

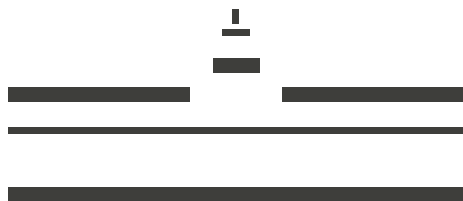
Matching NLO calculations with PS: recent developments in POWHEG BOX

Tomáš Ježo

Università di Milano-Bicocca

INFN, Sezione di Milano-Bicocca

Münster 25 November 2015



WESTFÄLISCHE
WILHELMS-UNIVERSITÄT
MÜNSTER



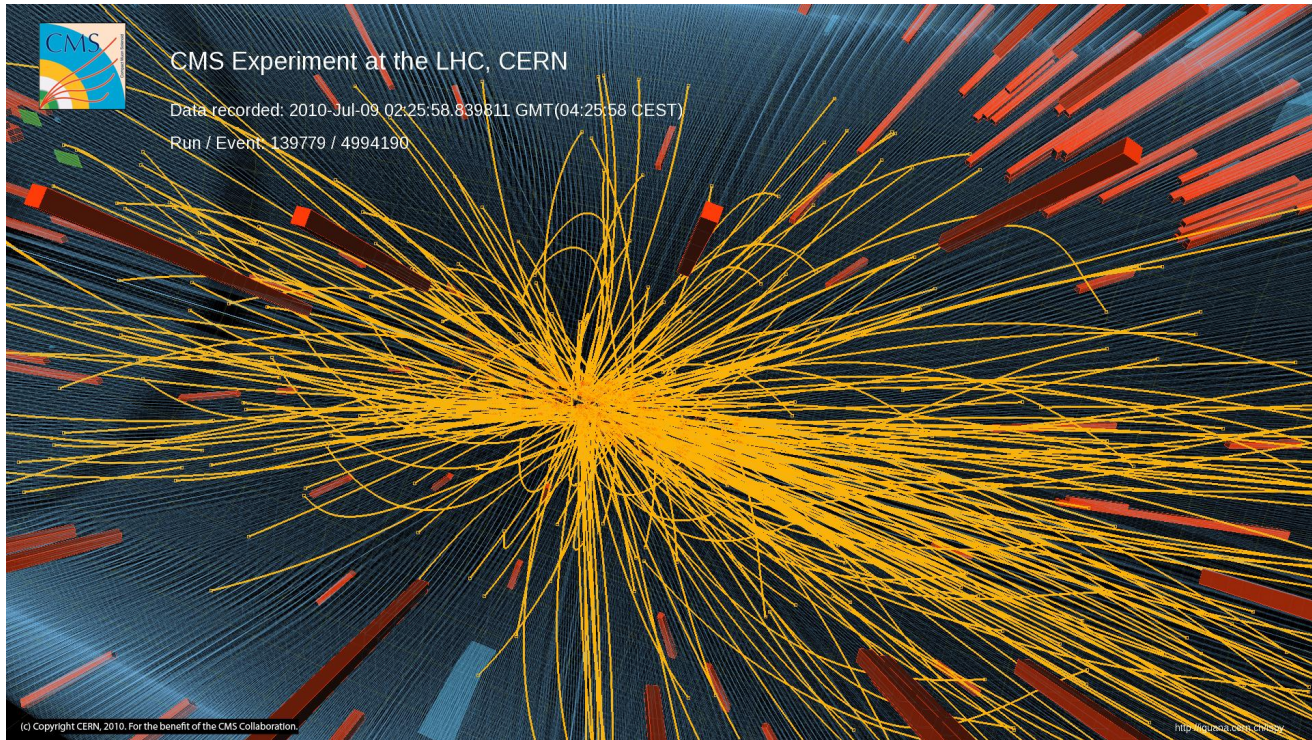
Matching NLO calculations with PS: recent developments in POWHEG BOX

- ▶ Matching fixed **N**ext-to-**L**eading **O**rders (NLO) calculations with **P**arton **S**hower (PS): **NLO+PS**
- ▶ Explain all the ingredients of a calculation at NLO+PS accuracy:
 - ▷ fixed order (FO): LO, NLO (real/virtual corrections)
 - ▷ parton shower (PS)
 - ▷ PS applied to NLO: NLO+PS
- ▶ Recent developments:
 - ▷ core POWHEG BOX: treatment of resonances
 - ▷ new processes: top-pair beyond SM, photon production



