall f \in H^{-1}(\Omega)$. \square

Thm: Let $A \in L^\infty(\Omega; \mathbb{R}^{n \times n})$, $b \in L^q(\Omega; \mathbb{R}^n)$, $c \in L^1(\Omega)$. The differential operator $L: H_0^1(\Omega) \rightarrow H^{-1}(\Omega)$,
distributional derivative $L u = -\operatorname{div}(A \nabla u) + b \cdot \nabla u + c u$ is bounded *weak derivative*

Pf: Let $u \in H_0^1(\Omega)$, $v \in C_c^\infty(\Omega)$. $\langle Lu, v \rangle_{H^{-1}, H_0^1} = \int_\Omega \nabla v \cdot A \nabla u + v b \cdot \nabla u + c u v \, dx \leq C \|u\|_{H^1} \|v\|_{H^1}$

Let now $v \in H_0^1(\Omega)$, $v_j \in C_c^\infty(\Omega)$, $v_j \rightarrow v$. $\langle Lu, v \rangle_{H^{-1}, H_0^1} = \lim_{j \rightarrow \infty} \langle Lu, v_j \rangle_{H^{-1}, H_0^1}$ exists (why?),
is linear & bounded in $v \in H_0^1(\Omega)$ (why?), and $\|L\| \leq C$ (why?). Hint: same argument as previous slide \square

Def: Let X, Y be normed spaces, $L: X \rightarrow Y$ linear and bounded. The adjoint operator

$L^*: Y^* \rightarrow X^*$ is defined by $\langle f, Lx \rangle_{Y^*, Y} = \langle L^*f, x \rangle_{X^*, X}$.

Ex: $L: H_0^1(\Omega) \rightarrow H^{-1}(\Omega)$, $L u = -\operatorname{div}(A \nabla u) + b \cdot \nabla u + c u \Rightarrow L^* v = -\operatorname{div}(A^T \nabla v) - \operatorname{div}(b v) + c v$;

since $\int_\Omega (Lu) v \, dx = \int_\Omega u L^* v \, dx$ after integrating by parts twice.

Finite Elements: elliptic PDEs (2D Poisson equation)

Adaptivity: idea of adaptive FE-solutions

a priori error estimate $\|u - u_h\| \leq \eta_h(u)$

a posteriori error estimate $\|u - u_h\| \leq \eta_h(u_h)$

Def: An error estimator $\eta_h(u_h)$ is called reliable, if $\|u - u_h\| \leq \eta_h(u_h)$.

It is called efficient with constant c , if $c \eta_h(u_h) \leq \|u - u_h\|$.

An error estimator can often be written as a sum $\eta_h(u_h)^2 = \sum_{T \in \mathcal{T}_h} \eta_{T_e}(u_h)^2$ over all FE.

Given $TOL > 0$: $\eta_{T_e}(u_h) < \frac{TOL}{\sqrt{M}}$ (\rightarrow refinement criterion) $\Rightarrow \|u - u_h\| \leq \eta_h(u_h) < TOL$ (stopping criterion)

Alg. (Adaptive FE)

choose initial subdivision

while stopping criterion not satisfied

compute u_h & $\eta_h(u_h)$

for $l = 1, \dots, M$

if refinement criterion satisfied for T_e

refine T_e into smaller elements

Finite Elements: elliptic PDEs (2D Poisson equation)

Adaptivity: a posteriori error estimate

Def: The residual of an elliptic PDE $Lu = f$ with $L: H_0^1(\Omega) \rightarrow H^{-1}(\Omega)$, $f \in L^2(\Omega)$ is
 $R: H_0^1(\Omega) \rightarrow H^{-1}(\Omega)$, $R(v) = Lv - f$.

Let $a(u, v) = \langle Lu, v \rangle_{H^{-1}, H_0^1}$, $\ell(v) = (f, v)_{L^2}$, $a(v, v) \geq c \|v\|_{H^1}^2$, $a(u, v) \leq C \|u\|_{H^1} \|v\|_{H^1} \forall u, v \in H_0^1(\Omega)$
 $a(u, v) = \ell(v) \forall v \in H_0^1(\Omega)$, $a(u_h, v_h) = \ell(v_h) \forall v_h \in X_h \subset H_0^1(\Omega)$
e.g. $L = (-\Delta + id)$, $a(u, v) = \int_{\Omega} \nabla u \cdot \nabla v + uv \, dx$

Thm: $c \|u - u_h\|_{H^1(\Omega)} \leq \|R(u_h)\|_{H^{-1}(\Omega)} \leq C \|u - u_h\|_{H^1(\Omega)}$

Pf: $\langle R(u_h), v \rangle_{H^{-1}, H_0^1} = \langle Lu_h - f, v \rangle_{H^{-1}, H_0^1} = \langle L(u_h - u), v \rangle_{H^{-1}, H_0^1} = a(u_h - u, v) \quad \forall v \in H_0^1(\Omega)$
 $\|R(u_h)\|_{H^{-1}} = \sup_{v \neq 0} \langle R(u_h), v \rangle_{H^{-1}, H_0^1} / \|v\|_{H^1} \leq C \|u_h - u\|_{H^1}$
 $\|R(u_h)\|_{H^{-1}} \geq \langle R(u_h), u_h - u \rangle_{H^{-1}, H_0^1} / \|u_h - u\|_{H^1} \geq c \|u_h - u\|_{H^1} \quad \square$

Thus, $\frac{1}{c} \|R(u_h)\|_{H^{-1}(\Omega)}$ is a reliable & efficient error estimator.

Finite Elements: elliptic PDEs (2D Poisson equation)

Adaptivity: a posteriori error estimate in L^2

$u \in X$ solves $a(u, v) = l(v) \forall v \in X$; a bilinear, coercive, bounded; l linear, bounded

$$a(u, v) = \langle Lu, v \rangle_{H^1, H^1} = \langle u, L^*v \rangle_{H_0^1, H_0^1} \text{ for a differential operator } L: H_0^1(\Omega) \rightarrow H^{-1}(\Omega)$$

a & l such that $u \in H^2$ & $a(v, w) = (f, v)_{L^2} \forall v \in H_0^1(\Omega)$ implies $\|w\|_{H^2} \leq C \|f\|_{L^2}$

e.g. $\Omega = [0, 1]^2$, $a(u, v) = \int_{\Omega} \nabla u \cdot \nabla v + uv \, dx$, $l(v) = \int_{\Omega} f v \, dx$, $f \in L^2(\Omega)$, $L = (-\Delta + id)$

$u_h \in X_h \subset X$ solves $a(u_h, v_h) = l(v_h) \forall v_h \in X_h$ (finite elements on quasiuniform triangulation \mathcal{T})

Thm: $\|u - u_h\|_{L^2} \leq Ch \|R(u_h)\|_{H^{-1}}$

Pf \because Let w solve the dual problem $a(v, w) = (u - u_h, v)_{L^2} \forall v \in X$ (see Aubin-Nitsche)

$$\Rightarrow \|u - u_h\|_{L^2}^2 = (u - u_h, u - u_h)_{L^2} = a(u - u_h, w) = a(u - u_h, w - w_h) \quad \forall w_h \in X_h$$

$$= \langle Lu - Lu_h, w - w_h \rangle_{H^1, H_0^1} = R(u_h)(w_h - w) \leq \|R(u_h)\|_{H^{-1}} \|w_h - w\|_{H^1}$$

$$\stackrel{w_h = I_h w}{\leq} C \|R(u_h)\|_{H^{-1}} h \|w\|_{H^2} \leq C \|R(u_h)\|_{H^{-1}} h \|u - u_h\|_{L^2} \quad \square$$

Using the following techniques, this can be further refined so that on each triangle T the local h_T is used instead of h .

