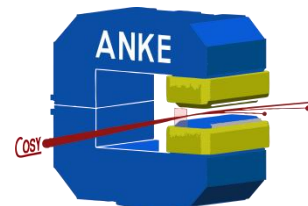


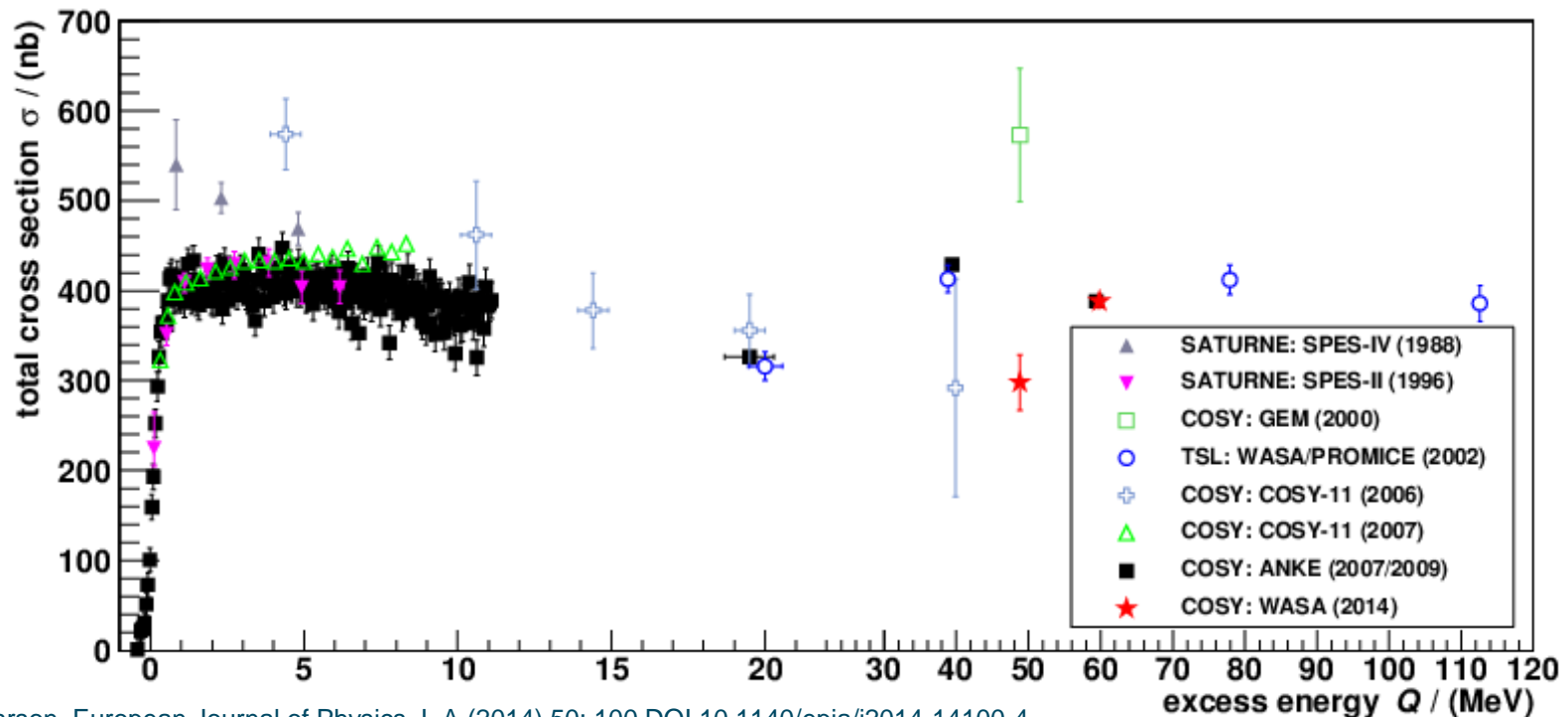
# Investigation of different normalization reactions for $dp$ collisions at ANKE

**Christopher Fritzscht for the ANKE collaboration**

Westfälische Wilhelms-Universität Münster, Institut für Kernphysik  
DPG Heidelberg (HK 37.5)