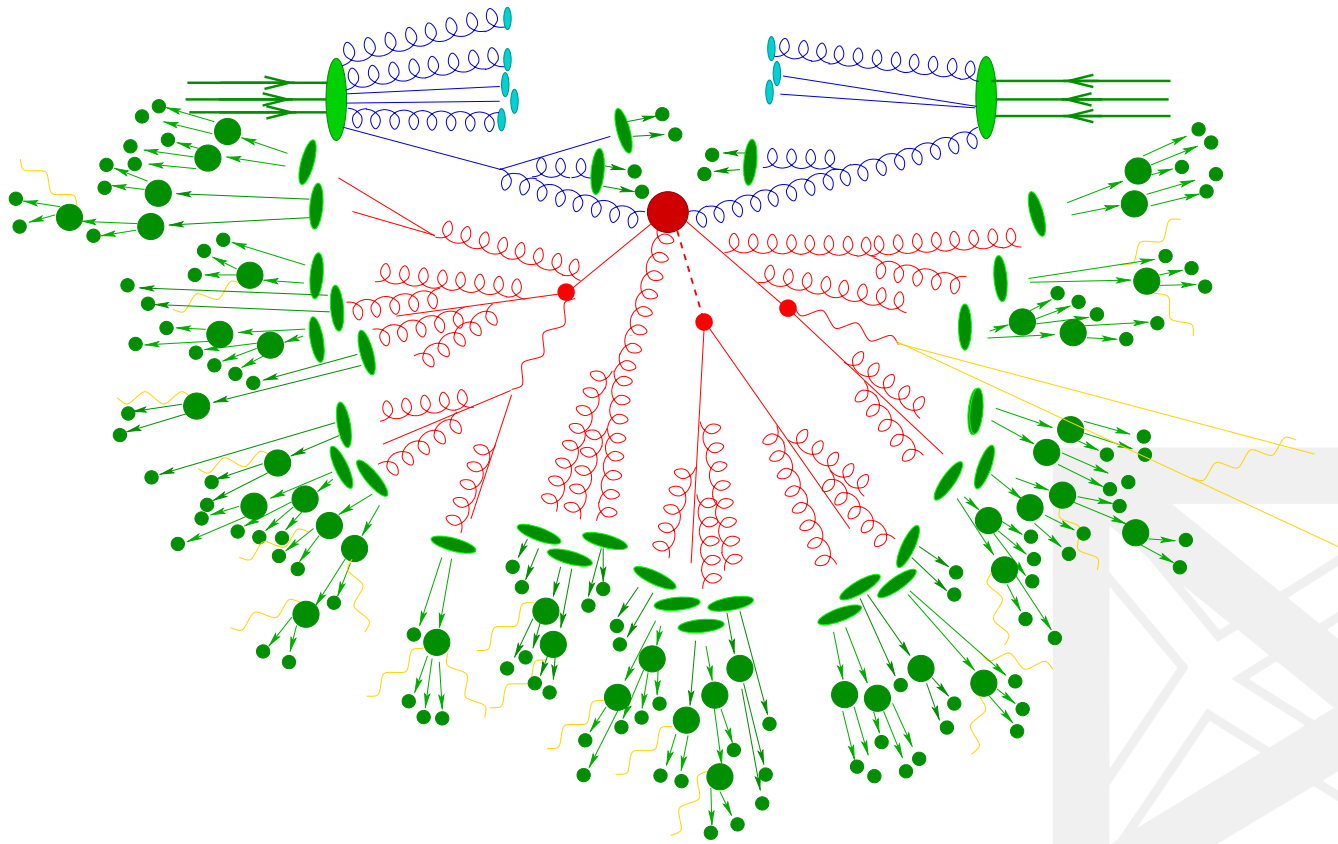
Hadronic collisions

- ▶ Typical proton-proton collision



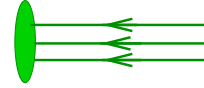
Hadronic collisions

- ▶ Typical proton-proton collision



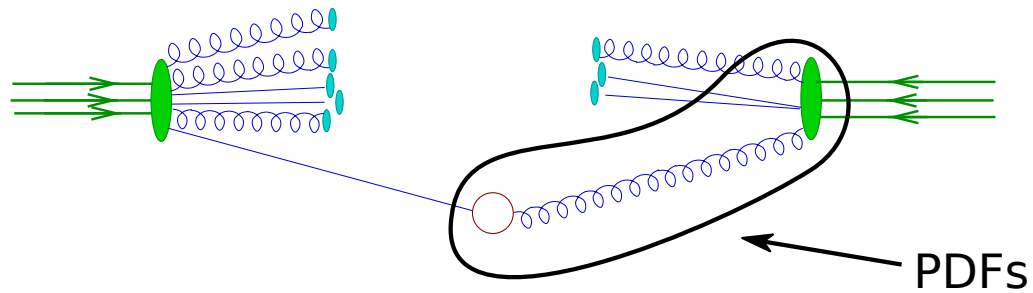
Hadronic collisions

- ▶ Typical proton-proton collision



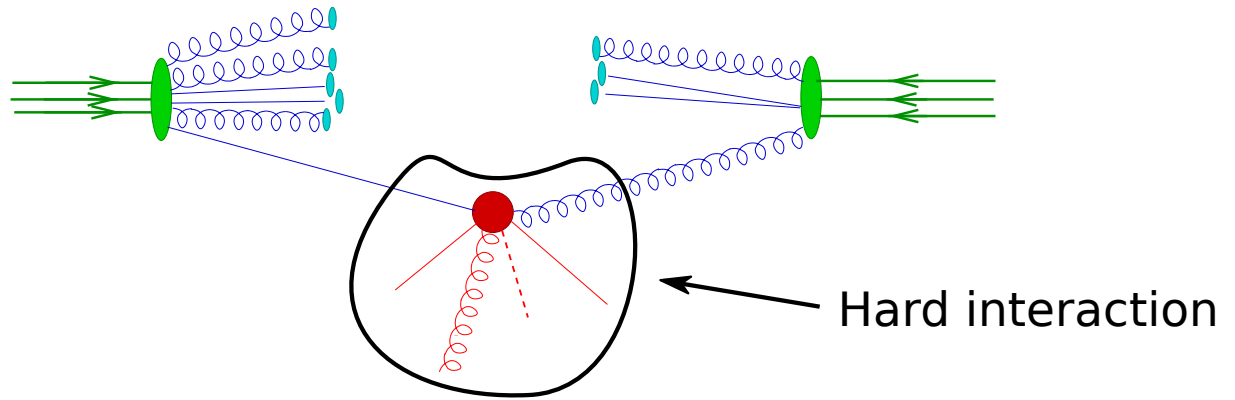
Hadronic collisions

- ▶ Typical proton-proton collision



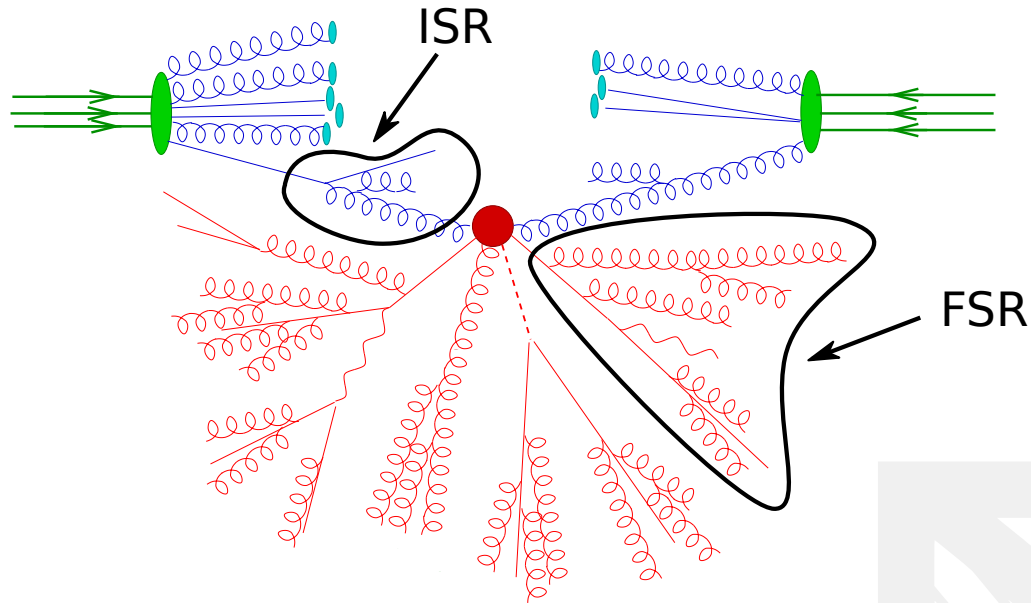
Hadronic collisions

- ▶ Typical proton-proton collision



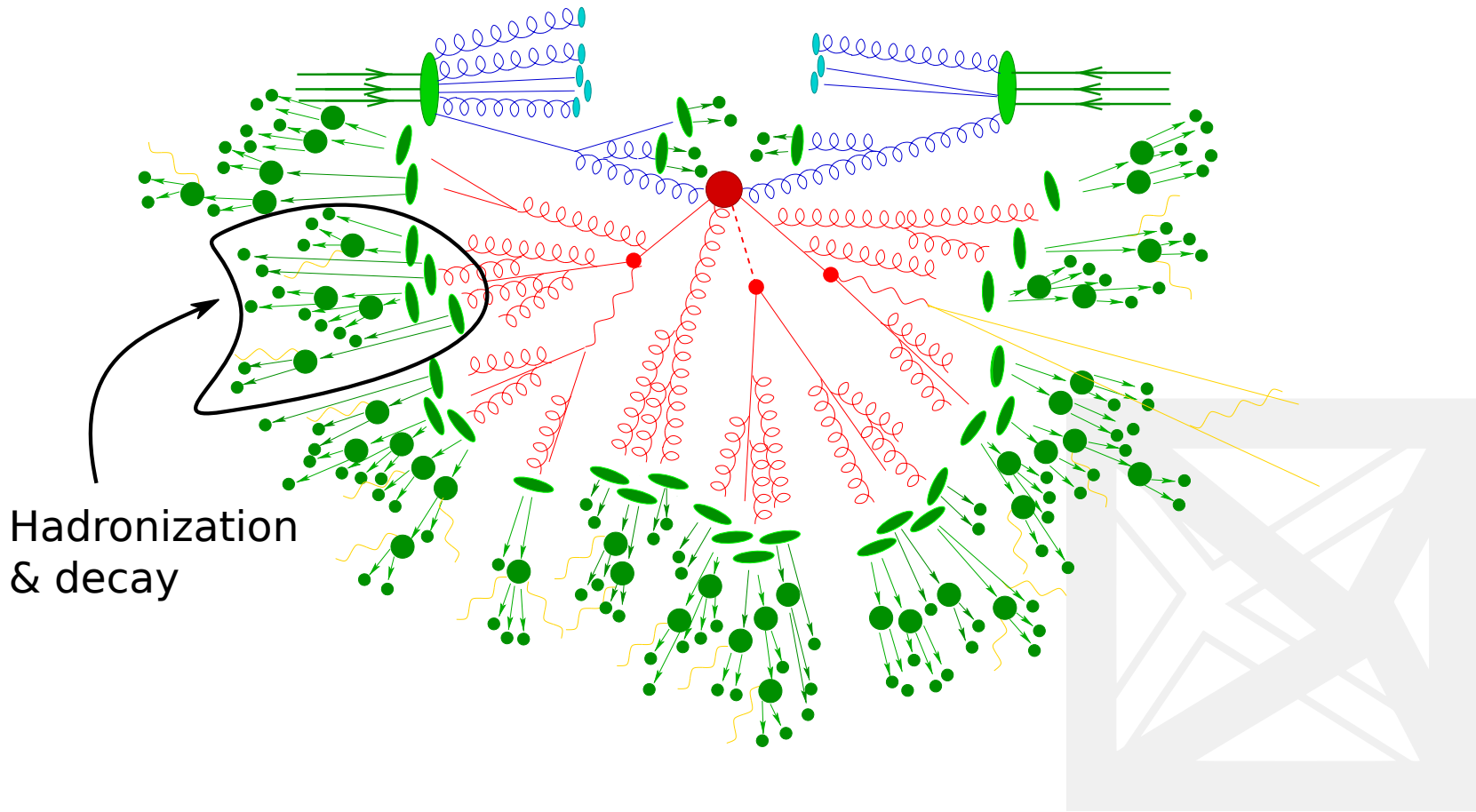
Hadronic collisions

- ▶ Typical proton-proton collision



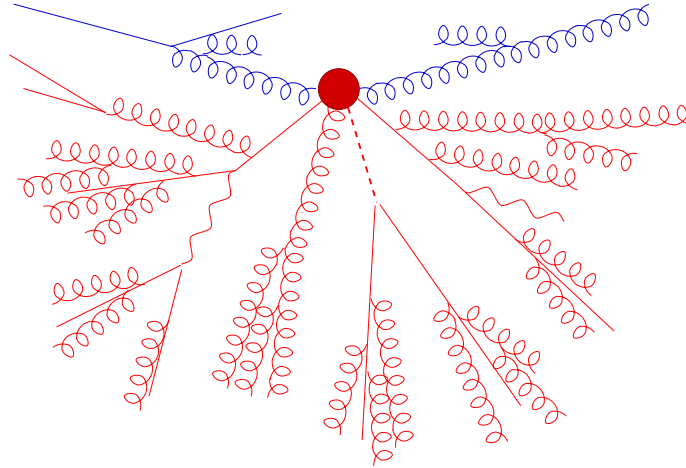
Hadronic collisions

- ▶ Typical proton-proton collision



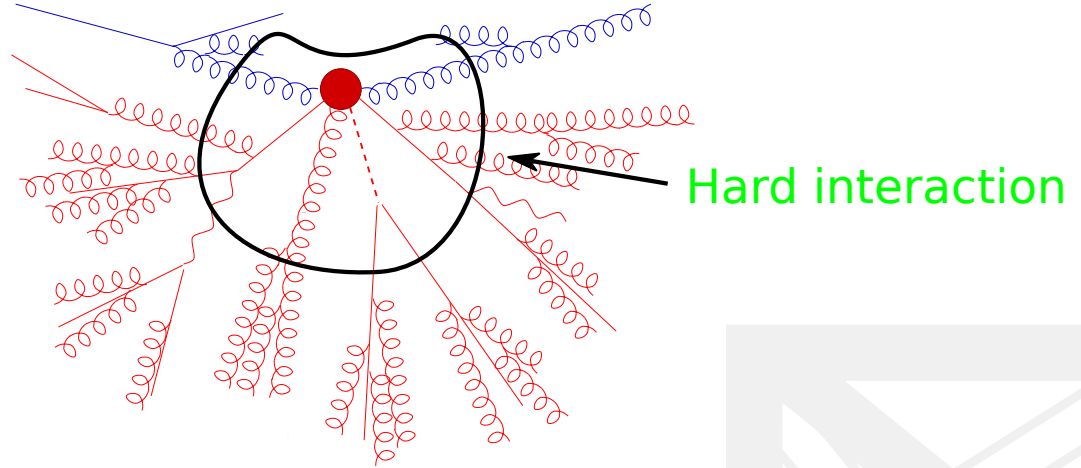
Hadronic collisions

- ▶ Typical proton-proton collision
 - ▷ I focus on interplay of FO calculations and PS @ NLO



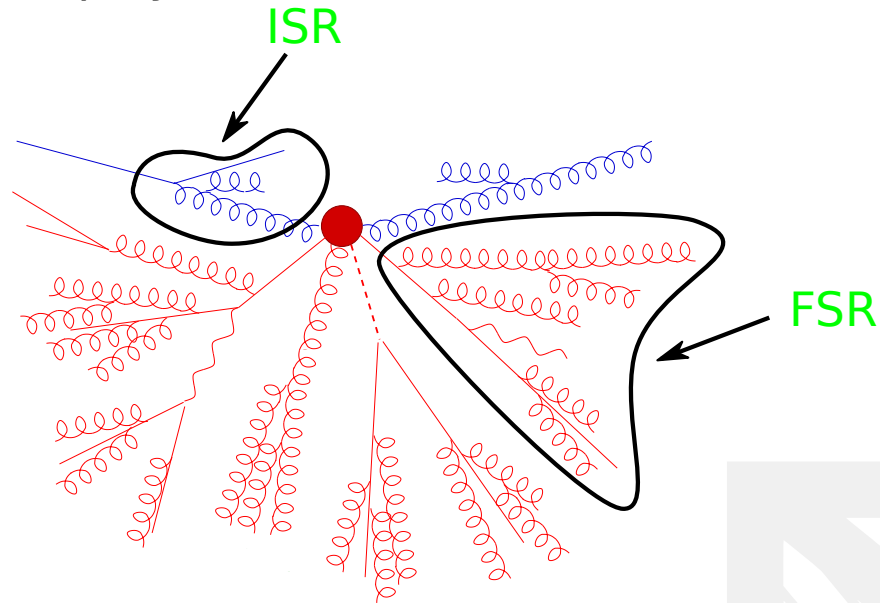
Hadronic collisions

- ▶ Typical proton-proton collision
 - ▷ I focus on interplay of **FO calculations** and PS @ NLO



Hadronic collisions

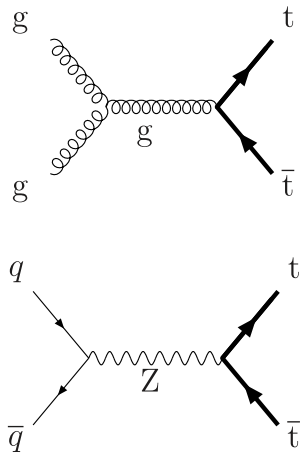
- ▶ Typical proton-proton collision
 - ▷ I focus on interplay of FO calculations and PS @ NLO