Finite Elements: elliptic PDEs (2D Poisson equation)

Adaptivity: local error estimator

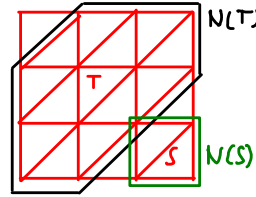
Let $a(u,v) = (f,v)_{L^2} \forall v \in H_0^1(\Omega)$ with $a(u,v) = \langle Lu, v \rangle_{H^{-1}, H_0^1}$ coercive & bounded, $Lu = -\operatorname{div}(A \nabla u) + b \cdot \nabla u + cu, f \in L^2(\Omega)$.

- 1) $\|R(u_h)\|_{H^{-1}}$ difficult to compute 2) we want error estimator for each element

\Rightarrow we will replace $\|R(u_h)\|_{H^{-1}}^2$ by $\eta_h(u_h)^2 = \sum_{T \in \mathcal{T}} \eta_T(u_h)^2$ with

$$\eta_T(u_h)^2 := h_T^2 \|Lu_h - f\|_{L^2(T)}^2 + h_T \| [A \nabla u_h \cdot n] \|_{L^2(\partial T)}^2 \quad [\cdot] = \text{jump across edge}$$

and show reliability and (almost) efficiency. Ingredients:



(a) $\frac{2}{\rho}$ -quasi-uniform triangulation \mathcal{T} ,

$$N(T) = \{T_2 \in \mathcal{T} \mid \overline{T_2} \cap \overline{T} \neq \emptyset\} \text{ for } T \in \mathcal{T}, \quad N(S) = \{T_2 \in \mathcal{T} \mid S = \overline{T_2} \cap \overline{T_j}\} \text{ for } S = \overline{T_2} \cap \overline{T_j}, \quad h_S = \max_{T \in N(S)} h_T$$

(b) Galerkin-orthogonality $\langle R(u_h), \varphi \rangle_{H^{-1}, H_0^1} = a(u_h - u, \varphi) = a(u_h - u, \varphi - \varphi_h)$

(c) \exists (quasi-)interpolation $J_h: H_0^1(\Omega) \rightarrow X_h \subset H_0^1(\Omega)$ such that $\forall K \in \mathcal{T}$ i.e. $\leq C(\rho) \dots$

$$\|J_h v - v\|_{H^{-1}(K)} \lesssim \|\nabla v\|_{L^2(N(K))} \quad \cdot \|J_h v - v\|_{L^2(K)} \lesssim h_K \|\nabla v\|_{L^2(N(K))} \quad (\text{see later})$$

(d) trace estimate $\|v\|_{L^2(S)} \lesssim h_T^{-\frac{1}{2}} \|v\|_{L^2(T)} + h_T^{\frac{1}{2}} \|\nabla v\|_{L^2(T)} \quad \forall T \in \mathcal{T}, S \in \partial T, v \in H^1(T)$ homework

(e) Poincaré-inequality $\|v\|_{L^2(T)} \lesssim h_T \|v\|_{H^1(T)} \quad \forall T \in \mathcal{T}, v \in H_0^1(T)$ homework

Finite Elements: elliptic PDEs (2D Poisson equation)

Adaptivity: reliability local error estimator

Theorem: Let u_h be the FE solution. There exists $C = C(\Omega, p) > 0$ with $\|R(u_h)\|_{H^{-1}}^2 \leq C \sum_{K \in \mathcal{T}} \eta_K(u_h)^2$.

Pf : $\langle R(u_h), \varphi \rangle_{H^{-1}, H_0^1} \stackrel{(b)}{=} a(u_h - u, \varphi - \varphi_h) = a(u_h, \varphi - \varphi_h) - (f, \varphi - \varphi_h)_{L^2(\Omega)}$

$$= \sum_{K \in \mathcal{T}} \int_K \nabla(\varphi - \varphi_h) \cdot A \nabla u_h + (\varphi - \varphi_h)(b \cdot \nabla u_h + c u_h - f) dx$$

$$= \sum_{K \in \mathcal{T}} (L u_h - f, \varphi - \varphi_h)_{L^2(K)} + \int_{\partial K} (\varphi - \varphi_h) (A \nabla u_h) \cdot n dx$$

Hölder

$$= \sum_{K \in \mathcal{T}} (L u_h - f, \varphi - \varphi_h)_{L^2(K)} + \frac{1}{2} \int_{\partial K} (\varphi - \varphi_h) [(A \nabla u_h) \cdot n] dx$$

$$\leq \sum_{K \in \mathcal{T}} \|L u_h - f\|_{L^2(K)} \|\varphi - \varphi_h\|_{L^2(K)} + \frac{1}{2} \|[A \nabla u_h \cdot n]\|_{L^2(\partial K)} \|\varphi - \varphi_h\|_{L^2(\partial K)}$$

pick $\varphi_h = f_h$

$$\stackrel{(a)}{\leq} \sum_{K \in \mathcal{T}} \|L u_h - f\|_{L^2(K)} \|\varphi - \varphi_h\|_{L^2(K)} + \|[A \nabla u_h \cdot n]\|_{L^2(\partial K)} (h_K^{-\frac{1}{2}} \|\varphi - \varphi_h\|_{L^2(K)} + h_K^{\frac{1}{2}} \|\varphi - \varphi_h\|_{H^1(K)})$$

Cauchy-Schwarz

$$\stackrel{(c)}{\leq} \sum_{K \in \mathcal{T}} (\|L u_h - f\|_{L^2(K)} h_K + \|[A \nabla u_h \cdot n]\|_{L^2(\partial K)} h_K^{\frac{1}{2}}) \|\nabla \varphi\|_{L^2(N(K))}$$

$$\leq \sqrt{\sum_{K \in \mathcal{T}} h_K^2 \|L u_h - f\|_{L^2(K)}^2 + h_K \|[A \nabla u_h \cdot n]\|_{L^2(\partial K)}^2} \sqrt{\sum_{K \in \mathcal{T}} \|\nabla \varphi\|_{L^2(N(K))}^2}$$

$$\sum_{K \in \mathcal{T}} \|\nabla \varphi\|_{L^2(N(K))}^2 \leq \sum_{K \in \mathcal{T}} \|\nabla \varphi\|_{L^2(K)}^2 = \|\nabla \varphi\|_{L^2(\Omega)}^2 \leq \|\varphi\|_{H^1(\Omega)}^2$$

$$\|R(u_h)\|_{H^{-1}} = \sup_{\varphi \in H_0^1(\Omega)} \frac{\langle R(u_h), \varphi \rangle_{H^{-1}, H_0^1}}{\|\varphi\|_{H^1}} \leq \sqrt{\sum_{K \in \mathcal{T}} \eta_K^2(u_h)}$$

