Past, current and future projects

GRK Annual Retreat 2015

Kevin Eckert

24. November 2015



INTRODUCTION

• born in Warendorf on 21.10.1987

- born in Warendorf on 21.10.1987
- enrolled at WWU at winter term 2009/10

- born in Warendorf on 21.10.1987
- enrolled at WWU at winter term 2009/10
- first contact with numerics: computer science as secondary subject

- born in Warendorf on 21.10.1987
- enrolled at WWU at winter term 2009/10
- first contact with numerics: computer science as secondary subject
- first application: Bachelor thesis on non-linear friction model ("Nichtlineare Dynamik im verallgemeinerten Prandtl-Tomlinson-Modell")

- born in Warendorf on 21.10.1987
- enrolled at WWU at winter term 2009/10
- first contact with numerics: computer science as secondary subject
- first application: Bachelor thesis on non-linear friction model ("Nichtlineare Dynamik im verallgemeinerten Prandtl-Tomlinson-Modell")
- wide-scale application: Master thesis in Lattice Field Theory ("Confinement-Kriterium für dynamische Materiefelder bei endlichen Temperaturen")

• applicable at strong coupling regimes

- applicable at strong coupling regimes
- non-perturbative effects can be studied

- applicable at strong coupling regimes
- non-perturbative effects can be studied

Formalism:

- applicable at strong coupling regimes
- non-perturbative effects can be studied

Formalism:

• replace continuous space-time by 4-dimensional lattice with euclidean action (via Wick-rotation of time $t \rightarrow i\tau$)

- applicable at strong coupling regimes
- non-perturbative effects can be studied

Formalism:

• replace continuous space-time by 4-dimensional lattice with euclidean action (via Wick-rotation of time $t \to i\tau$)



- applicable at strong coupling regimes
- non-perturbative effects can be studied

Formalism:

• replace continuous space-time by 4-dimensional lattice with euclidean action (via Wick-rotation of time $t \to i\tau$)



• matter fields are classical fields on lattice points

- applicable at strong coupling regimes
- non-perturbative effects can be studied

Formalism:

• replace continuous space-time by 4-dimensional lattice with euclidean action (via Wick-rotation of time $t \to i\tau$)



- matter fields are classical fields on lattice points
- gauge fields are realised on links connecting lattice points

Find operator, which can distinguish different phases in QCD.

• studied model is simplified from electro-weak sector of Standard Model

- studied model is simplified from electro-weak sector of Standard Model
- scalar Higgs-field ϕ with $SU(2)_L$ -charge does not couple to fermionic sector, electromagnetic interactions neglected

- studied model is simplified from electro-weak sector of Standard Model
- scalar Higgs-field ϕ with $SU(2)_L$ -charge does not couple to fermionic sector, electromagnetic interactions neglected
- SU(2)-Higgs-model is non-abelian, shows color-confinement, deconfinement-phase at high temperatures

- studied model is simplified from electro-weak sector of Standard Model
- scalar Higgs-field ϕ with $SU(2)_L$ -charge does not couple to fermionic sector, electromagnetic interactions neglected
- SU(2)-Higgs-model is non-abelian, shows color-confinement, deconfinement-phase at high temperatures
- "scalar quarks" require less computing resources

Construct operators to distinguish different phases; without dynamic matter fields Wilson-loop is good operator:

Construct operators to distinguish different phases; without dynamic matter fields Wilson-loop is good operator:

• consider static (infinitely heavy) quark-antiquark pair

Construct operators to distinguish different phases; without dynamic matter fields Wilson-loop is good operator:

• consider static (infinitely heavy) quark-antiquark pair



Construct operators to distinguish different phases; without dynamic matter fields Wilson-loop is good operator:

- consider static (infinitely heavy) quark-antiquark pair
- quark-antiquark pair is connected by color-gauge field



Construct operators to distinguish different phases; without dynamic matter fields Wilson-loop is good operator:

- consider static (infinitely heavy) quark-antiquark pair
- quark-antiquark pair is connected by color-gauge field



• Wilson-loop $W(\mathscr{C})$ is product of gauge links

Construct operators to distinguish different phases; without dynamic matter fields Wilson-loop is good operator:

- consider static (infinitely heavy) quark-antiquark pair
- quark-antiquark pair is connected by color-gauge field

• Wilson-loop $W(\mathscr{C})$ is product of gauge links





Construct operators to distinguish different phases; without dynamic matter fields Wilson-loop is good operator:

- consider static (infinitely heavy) quark-antiquark pair
- quark-antiquark pair is connected by color-gauge field

- Wilson-loop $W(\mathscr{C})$ is product of gauge links
- has good overlap with states corresponding to flux-tube
 → large signal









Problems for Wilson-loop criterion:



Problems for Wilson-loop criterion:

• Wilson-loop has poor overlap with mesonic states



Problems for Wilson-loop criterion:

- Wilson-loop has poor overlap with mesonic states
- Wilson-loop exhibits circumference-law behaviour in both phases

Proposal of K. Fredenhagen and M. Marcu: Construct operator with dynamical quark fields Ψ at the location of static ones

Proposal of K. Fredenhagen and M. Marcu: Construct operator with dynamical quark fields Ψ at the location of static ones

• consider series of dipole states

$$|\Phi_{xy}^{(n)}\rangle = \sum_{i} \overline{\Psi}_{i}(x) \hat{T}^{n} U(\mathscr{C}_{xy}) \hat{T}^{-n} \Psi_{i}(y) |0\rangle$$

Proposal of K. Fredenhagen and M. Marcu: Construct operator with dynamical quark fields Ψ at the location of static ones

• consider series of dipole states

$$|\Phi_{xy}^{(n)}\rangle = \sum_{i} \overline{\Psi}_{i}(x) \hat{T}^{n} U(\mathscr{C}_{xy}) \hat{T}^{-n} \Psi_{i}(y) |0\rangle$$

• pictographic representation:



• complete Fredenhagen-Marcu-operator (FM-operator):

$$\rho_{FM} = \lim_{R=|x-y|\to\infty} \frac{|\langle \Phi_{xy}^{(n)}|0\rangle|^2}{|\Phi_{xy}^{(n)}|^2}$$
$$= \lim_{R=|x-y|\to\infty} \frac{|\sum_i \langle 0|\overline{\Psi}_i(x)U^{(n)}(\mathscr{C}_{xy})\Psi_i(y)|0\rangle|^2}{\langle \mathscr{C}_{xy} \cdot \Theta \mathscr{C}_{xy} \rangle}$$





• in confinement phase numerator and denominator exhibit circumference-law behaviour $\mathbb{R} \rightarrow \infty$

$$\rightarrow \rho_{FM} \stackrel{n \to \infty}{=} const. \neq 0$$



• in confinement phase numerator and denominator exhibit circumference-law behaviour

 $\rightarrow \rho_{FM} \stackrel{\mathbf{R} \rightarrow \infty}{=} const. \neq 0$

• deconfinement phase: dipole state approximates free color-charge, orthogonal to color-neutral vacuum $\rightarrow \rho_{FM} \stackrel{R \to \infty}{=} 0$



• in confinement phase numerator and denominator exhibit circumference-law behaviour

 $\rightarrow \rho_{FM} \stackrel{R \to \infty}{=} const. \neq 0$

- deconfinement phase: dipole state approximates free color-charge, orthogonal to color-neutral vacuum $\rightarrow \rho_{FM} \stackrel{R \to \infty}{=} 0$
- FM-operator tests for color-confinement, able to distinguish between phases



circumference-law behaviour $\rightarrow \rho_{FM} \stackrel{R \to \infty}{=} const. \neq 0$

 $\rightarrow \rho_{FM} \stackrel{\mathbf{K} \rightarrow \infty}{=} 0$ FM-operator tests for color-confinement, able to di

• in confinement phase numerator and denominator exhibit

• deconfinement phase: dipole state approximates free color-charge, orthogonal to color-neutral vacuum

• FM-operator tests for color-confinement, able to distinguish between phases

BUT:

In confinement phase R must be larger than string-breaking distance!