Fixed order calculations

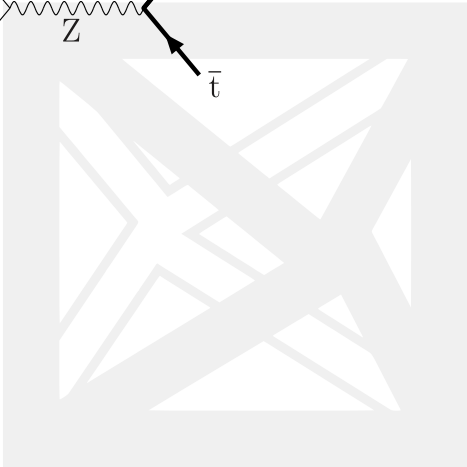
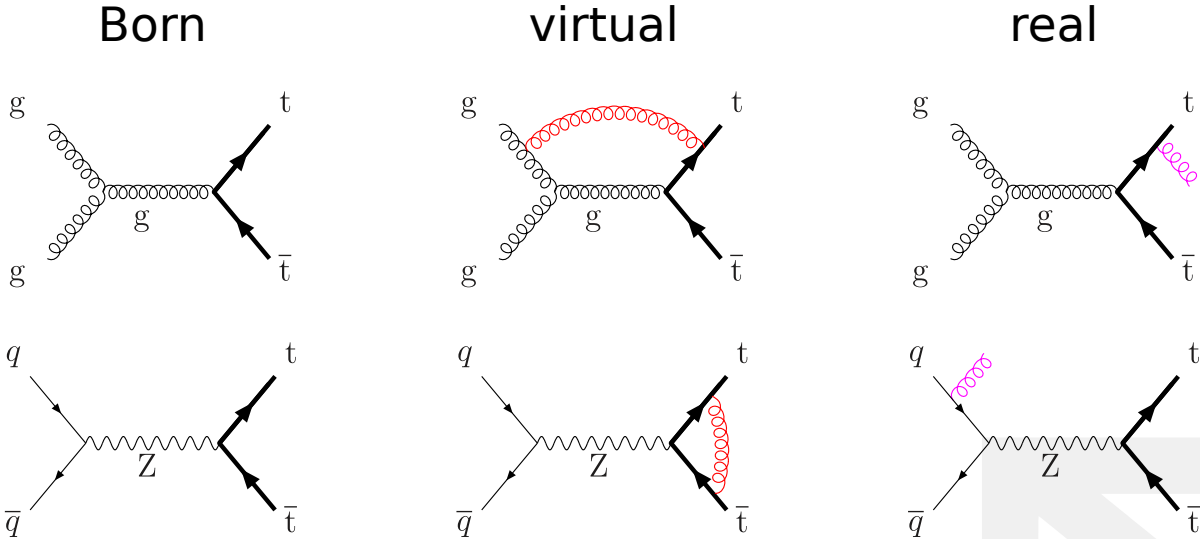
- Take for example top-pair production SM

Born



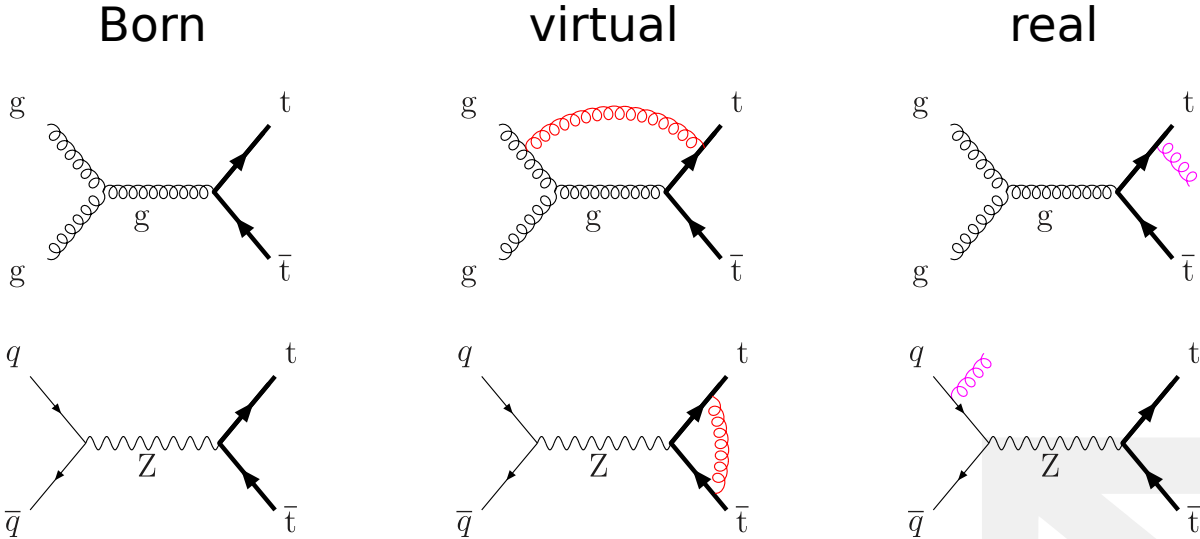
Fixed order calculations

► Take for example top-pair production SM

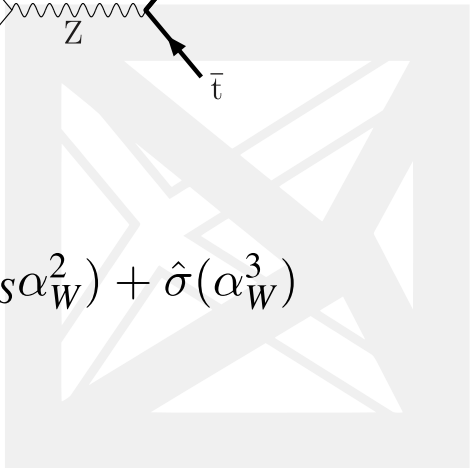


Fixed order calculations

► Take for example top-pair production SM

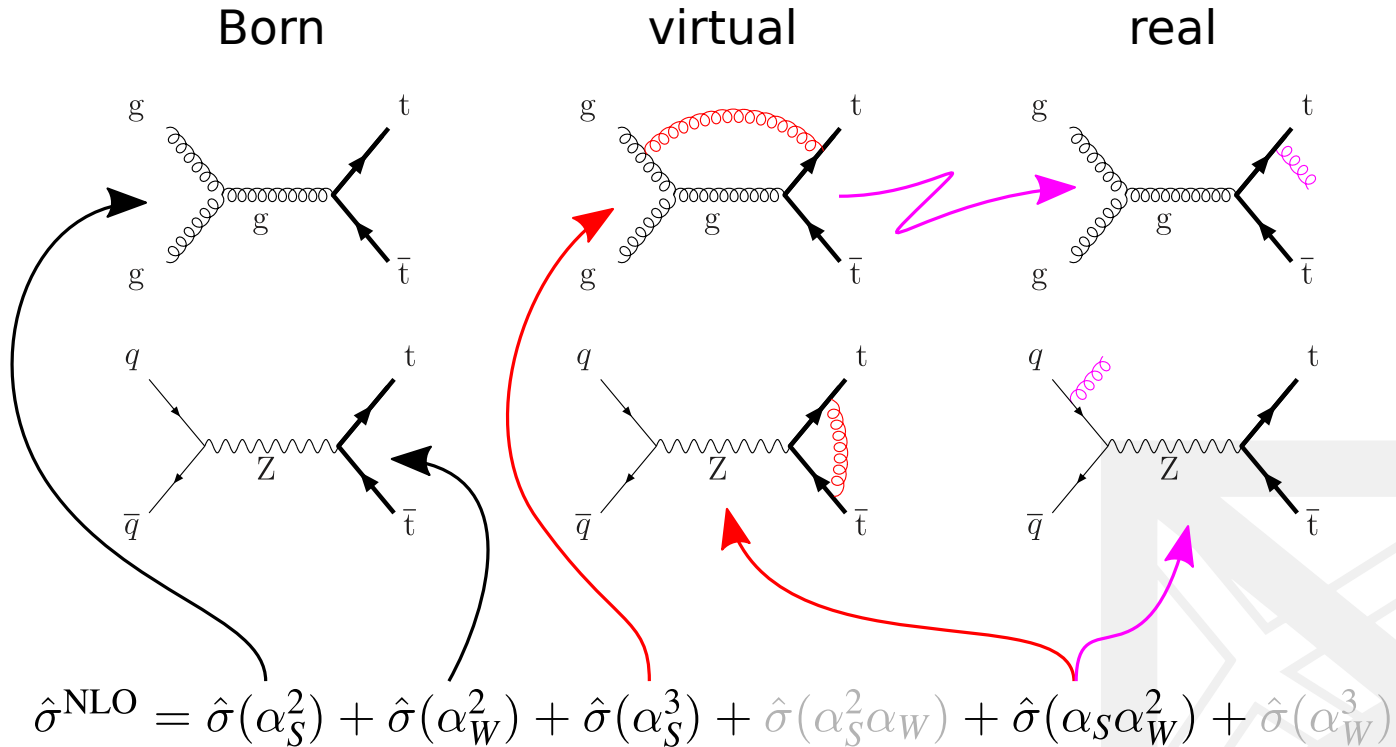


$$\hat{\sigma}^{\text{NLO}} = \hat{\sigma}(\alpha_S^2) + \hat{\sigma}(\alpha_W^2) + \hat{\sigma}(\alpha_S^3) + \hat{\sigma}(\alpha_S^2 \alpha_W) + \hat{\sigma}(\alpha_S \alpha_W^2) + \hat{\sigma}(\alpha_W^3)$$



Fixed order calculations

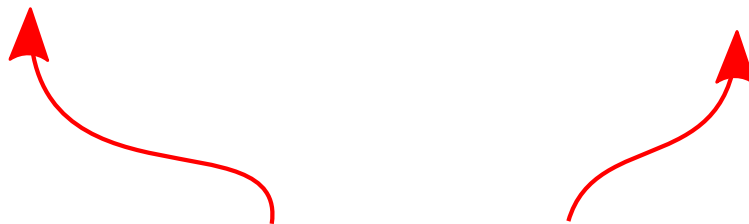
- Take for example top-pair production SM



Fixed order calculations @ NLO

- ▶ Total cross section for a $2 \rightarrow n$ scattering at NLO

$$\sigma_{\text{NLO}} = \int d\Phi_n \left[\mathcal{B}(\Phi_n) + \mathcal{V}(\Phi_n) \right] + \int d\Phi_{n+1} \mathcal{R}(\Phi_{n+1})$$



Separately infinite



Fixed order calculations @ NLO

- ▶ Total cross section for a $2 \rightarrow n$ scattering at NLO

$$\sigma_{\text{NLO}} = \int d\Phi_n \left[\mathcal{B}(\Phi_n) + \mathcal{V}(\Phi_n) \right] + \int d\Phi_{n+1} \mathcal{R}(\Phi_{n+1})$$

- ▶ **UV** divergences renormalized away $\mathcal{V} \rightarrow \mathcal{V}_b$; IR divergences cancel in the sum of real and virtual contributions



Fixed order calculations @ NLO

- ▶ Total cross section for a $2 \rightarrow n$ scattering at NLO

$$\sigma_{\text{NLO}} = \int d\Phi_n \left[\mathcal{B}(\Phi_n) + \mathcal{V}(\Phi_n) \right] + \int d\Phi_{n+1} \mathcal{R}(\Phi_{n+1})$$

- ▶ UV divergences renormalized away $\mathcal{V} \rightarrow \mathcal{V}_b$; IR divergences cancel in the sum of real and virtual contributions



Fixed order calculations @ NLO

- ▶ Total cross section for a $2 \rightarrow n$ scattering at NLO

$$\sigma_{\text{NLO}} = \int d\Phi_n \left[\mathcal{B}(\Phi_n) + \mathcal{V}(\Phi_n) \right] + \int d\Phi_{n+1} \mathcal{R}(\Phi_{n+1})$$