□

Finite Elements: elliptic PDEs (2D Poisson equation)

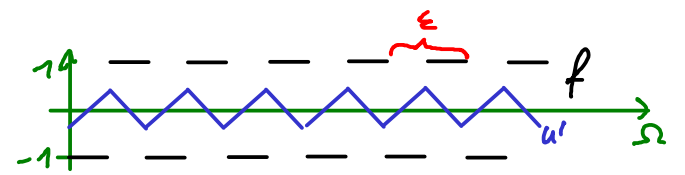
Adaptivity: efficiency local error estimator

Set $\cdot r \in L^2(\Omega)$, $r|_T = (Lu_h - f)|_T \quad \forall T \in \mathcal{T}$, $\cdot \bar{r} \in L^2(\Omega)$, $\bar{r}|_T = \frac{1}{|T|} \int_T r dx$
 $\cdot j \in L^2(S) \quad \forall \text{edges } S$, $j|_S = [A \nabla u_h \cdot n]$, $\cdot \bar{j} \in \mathcal{P}_1(S) \quad \forall S$ L^2 -projection of j
 $\Rightarrow \eta_T(u_h)^2 = h_T^2 \|r\|_{L^2(T)}^2 + h_T \sum_{S \in \partial T} \|j\|_{L^2(S)}^2$

Thm: $\eta_T(u_h)^2 \lesssim \|u_h - u\|_{H^1(N(T))}^2 + h_T^2 \|r - \bar{r}\|_{L^2(N(T))}^2 + h_T \|j - \bar{j}\|_{L^2(\partial T)}^2$. *local efficiency!*
 $\Rightarrow \eta_h(u_h)^2 = \sum_{T \in \mathcal{T}} \eta_T(u_h)^2 \lesssim \|u_h - u\|_{H^1(\Omega)}^2 + \sum_{T \in \mathcal{T}} (h_T^2 \|r - \bar{r}\|_{L^2(N(T))}^2 + h_T \|j - \bar{j}\|_{L^2(\partial T)}^2)$
 $\Rightarrow \eta_h(u_h)$ is efficient up to oscillation terms

The oscillation terms cannot be ignored:

Ex: $a(u, v) = \int_{\Omega} \nabla u \cdot \nabla v dx$, $a(u, v) = (f, v)_{L^2} \quad \forall v \in H_0^1(\Omega)$,



$a(u_h, v_h) = (f, v_h)_{L^2} \quad \forall v_h \in X_h = \text{linear FE-space}$.

Let $\epsilon \ll h_T \Rightarrow (f, v_h)_{L^2(\Omega)} = 0 \Rightarrow u_h = 0$, but $u'(x) = \begin{cases} x - \epsilon i, & x \in \epsilon i + [\epsilon/4, \epsilon/4], i \in \mathbb{N} \\ -x - \epsilon i/2, & \text{else} \end{cases}$

$\eta_h(u_h)^2 = h^2 \|f\|_{L^2(\Omega)}^2 = h^2 |\Omega|$

$\|u_h - u\|_{H^1}^2 \approx \|\nabla(u_h - u)\|_{L^2}^2 = \|\nabla u\|_{L^2}^2 = \frac{\epsilon^2}{3 \cdot 4^2} |\Omega| \not\approx \eta_h(u_h)^2$!

Finite Elements: elliptic PDEs (2D Poisson equation)

Adaptivity: efficiency local error estimator (proof)

Pf: $h_S^{\frac{1}{2}} \|j\|_{L^2(S)} \approx \|u_h - u\|_{H^1(N(S))} + h_S \|r - \bar{r}\|_{L^2(N(S))} + h_S^{\frac{1}{2}} \|j - \bar{j}\|_{L^2(S)}$, because...

• extend \bar{j} linearly from S to $N(S)$ such that $\|\bar{j}\|_{L^2(N(S))} + h_S \|\nabla \bar{j}\|_{L^2(N(S))} \approx h_S^{\frac{1}{2}} \|j\|_{L^2(S)}$ *homework*

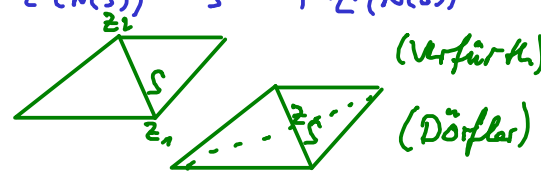
• Let $\varphi \in W_0^{1,\infty}(N(S))$ for edge S with

$$\varphi \geq 0, \int_S p^2 \varphi \, dx \sim \int_S p^2 \, dx \quad \forall p \in \mathcal{P}_1, \quad \|\varphi\|_{L^\infty(N(S))} + h_S \|\nabla \varphi\|_{L^\infty(N(S))} \lesssim 1$$

Ex: (edge-bubble) • $\varphi = \psi_{z_1}, \psi_{z_2}$ *(Verfürth)*

ψ_{z_2} = hat function in Z

• $\varphi = \psi_z$



$$\|\bar{j}\|_{L^2(S)}^2 \lesssim \int_S \bar{j} (\bar{j} \varphi) \, dx = \int_S j \psi_S \, dx + \int_S (\bar{j} - j) \psi_S \, dx$$

• $|\int_S j \psi_S \, dx| \lesssim \|r\|_{L^2(N(S))} \|\psi_S\|_{L^2(N(S))} + \|u_h - u\|_{H^1(N(S))} \|\psi_S\|_{H^1(N(S))}$, since

$$\int_S j \psi_S \, dx = a(u_h - u, \psi_S) - \int_{N(S)} r \psi_S \, dx$$

Hölder

• $\|\psi_S\|_{L^2(N(S))} \lesssim \|\bar{j}\|_{L^2(N(S))} \|\varphi\|_{L^\infty(N(S))} \lesssim h_S^{\frac{1}{2}} \|j\|_{L^2(S)}$ and

$$\begin{aligned} \|\psi_S\|_{H^1(N(S))} &\leq \|\psi_S\|_{L^2(N(S))} + \|\nabla \psi_S\|_{L^2(N(S))} \lesssim h_S^{\frac{1}{2}} \|j\|_{L^2(S)} + \|\bar{j} \nabla \varphi\|_{L^2(N(S))} + \|\varphi \nabla \bar{j}\|_{L^2(N(S))} \\ &\leq h_S^{\frac{1}{2}} \|j\|_{L^2(S)} + \|\bar{j}\|_{L^2(N(S))} \|\nabla \varphi\|_{L^\infty(N(S))} + \|\nabla \bar{j}\|_{L^2(N(S))} \|\varphi\|_{L^\infty(N(S))} \lesssim h_S^{-\frac{1}{2}} \|j\|_{L^2(S)} \end{aligned}$$

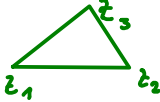
• $\|j\|_{L^2(S)} \leq \|j - \bar{j}\|_{L^2(S)} + \|\bar{j}\|_{L^2(S)} \leq \|j - \bar{j}\|_{L^2(S)} + \frac{\int_S j \psi_S \, dx}{\|\bar{j}\|_{L^2(S)}} + \|j - \bar{j}\|_{L^2(S)} \frac{\|\psi_S\|_{L^2(S)}}{\|\bar{j}\|_{L^2(S)}}$ *used (1)+(4), the above & following estimate for $\|r\|_{L^2(N(S))}$*


Adaptivity: efficiency local error estimator (proof cont.)

