



## Münsteranian Torturials on Nonlinear Science edited by Uwe Thiele, Oliver Kamps, Svetlana Gurevich

#### Continuation

# THFI: Sliding drops modelled by a thin film equation for a liquid layer or drop on an inclined homogeneous substrate

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For updates of this text and the accompanying programme files see www.uni-muenster.de/CeNoS/Lehre/Tutorials/auto.html

### 2 thfi: Sliding drops modelled by a thin film equation for a liquid layer or drop on an inclined homogeneous substrate

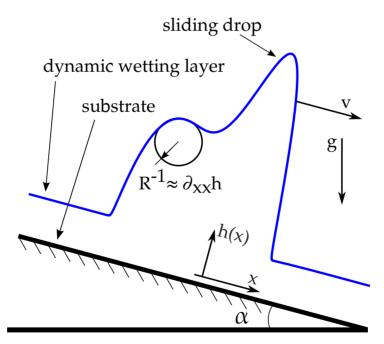
In the tutorial THFI a dimensionless thin film equation closely related to the one in DROP is used. In this case a lateral driving force is included and solutions for sliding drops on an inclined plane are examined. You will calculate steady solutions [in the comoving frame], in particular, drops moving at a constant velocity. The main treated control parameter is the inclination angle.

#### 2.1 Model

This demo illustrates the calculation of sliding drop and surface wave solutions of the dimensionless thin film equation

$$\partial_t h = -\partial_x \left\{ Q(h) \, \partial_x \left[ \partial_{xx} h - \partial_h f(h) \right] + \chi(h) \right\} \tag{2.1}$$

where  $Q(h) = h^3$  is the mobility function and  $\chi(h)$  is the lateral driving. For an inclined substrate it is  $\chi = \alpha Q(h)$  with  $\alpha$  being the inclination angle. For background information see [1, 2]. Examples of similar calculations with various f(h) can be found in [3, 4, 5]. The term in square brackets corresponds to a pressure. The pressure used here is the same as in the demo drop [6] where it is explained in detail. Our aim is to study sliding droplet or surface wave



**Figure 2.1:** The basic geometry for this problem. The profile is an actual solution to the equation of this problem.

solutions that are steady in some co-moving frame, i.e., the drops slide with constant speed v and shape. We introduce the coordinate in the frame moving with v by  $\tilde{x} = x - vt$  and obtain from Eq. (2.1) after dropping the tildes

$$-v\partial_x h = -\partial_x \left\{ Q(h) \partial_x \left[ \partial_{xx} h - \partial_h f(h) \right] + \chi(h) \right\}$$
 (2.2)

Eq. (2.2) is integrated once to obtain

$$0 = Q(h) \partial_x \left[ \partial_{xx} h - \partial_h f(h) \right] + \chi(h) - vh - C_0$$
 (2.3)

where the constant  $C_0$  corresponds to the flux in the co-moving frame and the unknown velocity v can be seen as an eigenvalue in a nonlinear eigenvalue problem. Note that  $C_0 + vh$  is the flux in the laboratory frame.

To use the continuation toolbox auto07p [7], we first write (2.3) as a system of first-order autonomous ordinary differential equations on the interval [0, 1]. Therefore, we introduce the variables  $u_1 = h - h_0$ ,  $u_2 = dh/dx$  and  $u_2 = d^2h/dx^2$ , use  $\chi(h) = \alpha Q(h)$ , and obtain from equation (2.3) the 3d dynamical system (NDIM = 3)

$$\dot{u}_{1} = Lu_{2} 
\dot{u}_{2} = Lu_{3} 
\dot{u}_{3} = L \left[ u_{2}f''(u_{1} + h_{0}) - \alpha + \frac{v(u_{1} + h_{0}) + C_{0}}{Q(u_{1} + h_{0})} \right].$$
(2.4)

where L is the physical domain size, and dots and primes denote derivatives with respect to  $\xi \equiv x/L$  and h, respectively. The advantage of the used form is that the fields  $u_1(\xi)$ ,  $u_2(\xi)$  and  $u_3(\xi)$  correspond to the correctly scaled physical fields  $h(L\xi)$ ,  $\partial_x h(L\xi)$  and  $\partial_{xx} h(L\xi)$ . We use periodic boundary conditions for all  $u_i$  (NBC = 3) that take the form

$$u_1(0) = u_1(1), (2.5)$$

$$u_2(0) = u_2(1), (2.6)$$

$$u_3(0) = u_3(1), (2.7)$$

and integral conditions for mass conservation and computational pinning (to break the translational symmetry that the solutions have on the considered homogeneous substrate) (NINT = 2). The integral condition for mass conservation takes the form

$$\int_0^1 u_1 \, \mathrm{d}\xi = 0. \tag{2.8}$$

There are two ways to start the continuation. Either (i) one sets  $\alpha=0$  and uses as in the demo drop [6] the starting solution consisting of small amplitude harmonic modulation of wavelength  $L_c=2\pi/k_c$  where  $k_c=\sqrt{-f''(h_0)}$  is the critical wavenumber for the linear instability of a flat film of thickness  $h_0$  and also sets initially v=0 and  $C_0=0$ ; or (ii) one starts at some  $\alpha\neq 0$ , uses small amplitude harmonic starting solution with  $L_c=2\pi/k_c$  and initialises  $v=\alpha Q'(h_0)$  and  $C_0=\alpha Q(h_0)-vh_0$ . In the present demo we use the former option.