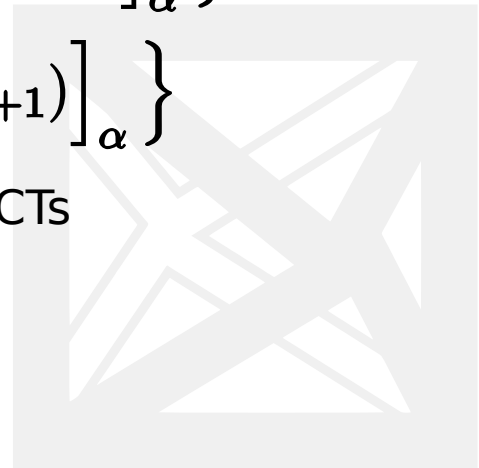
- ▶ UV divergences renormalized away $\mathcal{V} \rightarrow \mathcal{V}_b$; IR divergences cancel in the sum of real and virtual contributions

- ▶ In a general subtraction framework

$$\sigma_{\text{NLO}} = \int d\Phi_n \left\{ \mathcal{B}(\Phi_n) + \mathcal{V}_b(\Phi_n) + \sum_{\alpha} \left[\bar{\mathcal{C}}(\Phi_n) \right]_{\alpha} \right\} \\ + \int d\Phi_{n+1} \left\{ \mathcal{R}(\Phi_{n+1}) - \sum_{\alpha} \left[\mathcal{C}(\Phi_{n+1}) \right]_{\alpha} \right\}$$

- ▶ $\left[\mathcal{C}(\Phi_{n+1}) \right]_{\alpha}$: real CTs; $\left[\bar{\mathcal{C}}(\Phi_n) \right]_{\alpha}$: integrated CTs

- ▶ α labels singular regions



Fixed order calculations @ NLO

- ▶ Total cross section for a $2 \rightarrow n$ scattering at NLO

$$\sigma_{\text{NLO}} = \int d\Phi_n \left[\mathcal{B}(\Phi_n) + \mathcal{V}(\Phi_n) \right] + \int d\Phi_{n+1} \mathcal{R}(\Phi_{n+1})$$

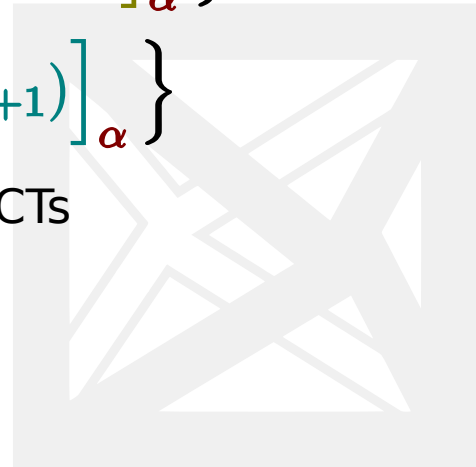
- ▶ UV divergences renormalized away $\mathcal{V} \rightarrow \mathcal{V}_b$; IR divergences cancel in the sum of real and virtual contributions

- ▶ In a general subtraction framework

$$\sigma_{\text{NLO}} = \int d\Phi_n \left\{ \mathcal{B}(\Phi_n) + \mathcal{V}_b(\Phi_n) + \sum_{\alpha} \left[\bar{\mathcal{C}}(\Phi_n) \right]_{\alpha} \right\} \\ + \int d\Phi_{n+1} \left\{ \mathcal{R}(\Phi_{n+1}) - \sum_{\alpha} \left[\mathcal{C}(\Phi_{n+1}) \right]_{\alpha} \right\}$$

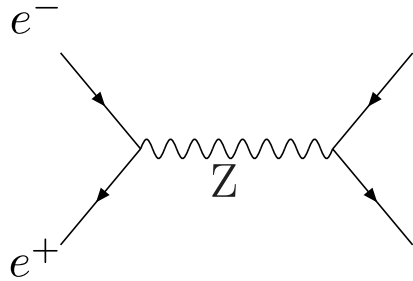
- ▶ $\left[\mathcal{C}(\Phi_{n+1}) \right]_{\alpha}$: real CTs; $\left[\bar{\mathcal{C}}(\Phi_n) \right]_{\alpha}$: integrated CTs

- ▶ α labels singular regions



Parton shower

- ▶ Parton showers can be automated

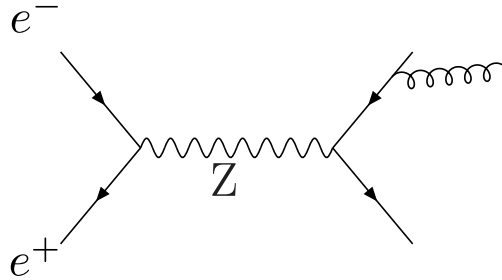


$$W = W_B$$



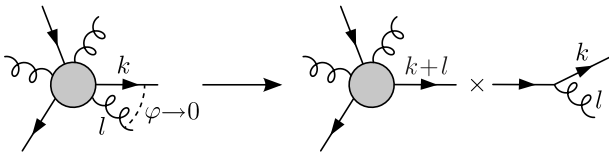
Parton shower

- ▶ Parton showers can be automated

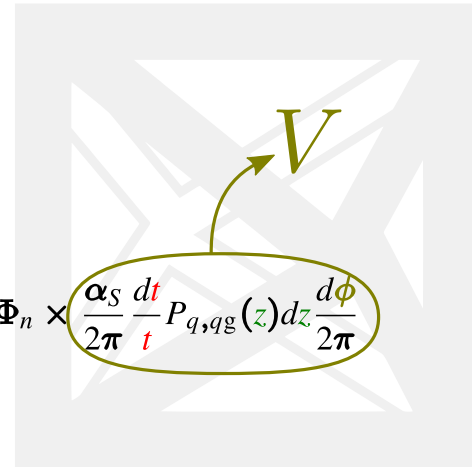


$$W = W_B \times V$$

- ▶ Real corrections in collinear approximation:

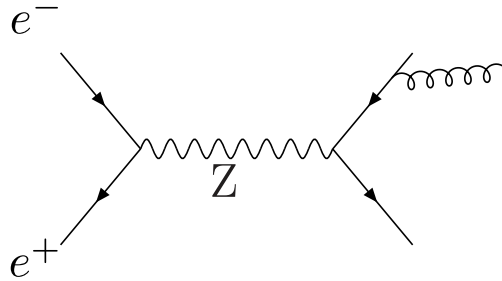


$$|\mathcal{M}_{n+1}|^2 d\Phi_{n+1} \rightarrow |\mathcal{M}_n|^2 d\Phi_n \times \frac{\alpha_S}{2\pi} \frac{dt}{t} P_{q,qg}(z) dz \frac{d\phi}{2\pi}$$



Parton shower

- ▶ Parton showers can be automated



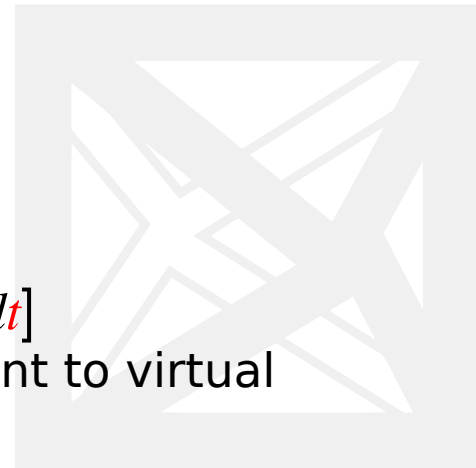
$$W = W_B \times V \times \Delta$$

- ▶ Virtual corrections in collinear approximation:



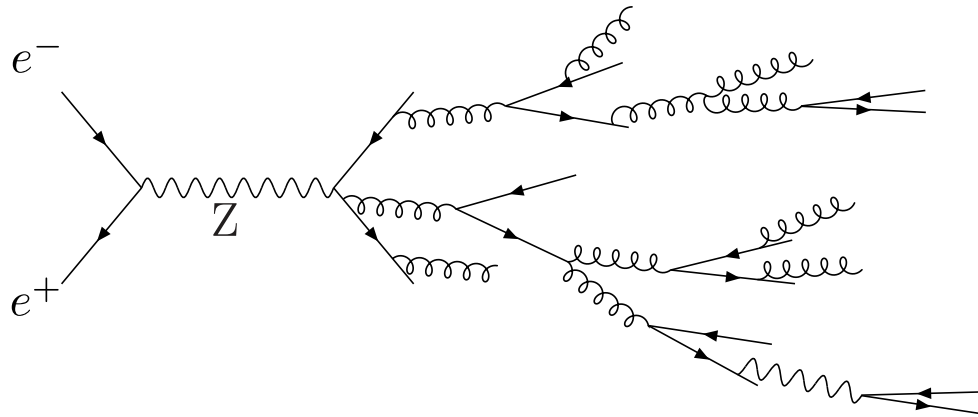
$$dP(t, t + dt) = \frac{\alpha_S}{2\pi} \frac{dt}{t} \int \frac{d\phi}{2\pi} \int P_{i,jl}(z) dz$$

- ▶ dP probability of $i \rightarrow jl$ splitting in $[t, t + dt]$
- ▶ $1 - dP$ probability of no radiation equivalent to virtual contribution



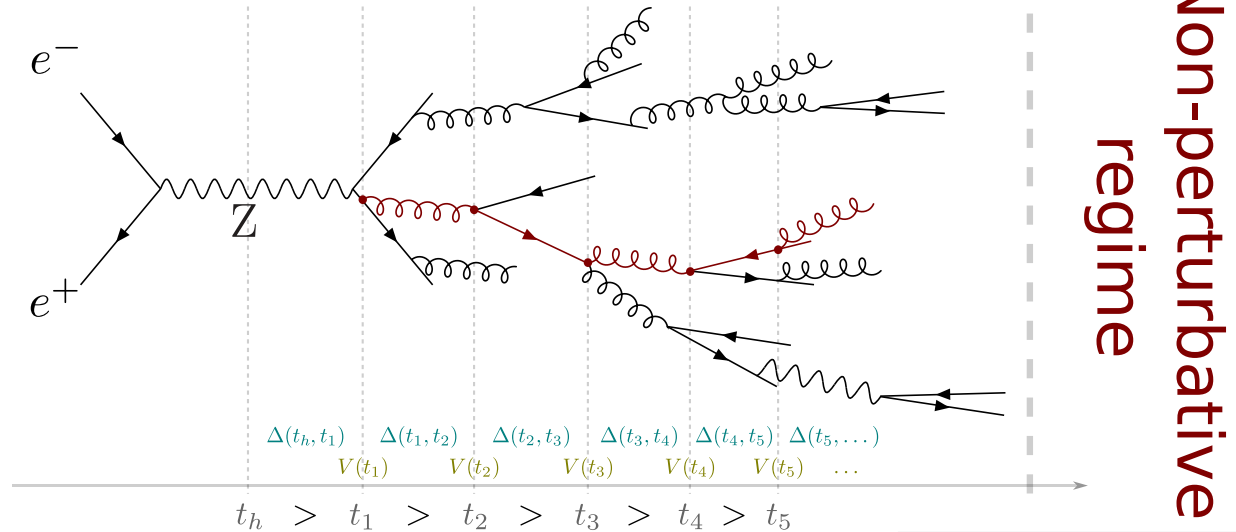
Parton shower

- ▶ Parton showers can be automated



Parton shower

- ▶ Parton showers can be automated



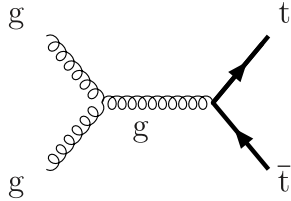
- ▶ Variable t measures hardness
 - ▷ vanishes in the collinear limit
- ▶ Weight of the event is the Born weight times the splitting and Sudakov factors