• $h_T \|r\|_{L^2(T)} \lesssim \|u_h - u\|_{H^1(T)} + h_T \|r - \bar{r}\|_{L^2(T)}$, because ...

• $h_T \|r\|_{L^2(T)} \leq h_T \|\bar{r}\|_{L^2(T)} + h_T \|r - \bar{r}\|_{L^2(T)}$

• Let $\psi \in W_0^{1,\infty}(T)$ with $\int_T \psi dx = 1$, $\|\nabla \psi\|_{L^\infty(T)} \lesssim h_T^{-1}$

Ex: (element bubble) • $\psi = c \psi_{z_1} \psi_{z_2} \psi_{z_3}$  (Verfärth)

• $\psi = c \psi_z$  (Dörfler)

$$\|\bar{r}\|_{L^2(T)}^2 = \int_T \bar{r} (\bar{r} \psi) dx \leq \|\bar{r}\|_{H^{-1}(T)} \|\bar{r} \psi\|_{H^1(T)} \lesssim \|\bar{r}\|_{H^{-1}(T)} \|\nabla(\bar{r} \psi)\|_{L^2(T)}$$

Hölder $\leq \|\bar{r}\|_{H^{-1}(T)} \|\bar{r}\|_{L^2(T)} \|\nabla \psi\|_{L^\infty(T)} \lesssim h_T^{-1} \|\bar{r}\|_{H^{-1}(T)} \|\bar{r}\|_{L^2(T)}$

$$\Rightarrow \|\bar{r}\|_{L^2(T)} \lesssim h_T^{-1} \|\bar{r}\|_{H^{-1}(T)} \leq h_T^{-1} \|r\|_{H^{-1}(T)} + h_T^{-1} \|r - \bar{r}\|_{H^{-1}(T)}$$

• $\|r\|_{H^{-1}(T)} = \sup_{\omega \in H_0^1(T)} \frac{\langle r, \omega \rangle_{H^{-1}, H^1}}{\|\omega\|_{H^1(T)}} = \sup_{\omega \in H_0^1(T)} \frac{a(u_h - u, \omega)}{\|\omega\|_{H^1(T)}} \lesssim \|u_h - u\|_{H^1(T)}$

• $\|r - \bar{r}\|_{H^{-1}(T)} = \sup_{\omega \in H_0^1(T)} \frac{\langle r - \bar{r}, \omega \rangle_{H^{-1}, H^1}}{\|\omega\|_{H^1(T)}} \leq \sup_{\omega \in H_0^1(T)} \frac{\|r - \bar{r}\|_{L^2(T)} \|\omega\|_{L^2(T)}}{\|\omega\|_{H^1(T)}} \stackrel{(e)}{\lesssim} h_T \|r - \bar{r}\|_{L^2(T)}$

• $\eta_T(u_h)^2 = h_T^2 \|r\|_{L^2(T)}^2 + h_T \sum_{S \in \partial T} \|j\|_{L^2(S)}^2$
 $\lesssim \|u_h - u\|_{H^1(T)}^2 + h_T^2 \|r - \bar{r}\|_{L^2(T)}^2 + \sum_{S \in \partial T} (\|u_h - u\|_{H^1(N(S))}^2 + h_S^2 \|r - \bar{r}\|_{L^2(N(S))}^2 + h_S \|j - \bar{j}\|_{L^2(S)}^2) \square$

Finite Elements: elliptic PDEs (2D Poisson equation)

Adaptivity: Clément interpolation

The Bramble-Hilbert lemma for nodal interpolation only holds if H^m (or more generally $W^{m,p}$) C^0 (where typically $l=0$). For all other cases, Clément interpolation was invented.

Thm: For $\Omega \subset \mathbb{R}^d$, $m \geq 0$, $p \in [1, \infty]$, $X_h = \{v \in C^0(\Omega) \mid v|_T \in \mathbb{P}_k(T) \forall T \in \mathcal{T}, v|_{\partial\Omega} = 0\}$

there exists a linear operator $\mathbb{I}_h : L^1(\Omega) \rightarrow X_h$ such that
 depends on quasiuniformity $\|D^t(v - \mathbb{I}_h v)\|_{L^q(T)} \lesssim h_T^{s - \frac{d}{p} - t + \frac{d}{q}} \|D^s v\|_{L^p(N(T))}$

$$\forall v \in W^{m,p}(\Omega), T \in \mathcal{T}, 0 \leq t \leq s \leq \min(m, k+1) \quad s - \frac{d}{p} \geq t - \frac{d}{q}.$$

In particular, $\|v - \mathbb{I}_h v\|_{L^2(T)} \lesssim h_T \|\nabla v\|_{L^2(N(T))}$ & $\|\nabla(v - \mathbb{I}_h v)\|_{L^2(T)} \lesssim \|\nabla v\|_{L^2(N(T))}$

Here for 1st order Lagrange finite elements for $H_0^1(\Omega)$:

Def: (Clément interpolation). M_i shall be the support of the piecewise affine hat function in node z_i :

$\cdot \mathbb{P}_{M_i} : L^2(M_i) \rightarrow X_h(M_i) = \{v_h \mid_{M_i} \mid v_h \in X_h\}$ is defined as the L^2 -orthogonal projection, i.e.

$$\|v - \mathbb{P}_{M_i} v\|_{L^2(M_i)} = \inf_{v_h \in X_h(M_i)} \|v - v_h\|_{L^2(M_i)}$$

why? $\mathbb{P}_{M_i} v$ is the solution $u_h \in X_h(M_i)$ of $(u_h, v_h)_{L^2(M_i)} = (v, v_h)_{L^2(M_i)} \forall v_h \in X_h(M_i)$, which is well-defined by L^2

$\cdot \mathbb{I}_h v = \sum_{z_i \in \mathcal{V}_\Omega} \chi_{M_i}(z_i) (\mathbb{P}_{M_i} v)(z_i) \psi_i$ for ψ_i the hat function in z_i

Thm: The Clément interpolation satisfies $\|v - \mathbb{I}_h v\|_{L^2(T)} \lesssim h_T \|\nabla v\|_{L^2(N(T))}$ & $\|\nabla(v - \mathbb{I}_h v)\|_{L^2(T)} \lesssim \|\nabla v\|_{L^2(N(T))}$

Finite Elements: elliptic PDEs (2D Poisson equation)

Adaptivity: Clément interpolation (proof)

Pf :: For an inner element T_e we have

triangle & Hölder ineq.