• FM-operator can distinguish between symmetric and Higgs phase at high physical temperature



• FM-operator can distinguish between symmetric and Higgs phase at high physical temperature



• distinction between confinement and deconfinement phases very difficult

• FM-operator can distinguish between symmetric and Higgs phase at high physical temperature



- distinction between confinement and deconfinement phases very difficult
- FM-operator shows signs of phase-transition, but no convincing evidence

• QCD corrections to D-meson decays still not very accurate

- QCD corrections to D-meson decays still not very accurate
- decay width of purely leptonic D-meson decays can be written as

$$\Gamma(D \to l\nu) = \frac{G_F^2}{8\pi} |V_{cq}|^2 m_l^2 \left(1 - \frac{m_l^2}{m_D^2}\right)^2 m_D f_D^2$$

- QCD corrections to D-meson decays still not very accurate
- decay width of purely leptonic D-meson decays can be written as

$$\Gamma(D \to l\nu) = \frac{G_F^2}{8\pi} |V_{cq}|^2 m_l^2 \left(1 - \frac{m_l^2}{m_D^2}\right)^2 m_D f_D^2$$

• decay constant f_D can be determined via LQCD

 $\langle 0 | \overline{q} \gamma_{\mu} \gamma_5 c(0) | D(\vec{p}) \rangle = i f_D p_{\mu}$

• \rightarrow determination of decay constants f_D allows for precise measurements of CKM-matrix elements $|V_{cq}|$

- \rightarrow determination of decay constants f_D allows for precise measurements of CKM-matrix elements $|V_{cq}|$
- $\bullet\,\,\rightarrow\,\,{\rm give}$ precise constraints on unitarity and CP-violation

- \rightarrow determination of decay constants f_D allows for precise measurements of CKM-matrix elements $|V_{cq}|$
- \rightarrow give precise constraints on unitarity and CP-violation
- may improve understanding of recent experimental findings of larger than expected CP-violation in the charm quark sector (e.g. by LHCb)
 - \rightarrow hint to new physics?

Use gauge configurations generated within the Coordinated Lattice Simulations (CLS) effort and Alpha collaboration

Use gauge configurations generated within the Coordinated Lattice Simulations (CLS) effort and Alpha collaboration

• simulations are based on openQCD program created by Martin Lüscher and Stefan Schaefer

Use gauge configurations generated within the Coordinated Lattice Simulations (CLS) effort and Alpha collaboration

- simulations are based on openQCD program created by Martin Lüscher and Stefan Schaefer
- meson correlators will be derived on PALMA

Use gauge configurations generated within the Coordinated Lattice Simulations (CLS) effort and Alpha collaboration

- simulations are based on openQCD program created by Martin Lüscher and Stefan Schaefer
- meson correlators will be derived on PALMA

Further possible extensions:

Use gauge configurations generated within the Coordinated Lattice Simulations (CLS) effort and Alpha collaboration

- simulations are based on openQCD program created by Martin Lüscher and Stefan Schaefer
- meson correlators will be derived on PALMA

Further possible extensions:

• modify meson measurement program, conduct charmed-meson spectroscopy

Use gauge configurations generated within the Coordinated Lattice Simulations (CLS) effort and Alpha collaboration

- simulations are based on openQCD program created by Martin Lüscher and Stefan Schaefer
- meson correlators will be derived on PALMA

Further possible extensions:

- modify meson measurement program, conduct charmed-meson spectroscopy
- compare data obtained from CLS configurations (with $n_f = 2 + 1$, no dynamical charm quark) with simulations including dynamical charm

 \rightarrow estimate corrections of dynamical charm to derived quantities

• PANDA experiment still under construction

- PANDA experiment still under construction
- WASA experiment at Jülich very similar, already extensive amount of data taken (lighter mesons like η)

- PANDA experiment still under construction
- WASA experiment at Jülich very similar, already extensive amount of data taken (lighter mesons like η)
- software program available for detector simulation

- PANDA experiment still under construction
- WASA experiment at Jülich very similar, already extensive amount of data taken (lighter mesons like η)
- software program available for detector simulation
- → calculation of experimentally important quantities
 (→ methods applicable to PANDA), but also get insight in detector physics (interplay of detector parts, resolution, limitations)

Thank you for your attention!