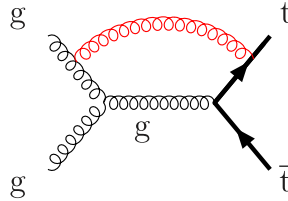
NLO & PS

- Fixed order calculation @ Next-to-Leading Order (NLO)

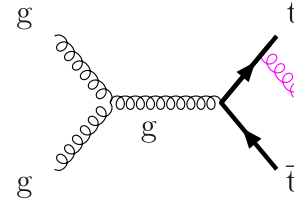
Born



virtual



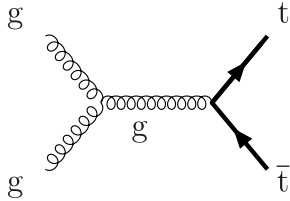
real



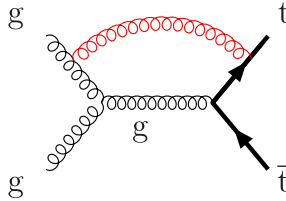
NLO & PS

- Fixed order calculation @ Next-to-Leading Order (NLO)

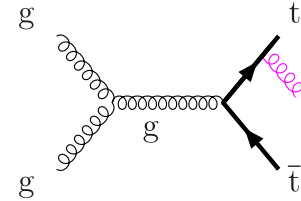
Born



virtual

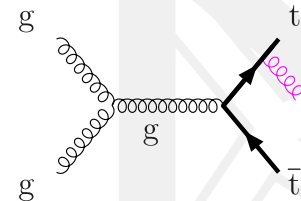
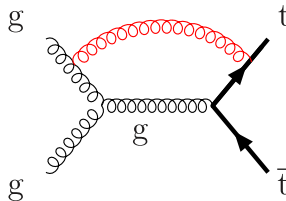
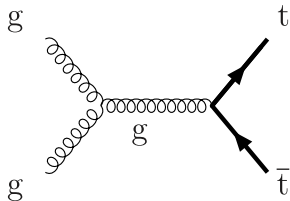


real



- Parton shower (PS)

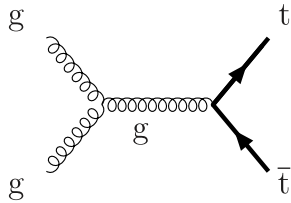
Born



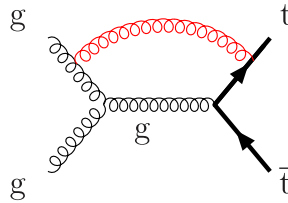
NLO & PS

- Fixed order calculation @ Next-to-Leading Order (NLO)

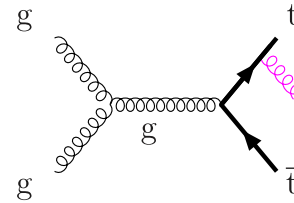
Born



virtual



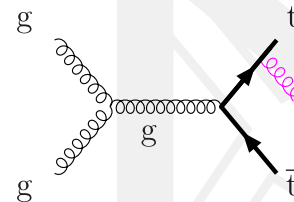
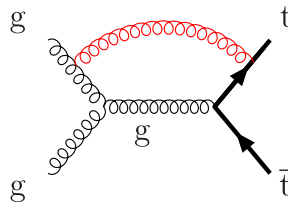
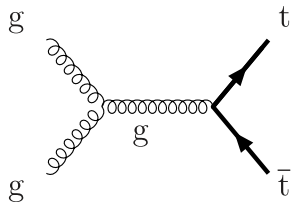
real



- Parton shower (PS)

aproximation

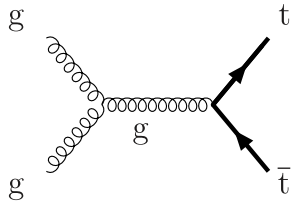
Born



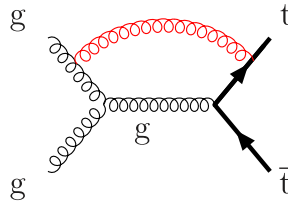
NLO+PS

► NLO merged with PS

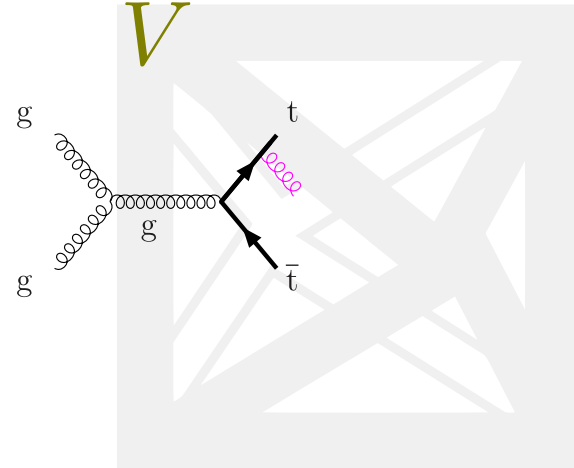
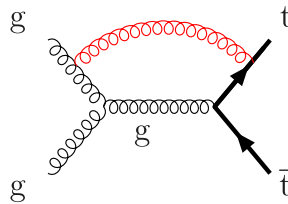
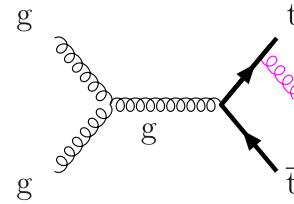
Born



virtual



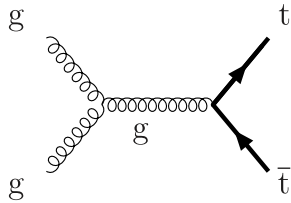
real



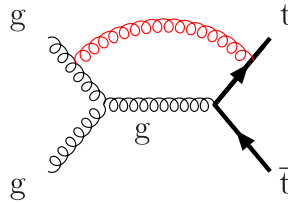
NLO+PS

► NLO merged with PS

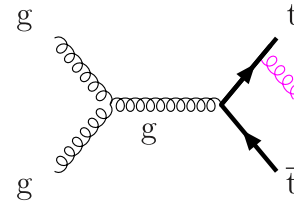
Born



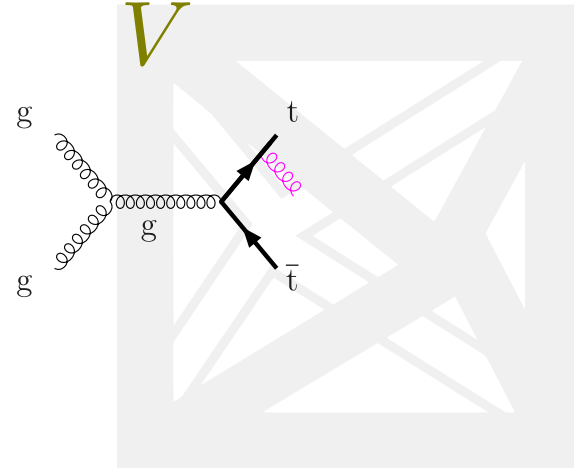
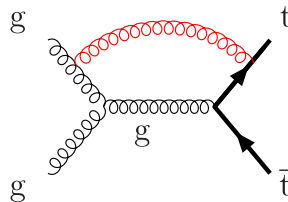
virtual



real

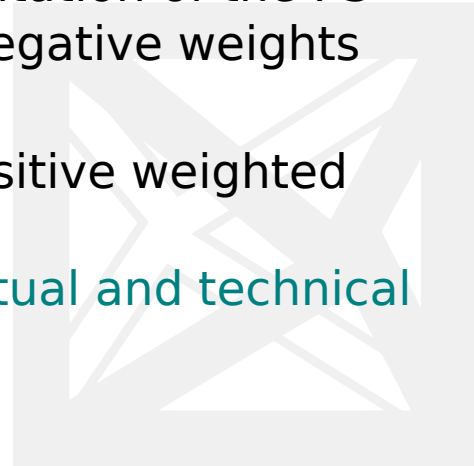


overcounted



NLO+PS

- ▶ Naive matching of NLO and PS doesn't work: **overcounting**
 - ▷ Both PS and NLO contain the real and virtual contributions in the collinear limit
- ▶ **Overcounting** can be **solved** for example by modifying the Sudakov form factor for the first radiation
- ▶ Multiple **solutions** exist
 - ▷ Matrix Element (ME) corrections, `Pythia`
 - available only for a few processes
 - ▷ MC@NLO: `mg5_aMC`, ...
 - procedure depends on the implementation of the PS algorithm, can lead to events with negative weights
 - ▷ POWHEG: `POWHEG BOX`, ...
 - independent of the PS algorithm, positive weighted events
- ▶ **However some problems remain, both conceptual and technical**



NLO+PS in POWHEG BOX

- ▶ POWHEG BOX automatically calculates everything down to the generation of the **hardest emission** provided the user specifies
 - ▷ Born matrix elements $\mathcal{B}(\Phi_n)$
 - ▷ Renormalized virtual matrix elements $\mathcal{V}_b(\Phi_n)$
 - ▷ Real matrix elements $\mathcal{R}(\Phi_{n+1})$
- ▶ **Consecutive emissions** can be generated by usual tools implementing PS (Pythia, Herwig, ...)

- ▶ It uses the FKS subtraction method

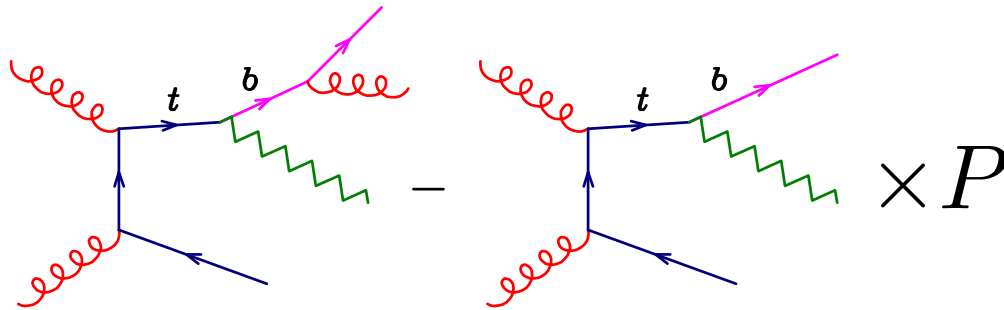
$$\sigma_{\text{NLO}} = \int d\Phi_n \left\{ \mathcal{B}(\Phi_n) + \hat{\mathcal{V}}(\Phi_n) \right\} + \int d\Phi_{n+1} \hat{\mathcal{R}}(\Phi_{n+1})$$

$$\hat{\mathcal{R}} \equiv \frac{1}{\xi} \left\{ \left(\frac{1}{\xi} \right)_+ \left(\frac{1}{1-y} \right)_+ [\xi^2 (1-y) \mathcal{R}] \right\} \quad \hat{\mathcal{V}} = \frac{\alpha_s}{2\pi} \left(Q\mathcal{B} + \sum_{\substack{i,j \in \mathcal{I} \\ i \neq j}} \mathcal{I}_{ij} \mathcal{B}_{ij} + \mathcal{V}_{\text{fin}} \right)$$

Treatment of resonances

in collaboration with P. Nason; arXiv:1509.09071

- ▶ Counterterm kinematics does not preserve the mass of the resonance - **spoiling IR cancelation**



- ▶ Mapping from real kinematics into underlying born kinematics does not preserve the mass of the resonance - leading to **distortion of radiation observables**

$$\Delta(p_T^2) = \exp \left[- \int \frac{R(\Phi_B, \Phi_{\text{rad}})}{B(\Phi_B)} \theta(k_T(\Phi_{\text{rad}}) - p_T) d\Phi_{\text{rad}} \right]$$