$$\|v - \mathcal{I}_h v\|_{L^2(T_e)} \stackrel{\text{partition of unity}}{=} \left\| \sum_i (v - P_{M_i} v(z_i)) \psi_i \right\|_{L^2(T_e)} \stackrel{!}{\leq} \sum_i \|v - P_{M_i} v(z_i)\|_{L^2(T_e)}$$

$$\stackrel{\text{sum over triangle nodes}}{\leq} \sum_i \|v - P_{M_i} v\|_{L^2(T_e)} + \|P_{M_i} v - P_{M_i} v(z_i)\|_{L^2(T_e)} \leq \sum_i \|v - P_{M_i} v\|_{L^2(M_i)} + \|P_{M_i} v - P_{M_i} v(z_i)\|_{L^2(T_e)}$$

• with $\bar{v}_M = \frac{1}{|M|} \int_M v \, dx$ one has *pull back to reference domain, Poincaré, transformation back to M*

$$\|v - P_M v\|_{L^2(M)}^2 \leq \|v - \bar{v}_M\|_{L^2(M)}^2 \stackrel{!}{\leq} h_M^2 \|\nabla(v - \bar{v}_M)\|_{L^2(M)}^2 = h_M^2 \|\nabla v\|_{L^2(M)}^2$$

$$\|P_{M_i} v - P_{M_i} v(z_i)\|_{L^2(T_e)} \leq \sum_k \|P_{M_i} v(z_k) - P_{M_i} v(z_i)\|_{L^2(T_e)} = \sum_k |T_e|^{1/2} |P_{M_i} v(z_k) - P_{M_i} v(z_i)|$$

$$\leq \sum_k |T_e|^{1/2} h_{T_e} \|\nabla P_{M_i} v\|_{L^\infty(T_e)} = \sum_k h_{T_e} \|\nabla P_{M_i} v\|_{L^2(T_e)} \lesssim h_{T_e} \|\nabla P_{M_i} v\|_{L^2(T_e)}$$

"inverse inequality" (homework)

$$\|\nabla P_M v\|_{L^2(T_e)} = \|\nabla(P_M v - \bar{v}_M)\|_{L^2(T_e)} = \|\nabla P_M(v - \bar{v}_M)\|_{L^2(T_e)} = \|\nabla P_M(v - \bar{v}_M)\|_{L^2(M)} \stackrel{!}{\leq} \frac{1}{h_M} \|P_M(v - \bar{v}_M)\|_{L^2(M)}$$

$$\|P_M\| \leq 1 \implies \|v - \bar{v}_M\|_{L^2(M)} / h_M \lesssim \|\nabla v\|_{L^2(M)}$$

$$\|\nabla(v - \mathcal{I}_h v)\|_{L^2(T_e)} = \left\| \sum_i (\nabla v) \psi_i + (v - P_{M_i} v(z_i)) \nabla \psi_i \right\|_{L^2(T_e)} \lesssim \sum_i \|\nabla v\|_{L^2(T_e)} + \|v - P_{M_i} v(z_i)\|_{L^2(T_e)} / h_{T_e}$$

$$\lesssim \sum_i \|\nabla v\|_{L^2(M_i)} \lesssim \|\nabla v\|_{L^2(N(T_e))} \quad \square$$

Finite Elements: elliptic PDEs (2D Poisson equation)

Perspective: nonlinear elliptic PDE - elasticity

A classical example of a nonlinear elliptic PDE comes from the description of the deformation of elastic (material) bodies.

Def: - In the undeformed, stress-free state the elastic body occupies a sufficiently smooth, open, bounded, connected domain $\Omega \subset \mathbb{R}^3$. This is called the reference configuration.

- The new position of a point $x \in \Omega$ after the deformation is denoted $y(x)$. $y: \Omega \rightarrow \mathbb{R}^3$ is called the deformation, $F = Dy: \Omega \rightarrow \mathbb{R}^{3 \times 3}$ is called the deformation gradient.



- The coordinate $x \in \Omega$ is called Lagrangian coordinate, i.e. every considered quantity (material density, elastic forces, etc.) at a position $y(x)$ in the deformed material is represented as a function of the original position x of the material point.

- The representation of considered quantities as a function of the new position $y(x)$ is called representation in Eulerian coordinates.

Perspective: nonlinear elliptic PDE - energy density

Def: A material is called hyperelastic if the total work performed to deform the material is stored in form of mechanical potential energy within the material and if the energy density at location x only depends on x and $F(x)$.

The stored energy thus has the form $\int_{\Omega} W(x, F(x)) dx$ with the stored energy function $W: \Omega \times \mathbb{R}^{3 \times 3} \rightarrow \mathbb{R}$.

W is called - homogeneous if $W(x, F) = W(F) \forall x \in \Omega, F \in \mathbb{R}^{3 \times 3}$

i.e. material is equally stiff in all directions

- orientation-preserving if $W(x, F) = \infty$ for $\det F \leq 0$

- isotropic if $W(x, F) = W(x, FR) \forall R \in SO(3)$

i.e. rigid rotation / translation costs no energy

- rigid body motion invariant if $W(x, I) = 0$ & $W(x, RF) = W(x, F) \forall R \in SO(3)$

- of p-growth if $\exists C > 0: W(x, F) \geq C \|F\|^p - C$

In the following, W shall have all these properties.

Frobenius norm $\|F\|^2 = \text{tr}(F^T F)$

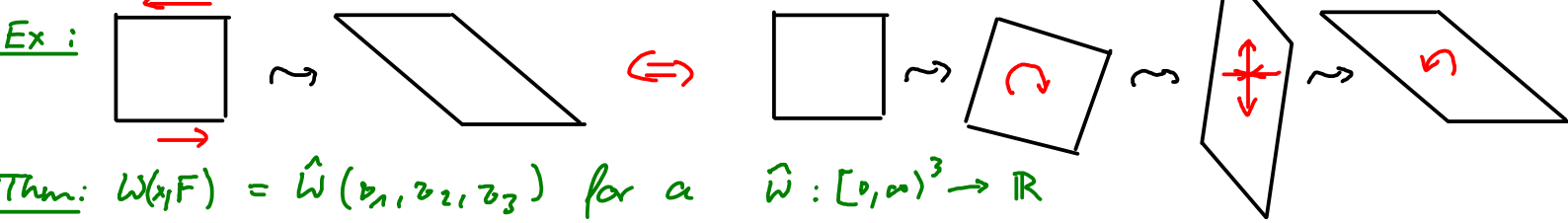
Ex: $W(F) = a \|F\|^p + b \det F^{-1}$ if $\det F > 0$, ∞ else

Finite Elements: elliptic PDEs (2D Poisson equation)

Perspective: nonlinear elliptic PDE - principal stretches

Rem: By first order Taylor expansion, any deformation y can be written as an affine transformation $y(z) \approx y(x) + F(x)(z-x)$ in a neighbourhood of a point x . $|F|$ describes the new length of a vector $v \in \mathbb{R}^2$ after the deformation, $|\text{cof } F|$ ($\text{cof } A = \det A A^{-T}$) the area change of the plane with normal n , $\det F$ the volume change.

Def: Let $F(x) = U \Sigma V^T$ be the singular value decomposition, then $y(z) \approx y(x) + F(x)(z-x)$ is a rotation V , subsequent stretch/compression along the coordinate directions by the singular values b_1, b_2, b_3 , a subsequent rotation U , and a final translation. b_1, b_2, b_3 are called the principal stretches.