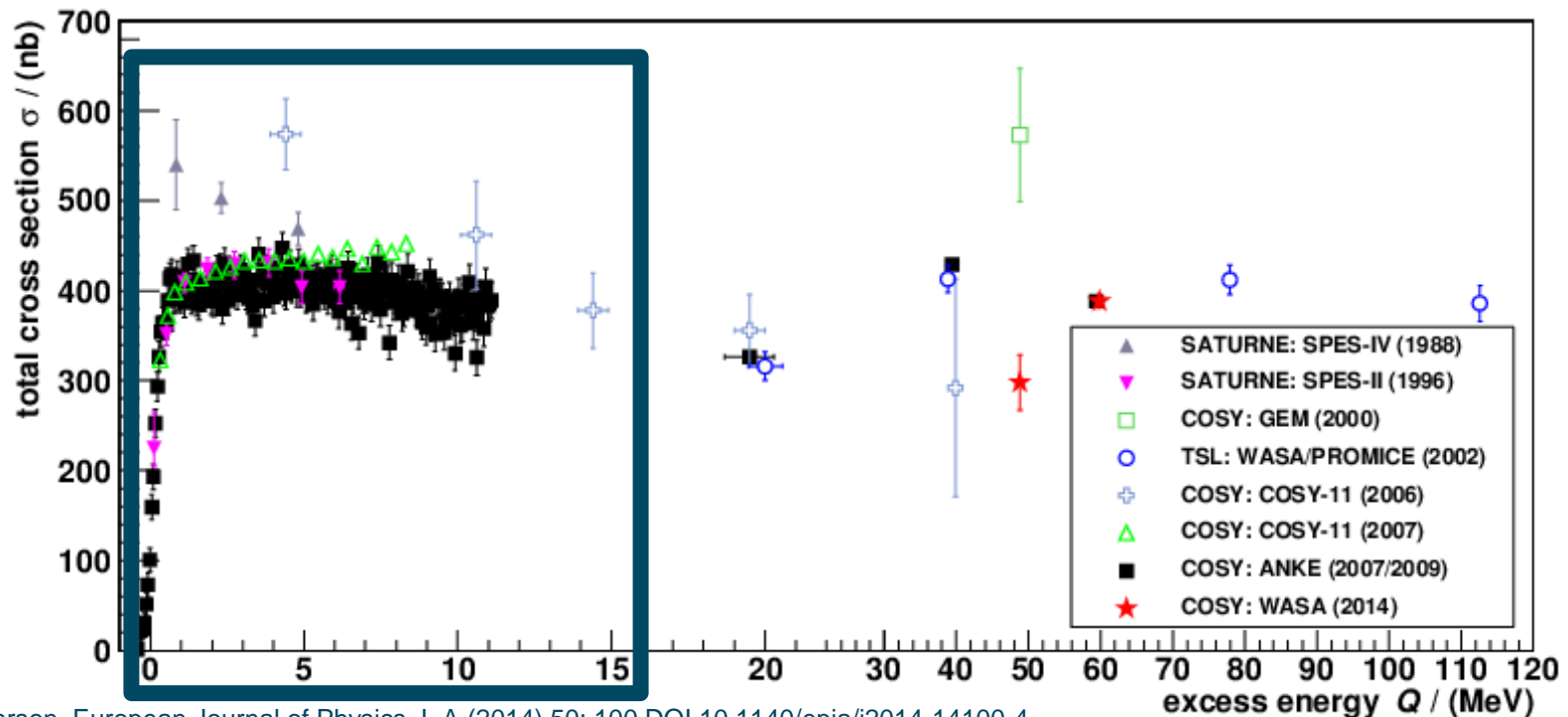
March 24th, 2015





P. Adlarson, European Journal of Physics J. A (2014) 50: 100 DOI 10.1140/epja/i2014-14100-4

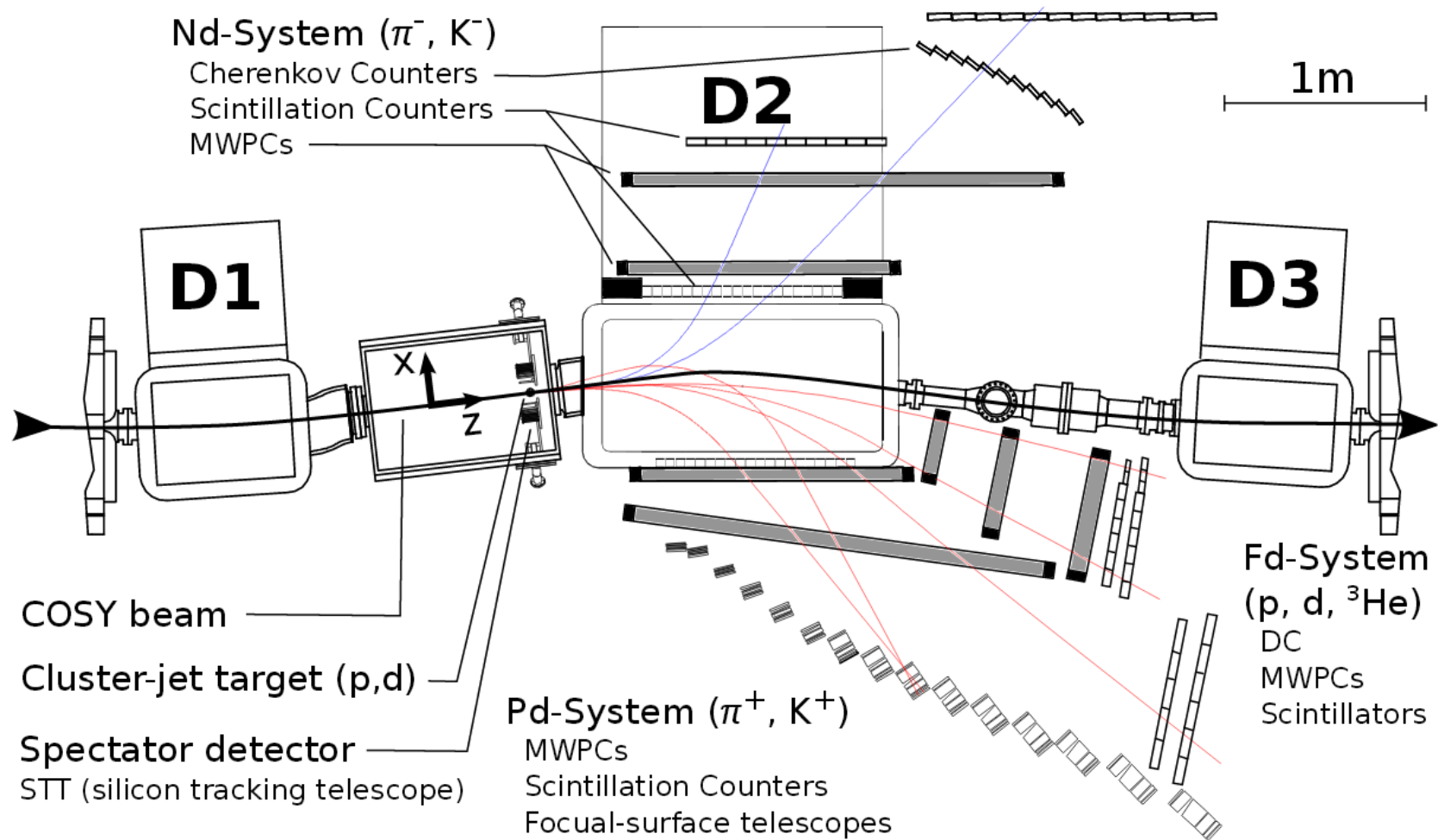
- total and differential cross sections of the reaction  $dp \rightarrow {}^3\text{He}\eta$  are of special interest
- indication of a quasi bound state of the  ${}^3\text{He}\eta$ -system
- high precision data from the ANKE spectrometer at COSY up to  $Q = 15$  MeV
  - ⇒ extraction of total and differential cross sections
  - ⇒ careful luminosity determination is necessary