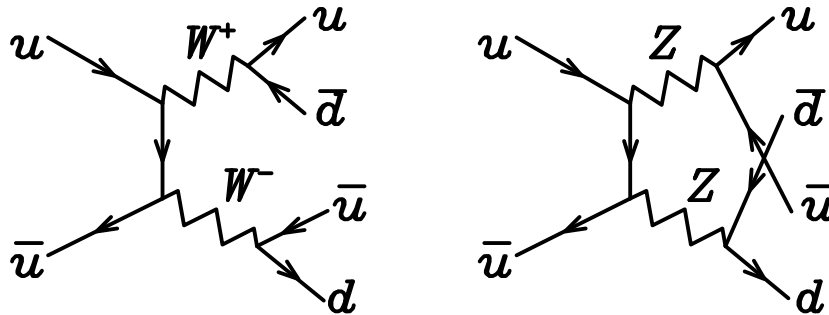
- ▷ R/B large violating the collinear approximation in region

$$m^2 \ll \Gamma E$$

Treatment of resonances

in collaboration with P. Nason; arXiv:1509.09071

- Mapping from the real to born kinematics has to **preserve the masses of the resonances**



- All contributions needs to be split up into regions with only **one dominant resonance structure**

$$\mathcal{B}_1 = \frac{P^1 \mathcal{B}}{P^1 + P^2}$$

$$P^1 = \frac{M_W^4}{(s_{34} - M_W^2)^2 + \Gamma_W^2 M_W^2} \times \frac{M_W^4}{(s_{56} - M_W^2)^2 + \Gamma_W^2 M_W^2}$$

$$P^2 = \frac{M_Z^4}{(s_{35} - M_Z^2)^2 + \Gamma_Z^2 M_Z^2} \times \frac{M_Z^4}{(s_{46} - M_Z^2)^2 + \Gamma_Z^2 M_Z^2}$$

Treatment of resonances

in collaboration with P. Nason; arXiv:1509.09071

► Implement:

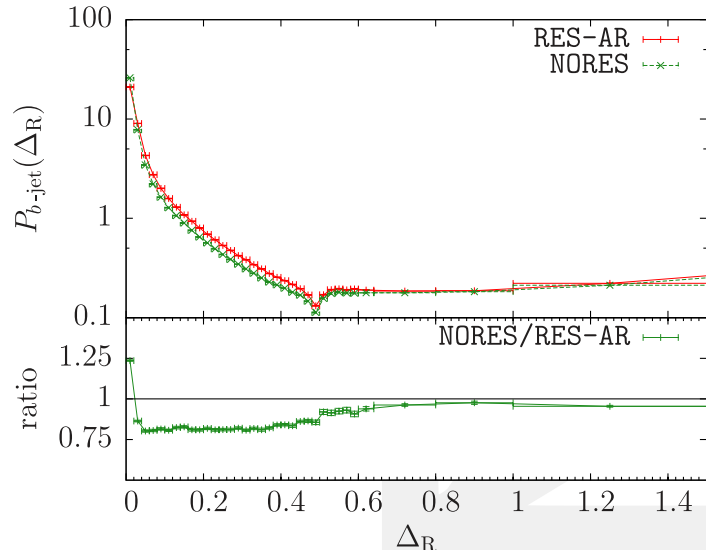
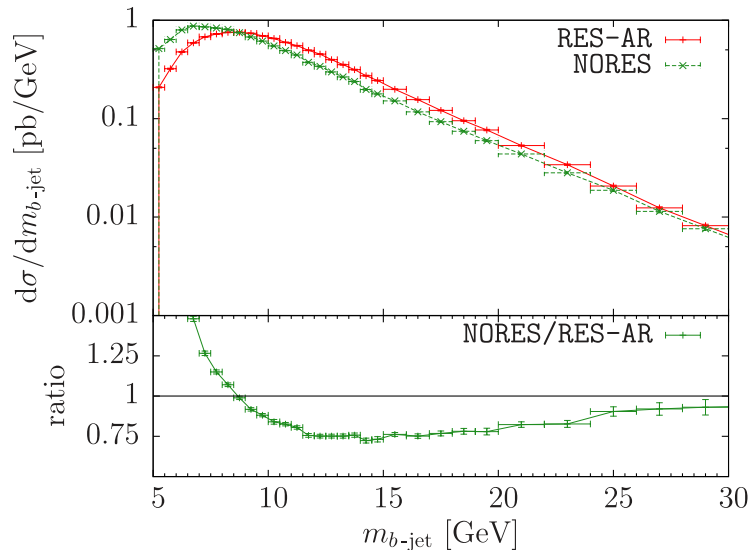
▷ $pp \rightarrow \mu^+ \nu_\mu j b j$ dominated by t -channel single top production at NLO QCD

▷ born and real: MadGraph4, virtual: MG5_aMC@NLO

► Study impact of proper resonance treatment:

▷ **NORES**: resonant treatment off

▷ **RES-AR**: resonant treatment on, 1 hardest emission from resonance + 1 hardest emission from the production



Treatment of resonances

in collaboration with P. Nason; arXiv:1509.09071

► Implement:

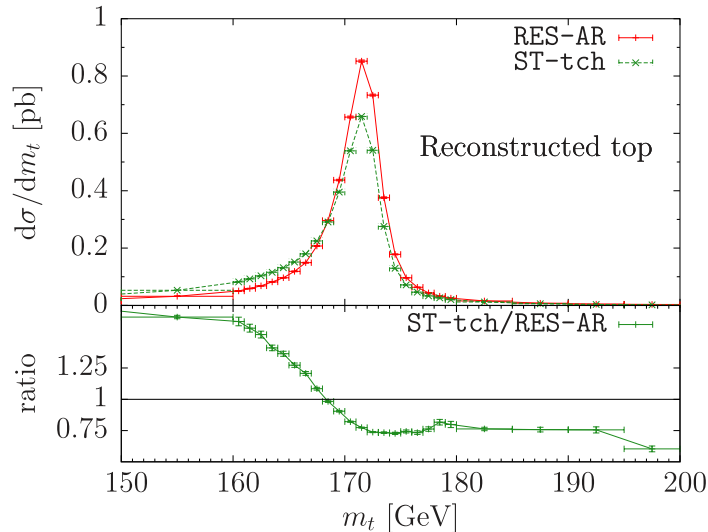
▷ $pp \rightarrow \mu^+ \nu_\mu j b j$ dominated by t -channel single top production at NLO QCD

▷ born and real: MadGraph4, virtual: MG5_aMC@NLO

► Study impact of proper resonance treatment:

▷ **ST-tch**: top-pair@NLO, top decay@LO

▷ **RES-AR**: resonant treatment on, 1 hardest emission from resonance + 1 hardest emission from the production



► average top mass in

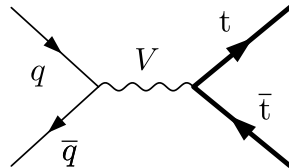
$$m_t = 172.5 \pm 15 \text{ GeV}$$

▷ **ST-tch**: $M_{\text{trec}} = 169.59(1) \text{ GeV}$

▷ **RES-AR**: $M_{\text{trec}} = 170.55(2) \text{ GeV}$

Electroweak top-pair @ NLO BSM

in collaboration with R. Bonciani, M. Klasen,
F. Lyonnet, I. Schienbein; arXiv:1511.xxxxx



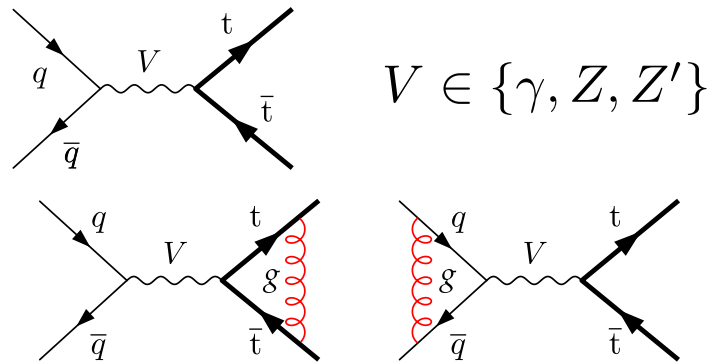
$$V \in \{\gamma, Z, Z'\}$$

- ▶ Z' couplings family diagonal but generic otherwise
 - ▷ suitable in particular for simulating models with extended gauge group, for example G(221) models
- ▶ amplitudes calculated analytically:
 - ▷ QGRAPH/DIANA + FORM tool chain



Electroweak top-pair @ NLO BSM

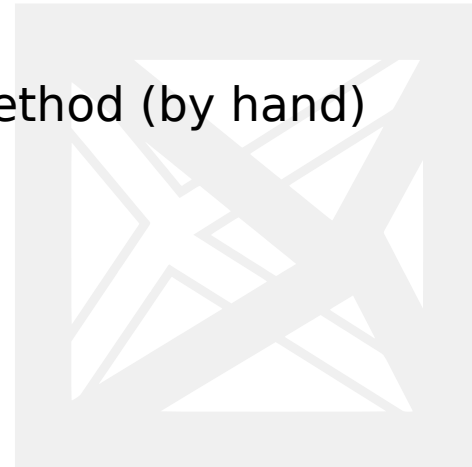
in collaboration with R. Bonciani, M. Klasen,
F. Lyonnet, I. Schienbein; arXiv:1511.xxxxx



► **virtual** corrections:

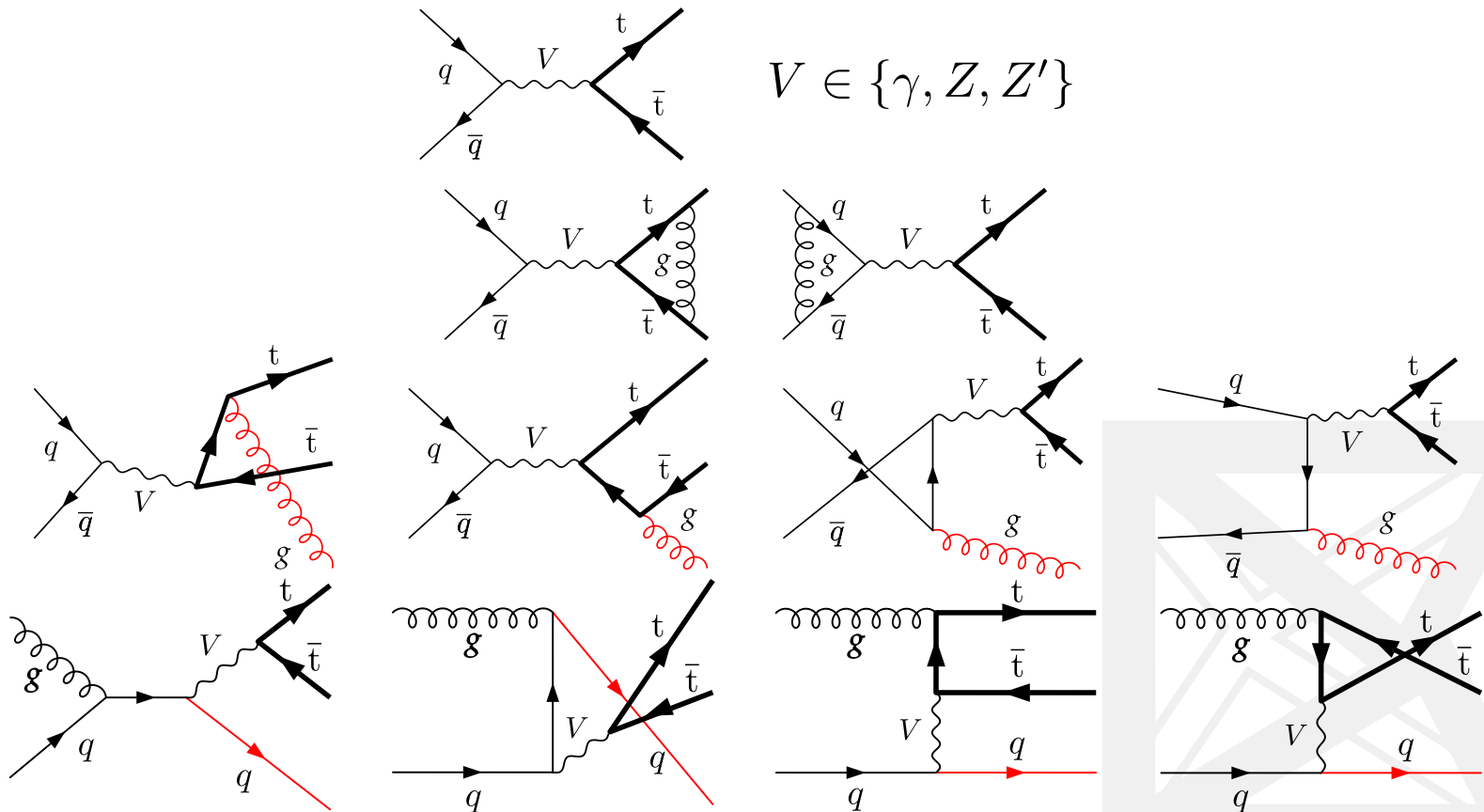
▷ loop integrals: IPBs (REDUZE)