Thm: $W(x, F) = \hat{W}(b_1, b_2, b_3)$ for a $\hat{W}: [0, \infty)^3 \rightarrow \mathbb{R}$

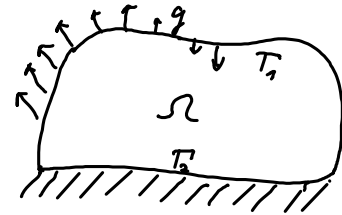
Pf: $W(x, F) = W(x, U^T F V^T) = W(x, \Sigma) = W(\Sigma)$ □

I.e. the stored energy only depends on the principal stretches!

Finite Elements: elliptic PDEs (2D Poisson equation)

Perspective: nonlinear elliptic PDE - variational problem

Let the reference configuration have Lipschitz boundary, the body shall be fixed at $\Gamma_2 \subset \partial\Omega$ with positive measure, at $\Gamma_1 = \partial\Omega \setminus \Gamma_2$ a force density (stress = force per area) $g: \Gamma_1 \rightarrow \mathbb{R}^3$ is applied.



if this is violated, set $E[y] = \infty$

The free energy of a deformation y with $y|_{\Gamma_2} = \text{id}|_{\Gamma_2}$ is $E[y] = \int_{\Omega} W(Dy(x)) dx - \int_{\Gamma_1} y(x) \cdot g(x) dx$.

The resulting deformation is a minimizer of the free energy.

minimization problem: $\min \{ E[y] \mid y|_{\Gamma_2} = \text{id}|_{\Gamma_2} \}$

most natural form of elliptic problems

weak formulation: $y \in W^{1,p}(\Omega)$ satisfies $y|_{\Gamma_2} = \text{id}|_{\Gamma_2}$ &

$A:B = \text{tr}(A^T B)$ — $\int_{\Omega} DW(Dy) : D\phi dx - \int_{\Gamma_1} \phi \cdot g dx = 0$

$\forall \phi \in W^{1,p}(\Omega)$ with $\phi|_{\Gamma_2} = 0$

strong formulation: y satisfies $\text{div}(DW(Dy)) = 0$

in Ω conservation of linear momentum

rowwise divergence

$y(x) = x$

on Γ_2

1. Piola-Kirchhoff-stress tensor

$DW(Dy) \cdot n = g$

on Γ_1

Perspective: nonlinear elliptic PDE - weak convergence

Def: Let X, X^* be normed vector space & dual space, $(x_i)_{i \in \mathbb{N}} \subset X, (y_i)_{i \in \mathbb{N}} \subset X^*$.

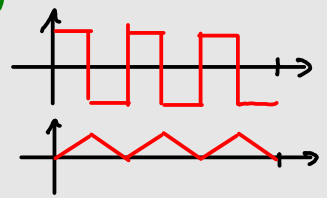
- x_i converges weakly to $x \in X$ ($x_i \rightharpoonup x$) if $\langle x_i, y \rangle_{X, X^*} \xrightarrow{i \rightarrow \infty} \langle x, y \rangle \forall y \in X^*$
- y_i converges weakly-* to $y \in X^*$ ($y_i \xrightarrow{*} y$) if $\langle x, y_i \rangle_{X, X^*} \xrightarrow{i \rightarrow \infty} \langle x, y \rangle \forall x \in X$
- X is called reflexive if $(X^*)^*$ is isometrically isomorphic to X .
- X is called separable if it contains a countable dense subset.

Ex: - $L^p(\Omega)$ is reflexive & separable for $p \in (1, \infty)$ ($(L^p)^{**} = (L^p)^* = L^q$ for $\frac{1}{p} + \frac{1}{q} = 1$)

• $W^{m,p}(\Omega)$ is reflexive & separable for $p \in (1, \infty)$

• $x_i = \text{sgn} \circ \sin(i \cdot) \rightarrow 0$ in $L^p(0,1)$

• $\tilde{x}_i = \int x_i dt \rightarrow 0$ in $W^{1,p}(0,1)$



Thm: $x_i \rightarrow x \Rightarrow x_i \rightharpoonup x$

• $x_i \rightarrow x \Rightarrow \{x_i\}$ is bounded

• $f_i \rightarrow f$ in $W^{1,p}(\Omega), p > \dim(\Omega) \Rightarrow f_i \rightarrow f$ in $C^0(\Omega)$ for subsequences

"Sobolev-embedding"



Finite Elements: elliptic PDEs (2D Poisson equation)

Perspective: nonlinear elliptic PDE - replacement for Lax-Milgram

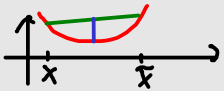
Thm: (Banach-Alaoglu) Let $A \subset X^*$ be a bounded subset of the dual space to a separable vector space X , then A contains a weakly-* converging sequence.

Cor: For X reflexive & X^* separable, any bounded $A \subset X^*$ contains a weakly converging sequence.

Thm: (Ball) Let $f_n \rightarrow f$ in $W^{1,p}(\Omega; \mathbb{R}^3)$ with $p > 3$. Then $\det Df_n \rightarrow \det Df$ in $L^{\frac{p}{2}}(\Omega)$ and $\text{cof } Df_n \rightarrow \text{cof } Df$ in $L^{p/2}(\Omega; \mathbb{R}^{3 \times 3})$.
 = collection of all 2x2 subdeterminants

Thm: (Mazur) Let $x_n \rightarrow x$, then there exists a sequence $\tilde{x}_n = \sum_{i=1}^n a_{in} x_i$ of convex combinations (i.e. $a_{in} \in [0,1], \sum_{i=1}^n a_{in} = 1$) with $\tilde{x}_n \rightarrow x$.

Thm: (Jensen) Let $f: X \rightarrow \mathbb{R}$ convex (i.e. $f(\frac{x+\tilde{x}}{2}) \leq \frac{f(x)+f(\tilde{x})}{2} \forall x, \tilde{x} \in X$) and $\sum_{i=1}^n a_i x_i$ a convex combination, then $f(\sum_{i=1}^n a_i x_i) \leq \sum_{i=1}^n a_i f(x_i)$.



Thm: Let $f_n \rightarrow f$ in $L^1(\Omega)$, then there exists a subsequence $f_{n_k} \rightarrow f$ pointwise almost everywhere.

Rem: Subsequence is important! $f_n^k(x) = \begin{cases} 1, & x \in [\frac{k}{n}, \frac{k+1}{n}] \\ 0 & \text{else} \end{cases}$; $f_1 = h_1, f_2 = f_2^0, f_3 = h_2^1, f_4 = h_3^0, \dots \xrightarrow{L^1} 0$

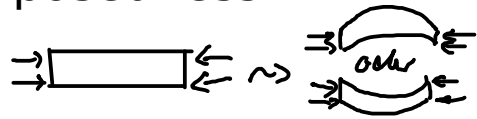
but not pointwise a.e.