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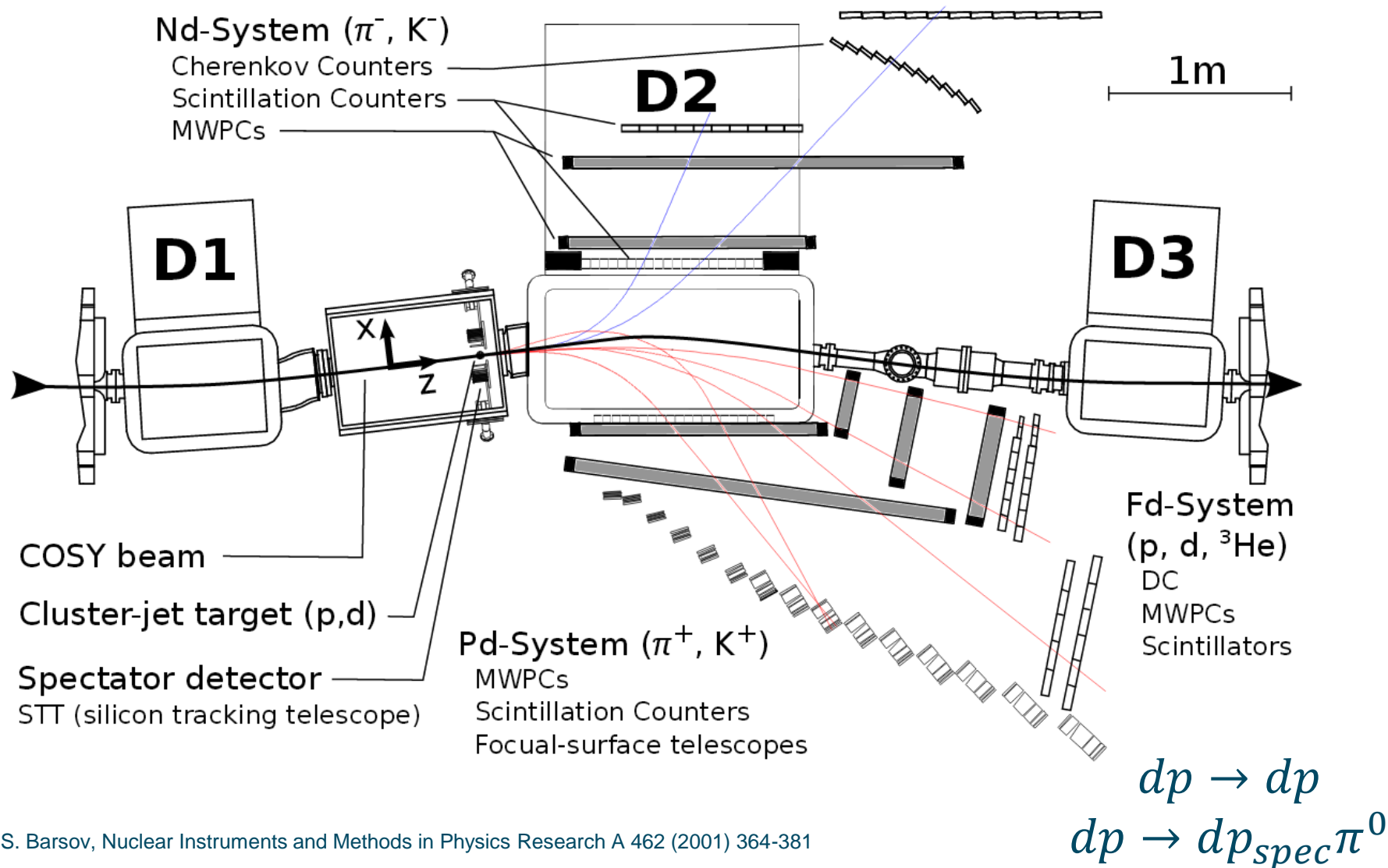
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# ANKE spectrometer