▷ master integrals: differential equation method (by hand)



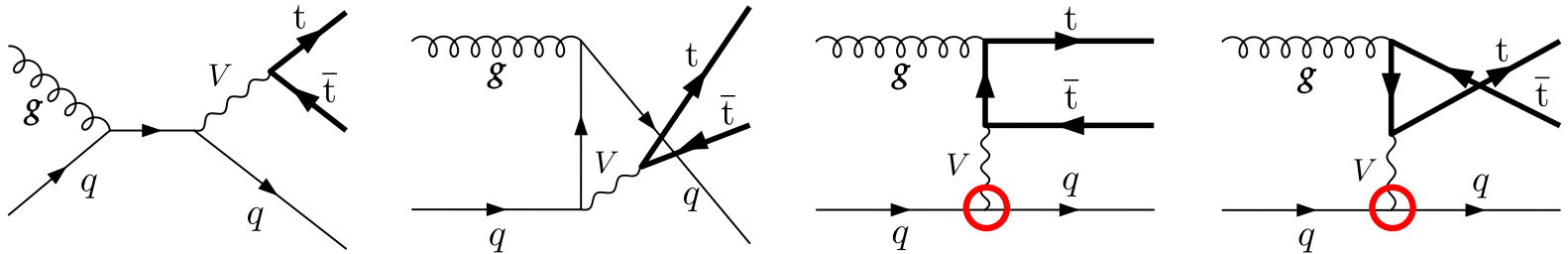
Electroweak top-pair @ NLO BSM

in collaboration with R. Bonciani, M. Klasen,
F. Lyonnet, I. Schienbein; arXiv:1511.xxxxx



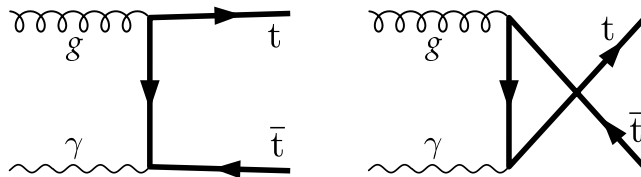
Electroweak top-pair @ NLO BSM

in collaboration with R. Bonciani, M. Klasen,
F. Lyonnet, I. Schienbein; arXiv:1511.xxxxx



► QED singularity:

- POWHEG BOX V2 should be ready to deal with
- we had to modify the routine searching for singular regions



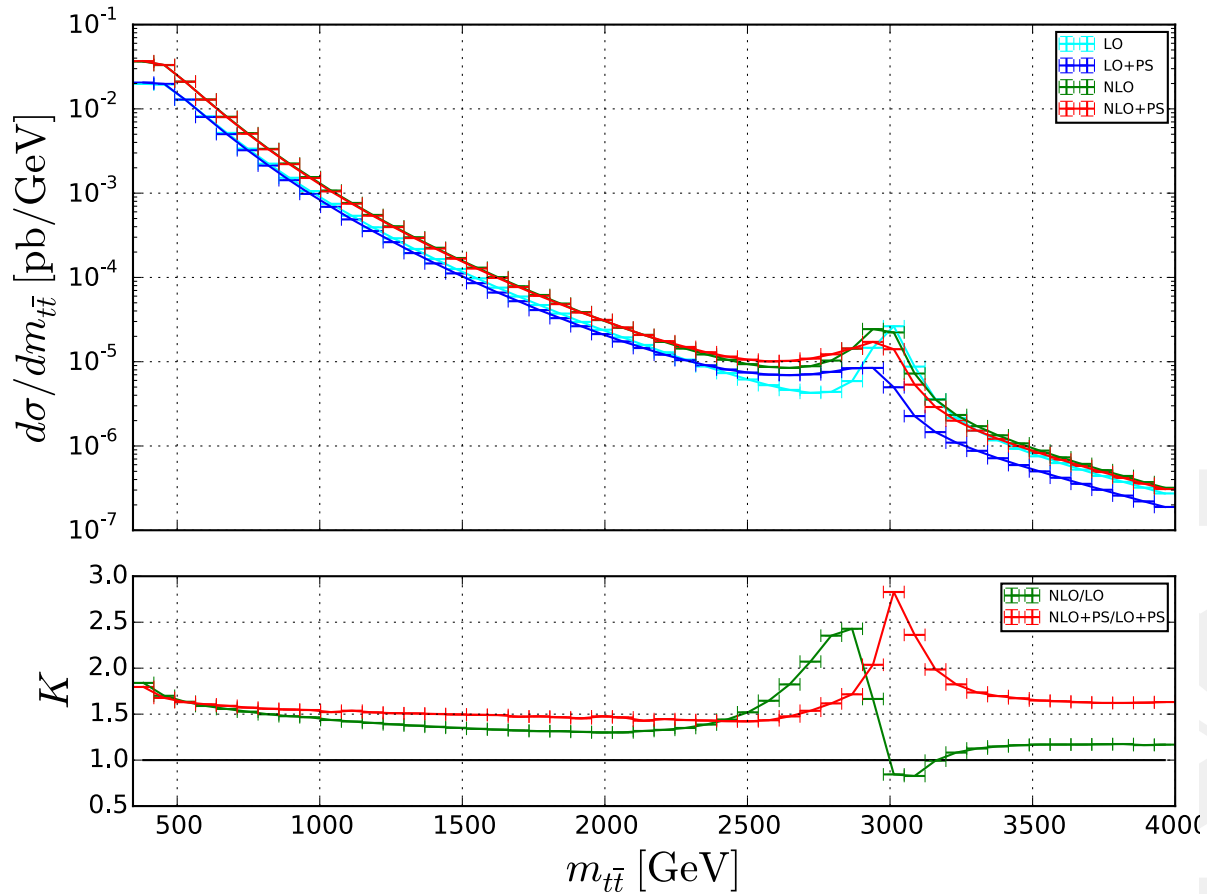
Electroweak top-pair @ NLO BSM

in collaboration with R. Bonciani, M. Klasen,
F. Lyonnet, I. Schienbein; arXiv:1511.xxxxx

Order	Processes	Model	σ [pb]	σ [pb] ($m_{t\bar{t}} > \frac{3}{4}m_{Z'}$)
LO	$q\bar{q}/gg \rightarrow t\bar{t}$		473.93(7)	0.15202(2)
NLO	$q\bar{q}/gg + qg \rightarrow t\bar{t} + q$		1261.0(2)	0.45255(7)
LO	$\gamma g + g\gamma \rightarrow t\bar{t}$		4.8701(8)	0.0049727(6)
LO	$\gamma g + g\gamma \rightarrow t\bar{t}$ (NLO α_s and PDFs)		5.1891(8)	0.004661(6)
LO	$q\bar{q} \rightarrow \gamma/Z \rightarrow t\bar{t}$	SM	0.36620(7)	0.00017135(3)
NLO	$q\bar{q} \rightarrow \gamma/Z \rightarrow t\bar{t}$	SM	0.5794(1)	0.00017174(5)
NLO	$q\bar{q}/qg \rightarrow \gamma/Z + q \rightarrow t\bar{t} + q$	SM	4.176(2)	0.001250(6)
LO	$q\bar{q} \rightarrow Z' \rightarrow t\bar{t}$	SSM	0.0050385(8)	0.0044848(7)
LO	$q\bar{q} \rightarrow \gamma/Z/Z' \rightarrow t\bar{t}$	SSM	0.35892(7)	0.0043464(7)
NLO	$q\bar{q} \rightarrow \gamma/Z/Z' \rightarrow t\bar{t}$	SSM	0.5676(1)	0.005155(3)
NLO	$q\bar{q}/qg \rightarrow \gamma/Z + q \rightarrow t\bar{t} + q$	SSM	4.172(2)	0.007456(9)

Electroweak top-pair @ NLO BSM

in collaboration with R. Bonciani, M. Klasen,
F. Lyonnet, I. Schienbein; arXiv:1511.xxxxx

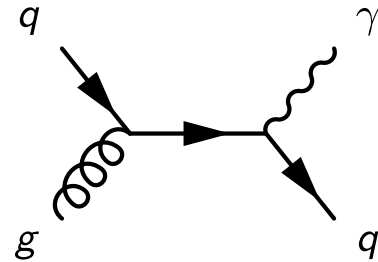
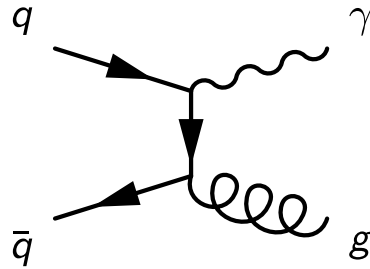