Thm: (Fatou) Let $f_n \geq 0$, then $\int_{\Omega} \liminf_{n \rightarrow \infty} f_n dx \leq \liminf_{n \rightarrow \infty} \int_{\Omega} f_n dx$

Perspective: nonlinear elliptic PDE - well-posedness

Rem: In general, a solution is not unique, e.g.



Def: W is called polyconvex if it can be written as a convex function of the subdeterminants of F , $W(F) = \hat{W}(F, \text{cof} F, \det F)$ with \hat{W} convex.

$f: X \rightarrow \mathbb{R}$ is called lower semi-continuous, if $f(x) \leq \liminf_{n \rightarrow \infty} f(x_n) \forall x_n \rightarrow x$



Assumption: W shall be nonnegative, lower semi-continuous, polyconvex, of p -growth with $p > 3$.

Thm: $y \mapsto \int_{\Omega} W(Dy) dx$ is lower semi-continuous with respect to weak convergence in $W^{1,p}(\Omega)$.

Pf: Let $y_n \rightarrow y$ in $W^{1,p}(\Omega)$; wlog let $\liminf_{n \rightarrow \infty} \int_{\Omega} W(Dy_n) dx = \lim_{n \rightarrow \infty} \int_{\Omega} W(Dy_n) dx$.

else consider subsequence

Ball
 $\Rightarrow \omega_n := (Dy_n, \text{cof} Dy_n, \det Dy_n) \xrightarrow{L^p \times L^{p/2} \times L^{p/3}} (Dy, \text{cof} Dy, \det Dy) =: \omega$

Mazur
 $\Rightarrow \tilde{\omega}_n = \sum_{i=1}^{N_n} a_{in} \omega_i \xrightarrow{L^p \times L^{p/2} \times L^{p/3}} \omega \Rightarrow \tilde{\omega}_n \rightarrow \omega$ pointwise a.e.

$\int_{\Omega} W(Dy) dx = \int_{\Omega} \hat{W}(\omega) dx \leq \int_{\Omega} \liminf_{i \rightarrow \infty} \hat{W}(\tilde{\omega}_n) dx \stackrel{\text{Fatou}}{\leq} \liminf_{i \rightarrow \infty} \int_{\Omega} \hat{W}(\tilde{\omega}_n) dx$

Fubini
 $\leq \liminf_{i \rightarrow \infty} \int_{\Omega} \sum_{j=1}^{N_n} a_{jn} \hat{W}(\omega_j) dx = \liminf_{i \rightarrow \infty} \sum_{j=1}^{N_n} a_{jn} \int_{\Omega} W(Dy_j) dx = \lim_{n \rightarrow \infty} \int_{\Omega} W(Dy_n) dx \square$

Finite Elements: elliptic PDEs (2D Poisson equation)

Perspective: nonlinear elliptic PDE - well-posedness II

Thm: Let $g \in L^q(\Gamma_2)$, $\frac{1}{q} + \frac{1}{p} = 1$. E has a minimizer y with $y|_{\Gamma_2} = \text{id}|_{\Gamma_2}$.

Pf: "direct method of the calculus of variations"

0) $E \neq \infty$ and E is bounded below

$$E[\text{id}] < \infty \quad \& \quad E[y] \geq C \|Dy\|_{L^p}^p - C \underbrace{\int_{\Gamma_2} g \cdot y \, dx}_{\text{Poincaré}} \geq \hat{c} \|y\|_{W^{1,p}}^p - C \hat{c} \|g\|_{L^q} \|y\|_{W^{1,p}} \quad (*)$$

$$\leq \|g\|_{L^q(\Gamma_2)} \|y\|_{L^p(\Gamma_2)} \leq \hat{c} \|g\|_{L^q(\Gamma_2)} \|y\|_{W^{1,p}(\Omega)} \geq -C \left(1 - \frac{\hat{c}}{p^{1/(p-1)}}\right)^{p-1} \sqrt{\frac{(\hat{c} \|g\|_{L^q})^p}{\hat{c}}}$$

1) Consider "minimizing sequena" y_1, y_2, \dots with $E[y_n] \rightarrow \inf_y E[y]$

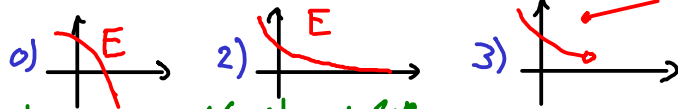
2) Show that a subsequence converges to some y (in a sense to be chosen adequately):

$$E[y_n] < C \stackrel{(*)}{\Rightarrow} \|y_n\|_{W^{1,p}} \text{ is bounded} \Rightarrow y_{n_i} \rightarrow y \text{ in } W^{1,p}(\Omega; \mathbb{R}^2)$$

3) Show lower semi-continuity w.r.t. the above convergence, i.e. $E[y] \leq \liminf_{i \rightarrow \infty} E[y_{n_i}] \leq \inf E$

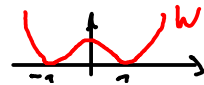
previous thm. + $\int_{\Gamma_2} g \cdot y_n \, dx \rightarrow \int_{\Gamma_2} g \cdot y \, dx$ (trace theorem) □

Ex: How can direct method fail?



Rem: (Poly-)convexity is essential! E.g. $y_n(x) = \frac{|nx - \text{round}(nx)|}{n} \xrightarrow{W^{1,p}} 0 \equiv y_1$

but $\int_0^1 W(y_n') \, dx = 0 \neq 1 = \int_0^1 W(y_1') \, dx$ for $W(a) = (a^2 - 1)^2$!



Finite Elements: elliptic PDEs (2D Poisson equation)

Perspective: nonlinear elliptic PDE - discretization

As usual for finite elements we choose a finite element space X_h and define the discrete approximation y_h to $y \in \operatorname{argmin} E$ as $y_h \in \operatorname{argmin}_{X_h} E$.

In the following we use first order Lagrange elements.

Thm: E has a minimizer in X_h .

Pf: Direct method in X_h ; it only remains to show that $y_h \in X_h, y_h \xrightarrow{w.r.t} y \Rightarrow y \in X_h$.
This follows from $y_h \rightarrow y$ in C^0 . □

Does y_h converge to y as $h \rightarrow 0$?

Define $E_h[y] = \begin{cases} E[y] & \text{if } y \in X_h \\ \infty & \text{else} \end{cases}$

$\Rightarrow y_h \in \operatorname{argmin} E_h[y]$

Perspective: nonlinear elliptic PDE - Gamma convergence

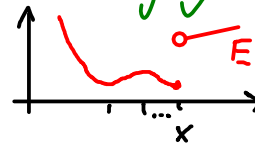
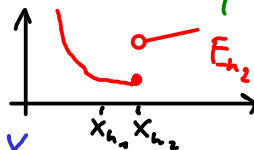
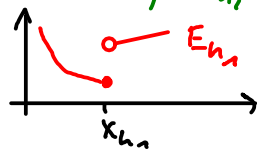
Def: Fix a notion of convergence on space X . A functional $E_h : X \rightarrow \mathbb{R}$ Γ -converges to $E : X \rightarrow \mathbb{R}$ ($\Gamma(X) - \lim_{h \rightarrow 0} E_h = E$), if $\forall h_i \rightarrow 0$

• $\forall x_i \xrightarrow{i \rightarrow \infty} x : \liminf_{i \rightarrow \infty} E_{h_i}[x_i] \geq E[x]$ "liminf-inequality"

• $\forall x \in X \exists x_i \xrightarrow{i \rightarrow \infty} x : \limsup_{i \rightarrow \infty} E_{h_i}[x_i] \leq E[x]$ "limsup-inequality"

• The functionals E_h are called equicoercive if $\exists C \subset X$ sequentially compact with $\text{argmin}_h E_h \subset C \forall h$.