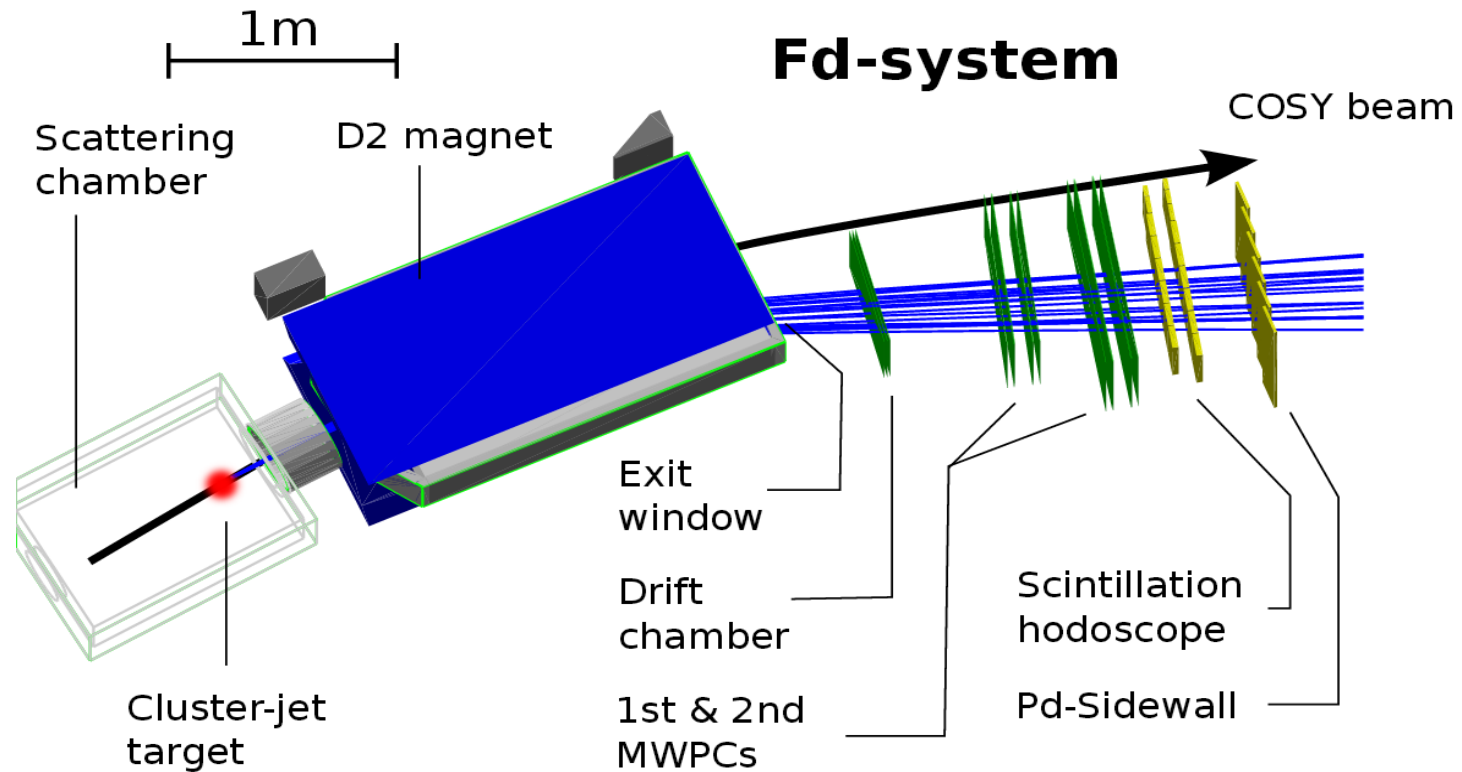


S. Barsov, Nuclear Instruments and Methods in Physics Research A 462 (2001) 364-381

# ANKE spectrometer



S. Barsov, Nuclear Instruments and Methods in Physics Research A 462 (2001) 364-381




- track reconstruction
  - ⇒ one multiwire drift chamber and two multiwire proportional chambers
- energy loss measurements
  - ⇒ two layers of scintillator hodoscopes

- main goals: high precision determination of the  $\eta$  meson mass ( $dp \rightarrow {}^3\text{He}\eta$ ) and studies on the two pion production ( $dp \rightarrow {}^3\text{He}\pi^+\pi^-$ )
- supercycle mode: one supercycle consists of up to 7 different beam momenta  
⇒ 19 beam momenta overall

flattop	1st supercycle $p_d / (\text{MeV}/c)$	2nd supercycle $p_d / (\text{MeV}/c)$	3rd supercycle $p_d / (\text{MeV}/c)$
0	3120.17(17)	3120.00(22)	3125
2	3146.41(17)	3147.35(17)	3146
3	3148.45(17)	3150.42(17)	-----
4	3152.45(17)	3154.49(17)	3157.48(22)
5	3158.71(17)	3162.78(17)	3160.62(22)
6	3168.05(17)	3172.15(17)	-----
7	3177.51(17)	3184.87(17)	3204.16(23)

below  $\eta$  threshold ( $dp \rightarrow {}^3\text{He}\eta$ ) for a smooth background description

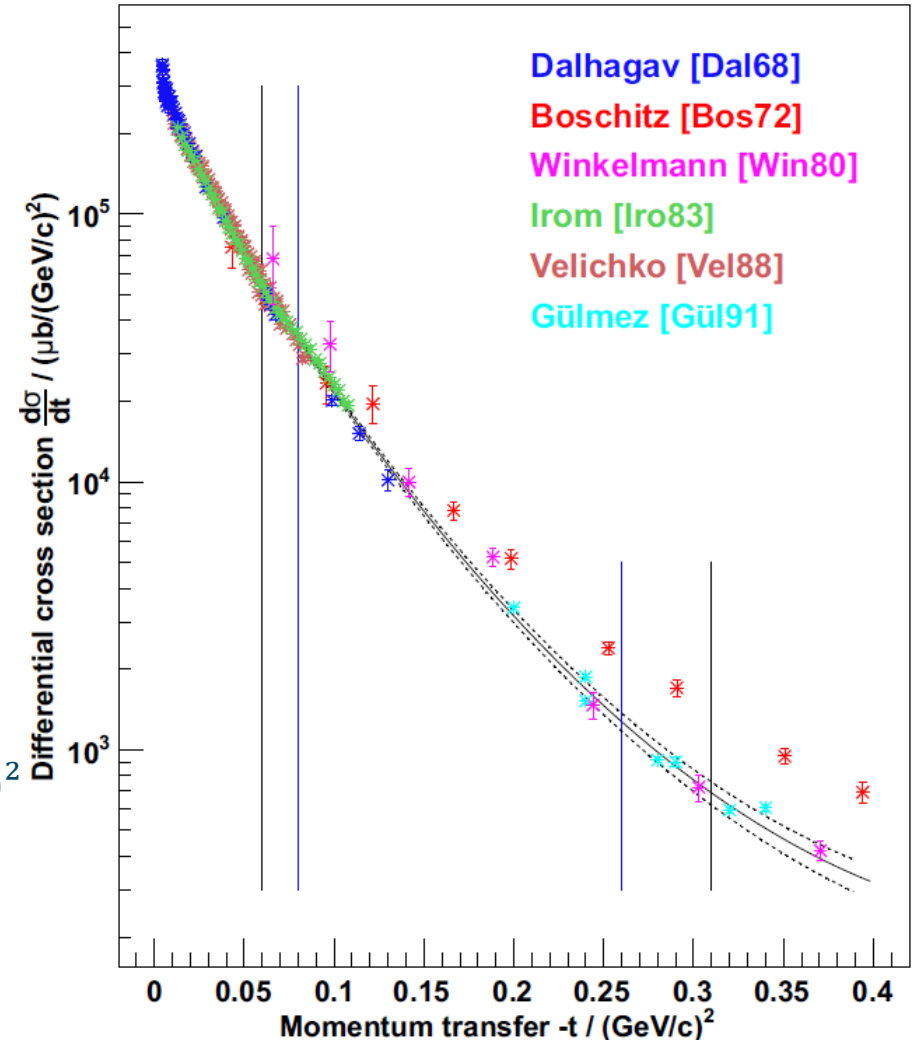


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7	3177.51(17)	3184.87(17)	3204.16(23)



# Luminosity via $dp \rightarrow dp$

- advantages of the  $dp$ -elastic scattering as normalization reaction:
  - ⇒ broad data base of available differential cross section
  - ⇒ high differential cross sections ensure good statistics
  - ⇒ excellent signal-to-background ratio
- differential cross section as function of momentum transfer
  - ⇒ independent of beam momentum
- ANKE acceptance:  $-t = (0.06 - 0.31) (\text{GeV}/c)^2$
- range for luminosity determination  
 $-t = (0.08 - 0.26) (\text{GeV}/c)^2$