The number of free (continuation) parameters is given by NCONT = NBC + NINT - NDIM + 1 and is here equal to 3. For more details see [2].

#### **2.2 Runs:**

Python interface command line	Terminal command line
auto	

**run 1:** Determine steady solutions on the horizontal substrate as a function of domain size L, starting at the critical  $L_c$  with a small amplitude sinusoidal solution. Mean thickness is fixed. Compute the branch of periodic solutions for  $h_0 = 5$  continued in L (PAR(5)) up to L = 400.

**Remaining true continuation parameters:**  $C_0$  (PAR(6)) and v (PAR(42));

**Parameters:** IPS= 4, ISP= 0, ISW= 1, ICP= [5, 6, 42],

Start data from function stpnt (IRS= 0)

save output-files as b.d1, s.d1, d.d1

r1 = run(e = 'thfi', c = 'thfi.1', sv = 'tf1') @@R thfi 1 @sv tf1

run 11: Restart at domain size L = 400, keep mean thickness  $h_0$  fixed and incline substrate to observe transition from sliding drops to surface waves.

Continued in inclination  $\alpha$  (PAR(41)) for fixed domain size L. Stop at  $\alpha = 0.01$ 

**Remaining true continuation parameters:**  $C_0$  (PAR(6)) and and v (PAR(42));

Other output: abs. value of minimal slope of h (PAR(46)), i.e., advancing dynamic contact angle  $\theta_{\rm adv}$ ; maximal slope of h (PAR(47)), i.e., receding dynamic contact angle  $\theta_{\rm rec}$ ; Parameter: IPS= 4, ISP= 0, ISW= 1, ICP= [41, 6, 42, 46, 47],

Start at final result of run 1: IRS= 7

save output-files as b.d11, s.d11, d.d11

**Table 2.1:** Commands for running demo thfi.

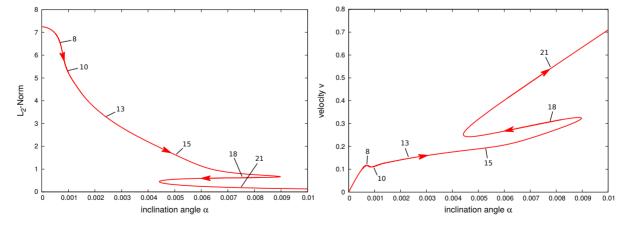
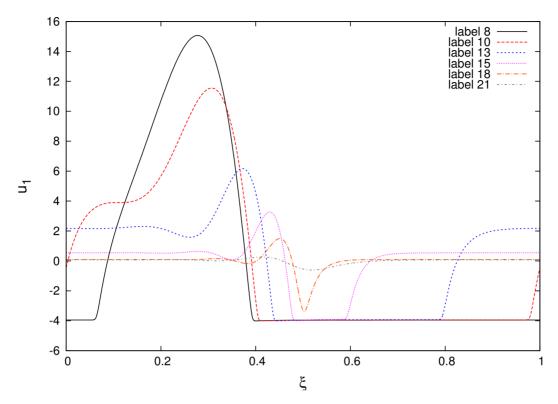


Figure 2.2: An illustration for run 11 of demo thfi is given. The  $L_2$ -norm of stationary solutions (steady in the frame moving with v) is shown in dependence of the principal continuation parameter inclination angle  $\alpha$  (par(41)) for fixed domain size L=400 (par(5)) and mean thickness  $h_0=5$  (par(1)).



**Figure 2.3:** Stationary film thickness profiles corresponding to Fig. (2.2) at inclination angles as given in the legend. The profiles are represented by  $u_1(\xi) = h(\xi L) - h_0$  where  $h_0 = 5$ .

#### 2.3 Remarks:

- Beside the NCONT true continuation parameters that have to be given as ICP in the c.\* parameter file, one may list other output parameters as defined in the subroutine PVLS in the \*.f90 file.
- As in the demo drop [6] one may define other integral conditions to determine integral measures one might be interested in
- Screen output and command line commands are provided in README file.

#### **2.4 Tasks:**

After running the examples, you should try to implement your own adaptations, e.g.:

- 1. Redo runs 1 and 11 for other values of  $h_0$ .
- 2. Try to run instead of run 11 a continuation with fixed  $C_0$ . You need to 'set free' another parameter. This will not work if  $\alpha = 0$  initially. Start from a solution of the original run 11.
- 3. Include additional integral condition(s), to measure characteristics of interest. These might be the surface energy, wetting energy, total energy dissipation.

- 4. Replace the used Derjaguin pressure by a different one that you get from the literature.
- 5. Replace the used mobility function Q(h) by a different one. An option is  $Q = h^2(h + l_s)/3\eta$  that incorporates slip of the liquid at the solid substrate ( $l_s$  is the slip length).

#### References

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