Photonproduction in POWHEG BOX

in collaboration with M. Klasen, K. Kovarik, F. König

► direct contribution

LO
 $\mathcal{O}(\alpha\alpha_s)$

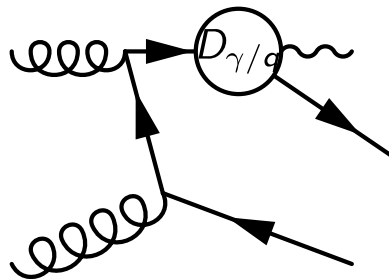


NLO
 $\mathcal{O}(\alpha\alpha_s^2)$

+ extra g or $q\bar{q}$ splitting

+ virtual corrections

► fragmentation contribution



Photonproduction in POWHEG BOX

in collaboration with M. Klasen, K. Kovarik, F. König

- ▶ At NLO γ radiation off of massless quarks

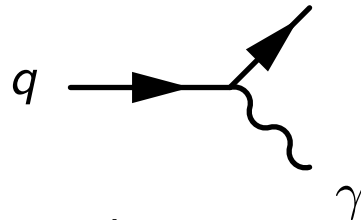
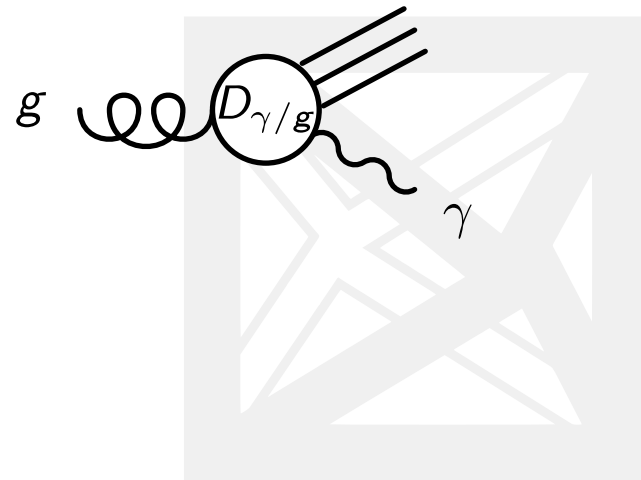
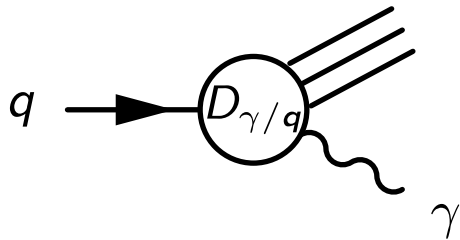


exhibit collinear singularities that cannot be cancelled by loops

- ▶ these have to be absorbed into γ fragmentation functions

$$D_{\gamma/k}(z, \mu_D)$$



due to γ coupling to QCD bound states

Photonproduction in POWHEG BOX

in collaboration with M. Klasen, K. Kovarik, F. König

- ▶ At NLO γ radiation off of massless quarks

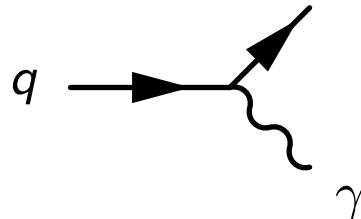
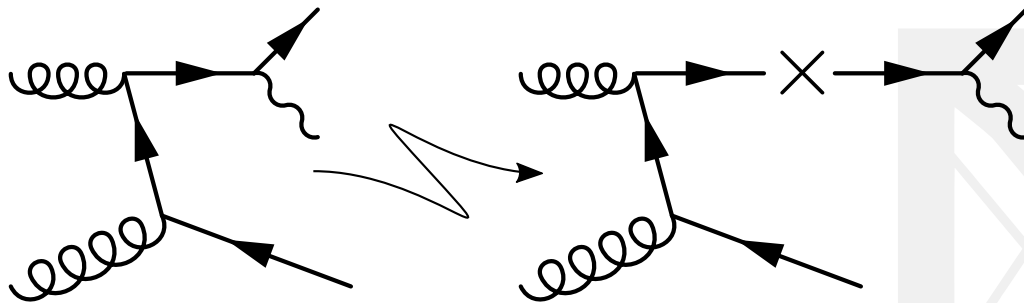


exhibit collinear singularities that cannot be cancelled by loops

- ▶ collinear singularity can also be treated by real counterterms generated by FKS

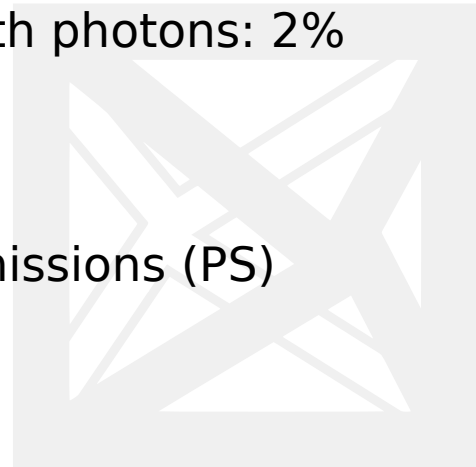


- ▶ perhaps rest of the fragmentation contribution can be recovered by parton shower

Photonproduction in POWHEG BOX

in collaboration with M. Klasen, K. Kovarik, F. König

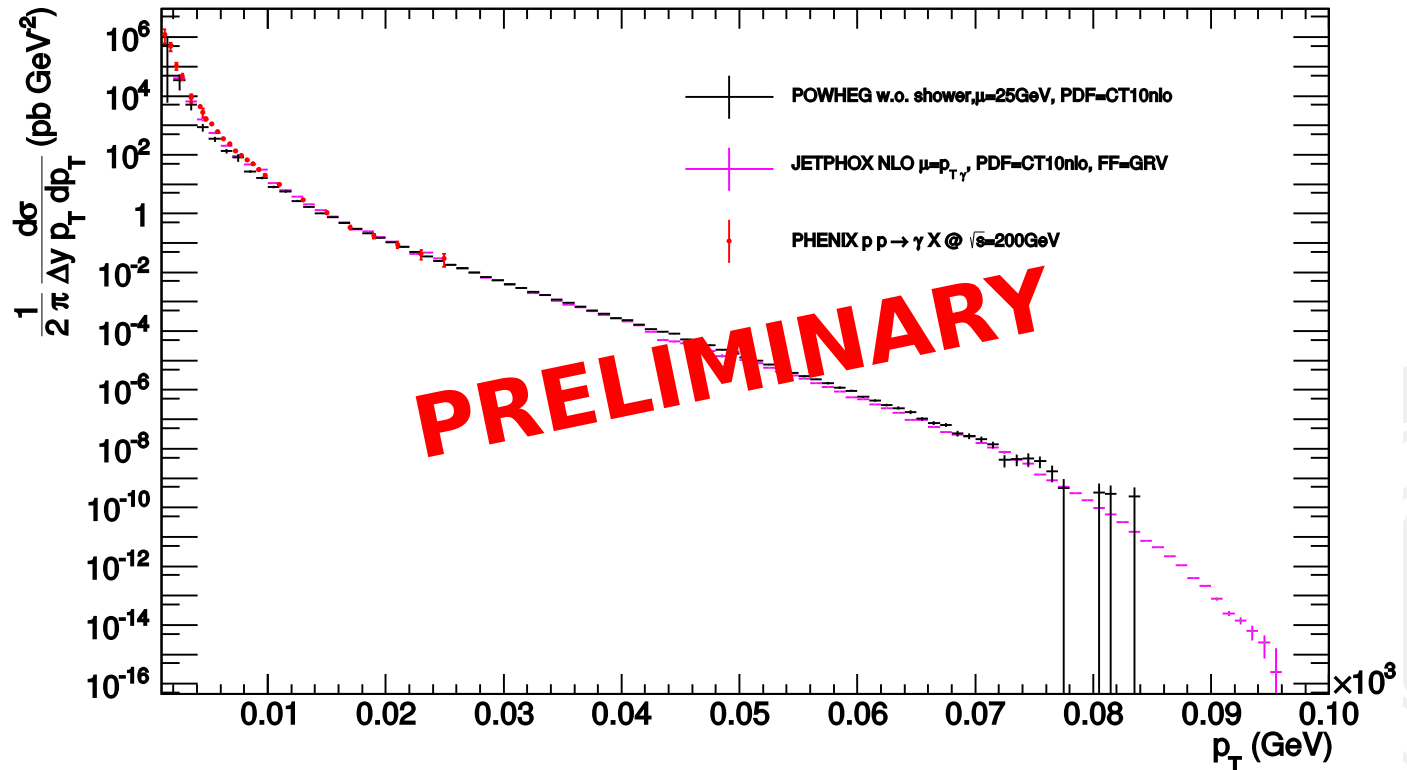
- ▶ Calculated:
 - ▷ Born, colour and spin-correlated Born $\mathcal{O}(\alpha_{em}\alpha_s)$
 - ▷ virtual and real $\mathcal{O}(\alpha_{em}\alpha_s^2)$
 - ▷ Born $\mathcal{O}(\alpha_s^2)$
- ▶ Extended POWHEG BOX
 - ▷ photon radiation off massless particles
 - ▷ reweighting with modified R/B ratio
low efficiency in generation of events with photons: 2%
with $\alpha_{em} \rightarrow 40\alpha_{em}$ the efficiency is: 50%
- ▶ To do:
 - ▷ consistent generation of consecutive emissions (PS)



Photonproduction in POWHEG BOX

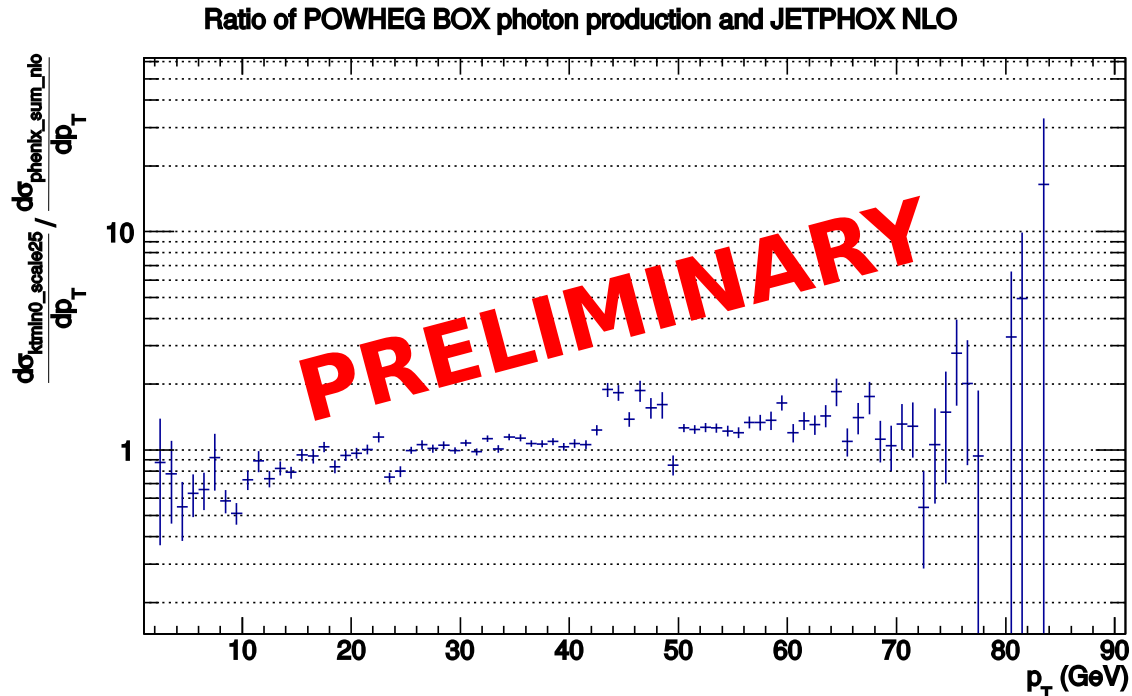
in collaboration with M. Klasen, K. Kovarik, F. König

p_T distribution of hardest photon



Photonproduction in POWHEG BOX

in collaboration with M. Klasen, K. Kovarik, F. König



Note that for JETPHOX $\mu = \mu_R = \mu_F = p_{T\gamma}$ and POWHEG $\mu = 25$ GeV (no low- p_T -cut)

Summary and Outlook

- ▶ I have reviewed NLO+PS matching
- ▶ NLO+PS is a very important and exciting research topic
- ▶ POWHEG BOX is leading tool for NLO+PS matching, developed on many fronts
- ▶ Research group of M. Klasen actively participates in this effort

- ▶ Treatment of resonances:
 - ▷ important for the shape of radiation or derived observables, most notably may affect the measurement of top mass
 - ▷ top-pair production, HV production still to come in 2015

- ▶ New processes:
 - ▷ EW top-pair production, single top production beyond SM
 - ▷ photonproduction