# Identification of $dp \rightarrow dp$

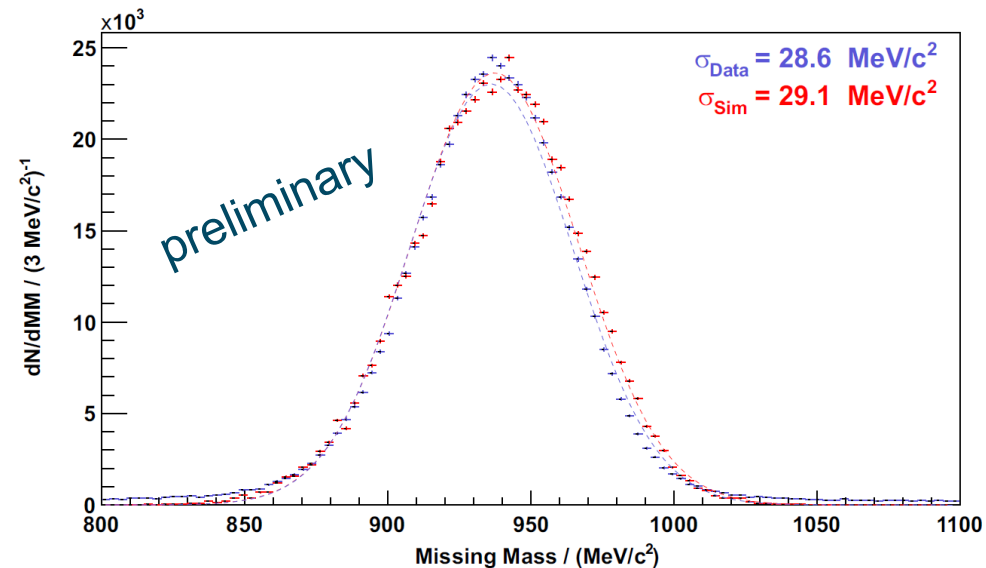
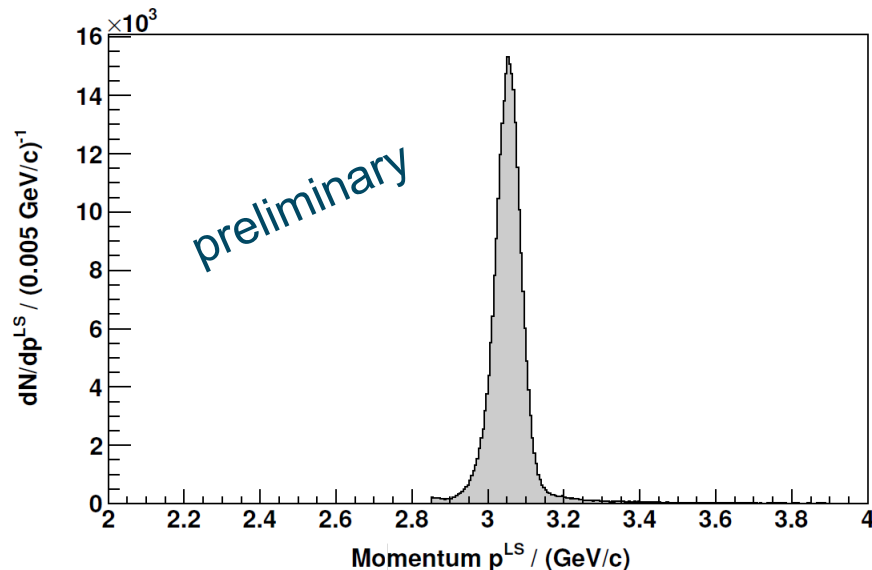
- identification of the reaction via the missing mass technique

$$\text{Missing Mass} = |\mathbb{P}_{p(\text{ejectile})}| = |\mathbb{P}_{d(\text{beam})} + \mathbb{P}_{p(\text{target})} - \mathbb{P}_{d(\text{ejectile})}|$$

⇒ deuterons in the Fd-system


⇒ momentum close to beam momentum

⇒ low momentum transfer onto the target proton



$$L_{int} = \frac{N_{det}(\Delta t)}{\int_{t_1}^{t_2} \left( \frac{d\sigma}{dt}(t) \right)_{Ref} dt}$$