Thm: If $\Gamma(X) - \lim_{h \rightarrow 0} E_h = E$ and the E_h are equicoercive, then any sequence x_{h_i} of minimizers of E_{h_i} possesses a subsequence converging to a minimizer of E .



↓
but no convergence rate!
This would depend on the shape of the energy in the minimum (e.g. locally quadratic or quartic).

Pf ∵ equicoercivity $\Rightarrow x_{h_i} \xrightarrow{i \rightarrow \infty} x \in X$ for a subsequence

• for all $\hat{x} \in X$ there exists $x_i \xrightarrow{i \rightarrow \infty} \hat{x}$ with

$$E[x] \geq \limsup_{i \rightarrow \infty} E_{h_i}[x_i] \geq \limsup_{i \rightarrow \infty} E_{h_i}[x_{h_i}] \geq \liminf_{i \rightarrow \infty} E_{h_i}[x_{h_i}] \geq E[x] \quad \square$$

Perspective: nonlinear elliptic PDE - convergence

- Thm: Let
- Ω Lipschitz, $\Gamma_1, \Gamma_2 = \partial\Omega \setminus \Gamma_1$ Lipschitz,
 - $W: \mathbb{R}^{3 \times 3} \rightarrow \mathbb{R}$ polyconvex, continuous, of p -growth with $p > 3$,
 - $W(F) \leq C|F|^p + c$ for a $C > 0$, ← usually we want $W(F) = \infty$ if $\det F \leq 0$, but this makes argument much more complicated
 - $g \in L^q(\Gamma_1)$, $\frac{1}{q} + \frac{1}{p} = 1$,
 - X_h $\frac{1}{p}$ -quasiuniform first order Lagrange finite element space with grid width h ,
 - $E: W^{1,p}(\Omega) \rightarrow \mathbb{R} \cup \{\infty\}$, $E[y] = \begin{cases} \int_{\Omega} W(\nabla y) dx - \int_{\Gamma_1} g \cdot y dx, & \text{if } y|_{\Gamma_2} = \text{id} \\ \infty & \text{else} \end{cases}$
 - $E_h: W^{1,p}(\Omega) \rightarrow \mathbb{R} \cup \{\infty\}$, $E_h[y] = \begin{cases} E[y], & \text{if } y \in X_h \\ \infty & \text{else} \end{cases}$.

Then Γ (weak convergence in $W^{1,p}(\Omega)$) $\text{-}\lim_{h \rightarrow 0} E_h = E$.

Cor: Every sequence $y_{n_i} \in \text{argmin } E_{h_i}$, $h_i \xrightarrow{i \rightarrow \infty} 0$, contains a subsequence $y_{h_{i_n}} \xrightarrow[h \rightarrow 0]{W^{1,p}} y \in \text{argmin } E$.

Pf: $E[\text{id}] \geq E_{h_i}[y_{h_i}] = E[y_{h_i}] \xrightarrow{(*)} \|y_{h_i}\|_{W^{1,p}}$ is uniformly bounded

$\Rightarrow E_{h_i}$ is equicoercive wrt. weak convergence in $W^{1,p}(\Omega)$. □

Perspective: nonlinear elliptic PDE - convergence proof

Pf : Let $h_i \rightarrow 0$.

liminf-ineq. : Let y_i be a sequence with $y_i \xrightarrow{W^{1,p}} y$, then the lower semi-continuity of E implies:

$$E[y] \leq \liminf_{i \rightarrow \infty} E[y_i] \leq \liminf_{i \rightarrow \infty} E_{h_i}[y_i]$$

limsup-ineq. : Let $y \in W^{1,p}(\Omega)$ be given.

- $W^{2,p}(\Omega)$ is dense in $W^{1,p}(\Omega) \Rightarrow \forall \epsilon \in \mathbb{N} \exists y^\epsilon \in W^{2,p}(\Omega)$ with $\|y^\epsilon - y\|_{W^{1,p}} \leq \frac{1}{\epsilon}$.
- Set $y_i^\epsilon = I_{h_i} y^\epsilon \Rightarrow \|y_i^\epsilon - y\|_{W^{1,p}} \leq \|y_i^\epsilon - y^\epsilon\|_{W^{1,p}} + \|y^\epsilon - y\|_{W^{1,p}} \leq Ch_i \|y^\epsilon\|_{W^{2,p}} + \frac{1}{\epsilon}$
- Let i_n such that $Ch_{i_n} \|y^\epsilon\|_{W^{2,p}} \leq \frac{1}{n}$, and define the sequence $y_i \in X_{h_i}$, $i \in \mathbb{N}$, as $(y_1^1, y_2^1, \dots, y_{i_2-1}^1, y_{i_2}^2, y_{i_2+1}^2, \dots, y_{i_3-1}^2, y_{i_3}^3, y_{i_3+1}^3, \dots)$
- $y_i \xrightarrow{W^{1,p}} y$ (thus also $y_i \rightarrow y$), since $\|y_i - y\|_{W^{1,p}} \leq \frac{2}{n}$ for all y_i after $y_{i_n}^\epsilon$
- choose subsequence y_{i_k} such that $\lim_{k \rightarrow \infty} E_{h_{i_k}}[y_{i_k}] = \limsup_{i \rightarrow \infty} E_{h_i}[y_i]$ and such that $D y_{i_k} \rightarrow Dy$ pointwise a. e.
- still to show : $\lim_{k \rightarrow \infty} E_{h_{i_k}}[y_{i_k}] \leq E[y]$

Perspective: nonlinear elliptic PDE - convergence proof (cont.)

We have $\cdot W(Dy_{i_k}) \rightarrow W(Dy)$ pointwise a.e.

$\cdot C|Dy_{i_k}|^p + C \xrightarrow{L^1} C|Dy|^p + C$, since

$$\int_{\Omega} |C|Dy_{i_k}|^p + C - C|Dy|^p - C| dx = C \int_{\Omega} ||Dy_{i_k}|^p - |Dy|^p| dx \xrightarrow{k \rightarrow \infty} 0$$

$y_{i_k} \rightarrow y$ in $W^{1,p}$

$\cdot \int_{\Omega} W(Dy_{i_k}) dx \rightarrow \int_{\Omega} W(Dy) dx$ by Lebesgue's dominated convergence

$\cdot \int_{\Gamma_n} g \cdot y_{i_k} dx \rightarrow \int_{\Gamma_n} g \cdot y dx$, since $y_{i_k}|_{\Gamma_n} \xrightarrow{L^p} y|_{\Gamma_n}$

$$\Rightarrow E_{h_{i_k}}[y_{i_k}] \rightarrow E[y]$$

□