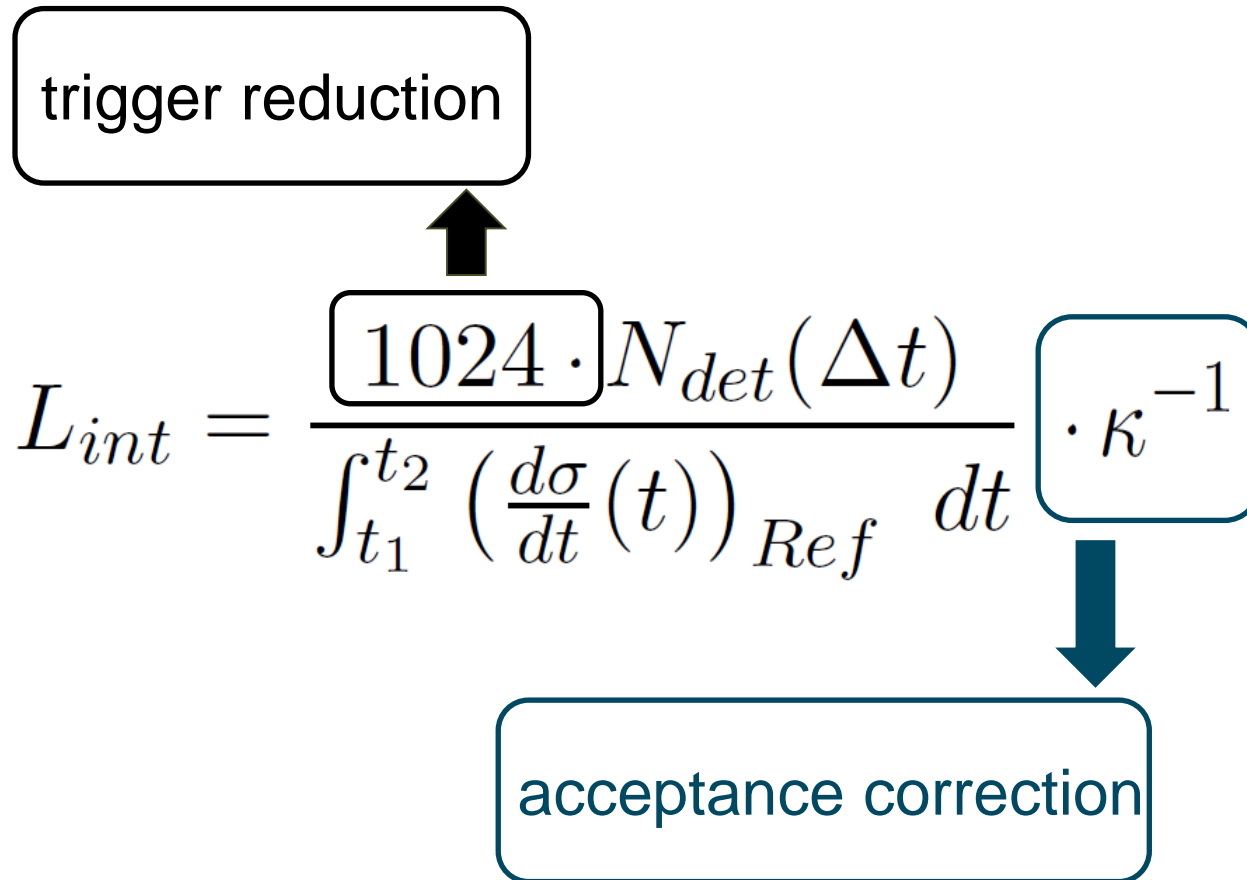
trigger reduction

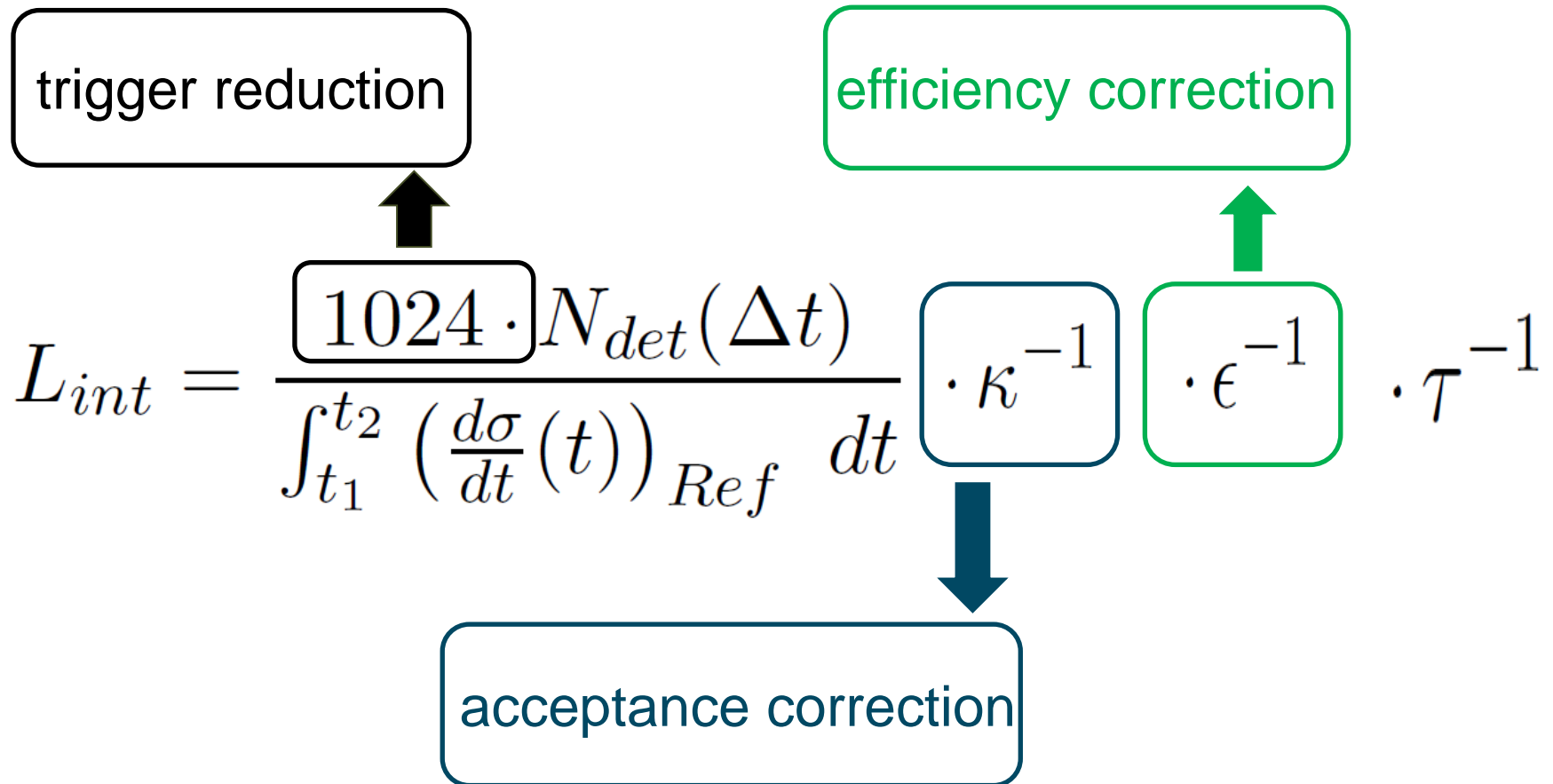

$$L_{int} = \frac{1024 \cdot N_{det}(\Delta t)}{\int_{t_1}^{t_2} \left( \frac{d\sigma}{dt}(t) \right)_{Ref} dt}$$

trigger reduction

$$L_{int} = \frac{1024 \cdot N_{det}(\Delta t)}{\int_{t_1}^{t_2} \left( \frac{d\sigma}{dt}(t) \right)_{Ref} dt} \cdot \kappa^{-1}$$

acceptance correction

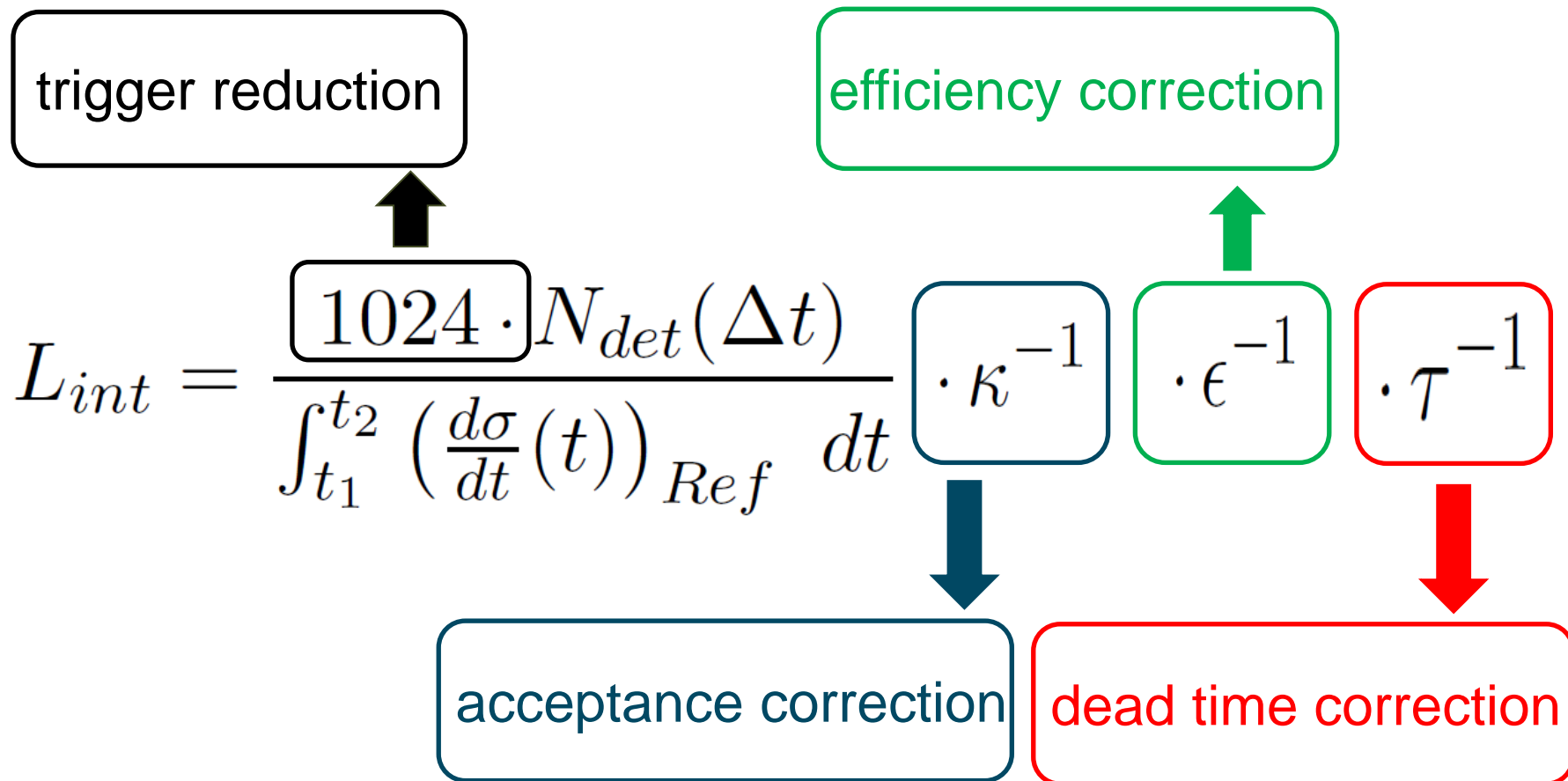




The diagram illustrates the luminosity determination formula with three annotations:

- trigger reduction**: A black box above the numerator  $1024 \cdot N_{det}(\Delta t)$  with a black arrow pointing to it.
- efficiency correction**: A green box above the term  $\cdot \epsilon^{-1}$  with a green arrow pointing to it.
- acceptance correction**: A blue box below the term  $\cdot \kappa^{-1}$  with a blue arrow pointing to it.

$$L_{int} = \frac{1024 \cdot N_{det}(\Delta t)}{\int_{t_1}^{t_2} \left( \frac{d\sigma}{dt}(t) \right)_{Ref} dt} \cdot \kappa^{-1} \cdot \epsilon^{-1} \cdot \tau^{-1}$$



The diagram illustrates the formula for luminosity determination,  $L_{int}$ , and the various correction factors applied to it. The formula is:

$$L_{int} = \frac{1024 \cdot N_{det}(\Delta t)}{\int_{t_1}^{t_2} \left( \frac{d\sigma}{dt}(t) \right)_{Ref} dt} \cdot \kappa^{-1} \cdot \epsilon^{-1} \cdot \tau^{-1}$$

The factors are represented in boxes with arrows indicating their relationship to the formula:

- trigger reduction** (black box) is indicated by a black arrow pointing up to the  $1024$  factor.
- efficiency correction** (green box) is indicated by a green arrow pointing up to the  $\epsilon^{-1}$  factor.
- acceptance correction** (blue box) is indicated by a blue arrow pointing down from the  $\kappa^{-1}$  factor.
- dead time correction** (red box) is indicated by a red arrow pointing down from the  $\tau^{-1}$  factor.

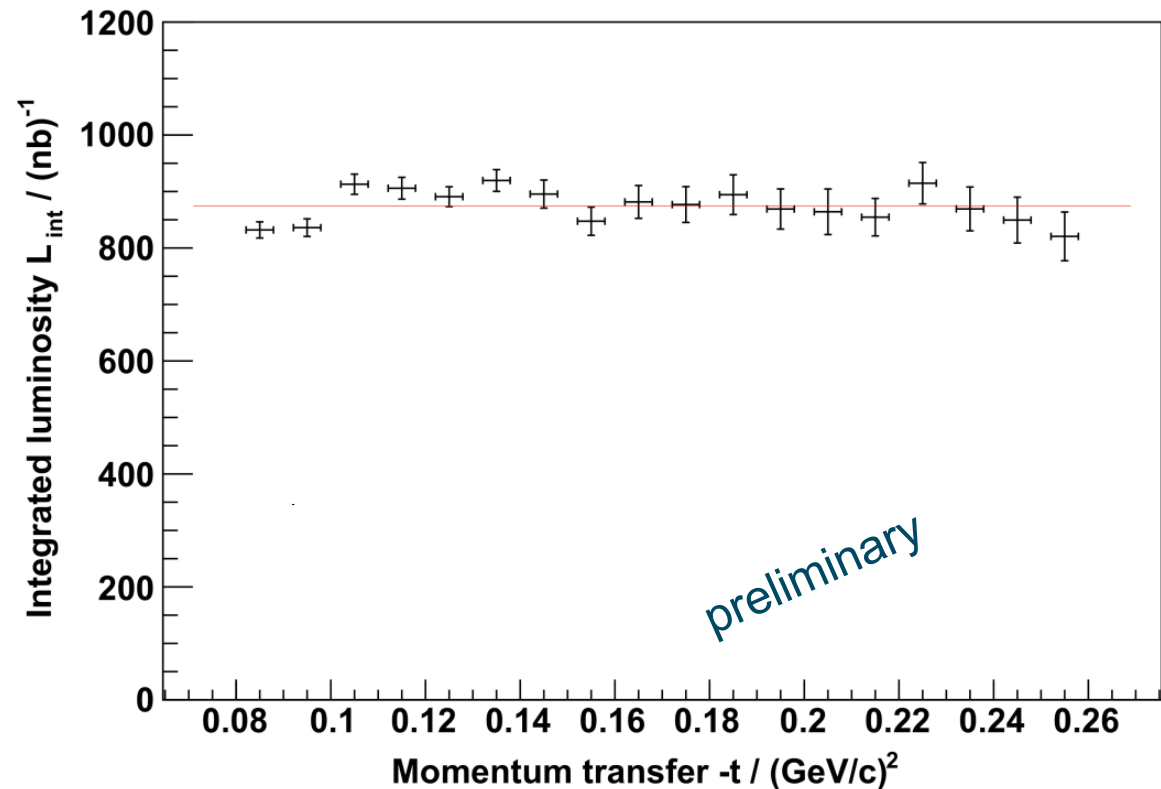
# Results for $dp \rightarrow dp$

- luminosity should be independent of momentum transfer
- determination performed for 18 momentum transfer bins ( $\Delta t = 0.01 \text{ (GeV/c)}^2$ )  
⇒ for each of the 19 beam momenta
- luminosity precision achieved:

$$\Delta L_{\text{stat}} = 1\%$$

$$\Delta L_{\text{sys}} = 6\%$$

- improvement by at least a factor of two
- relative normalization via the reactions  $dp \rightarrow p_{\text{spec}}X$  by determination of the number of spectator protons in the Fd-system





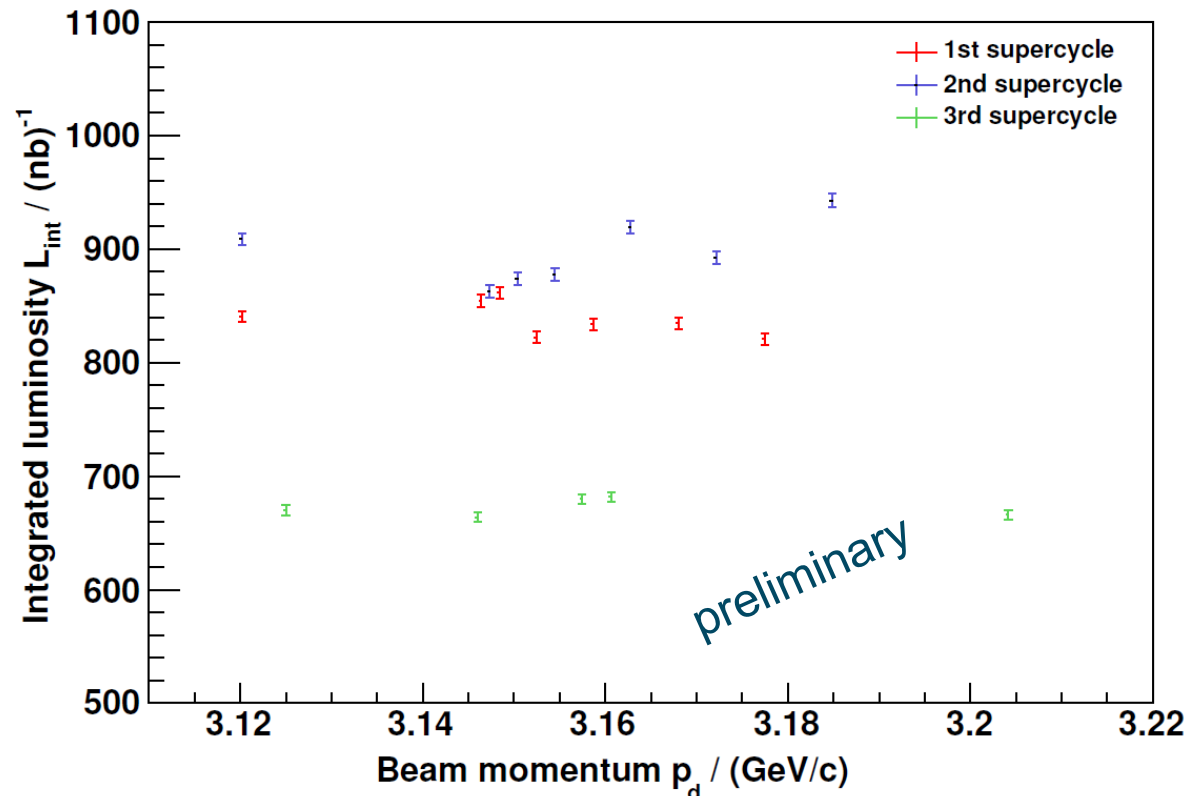
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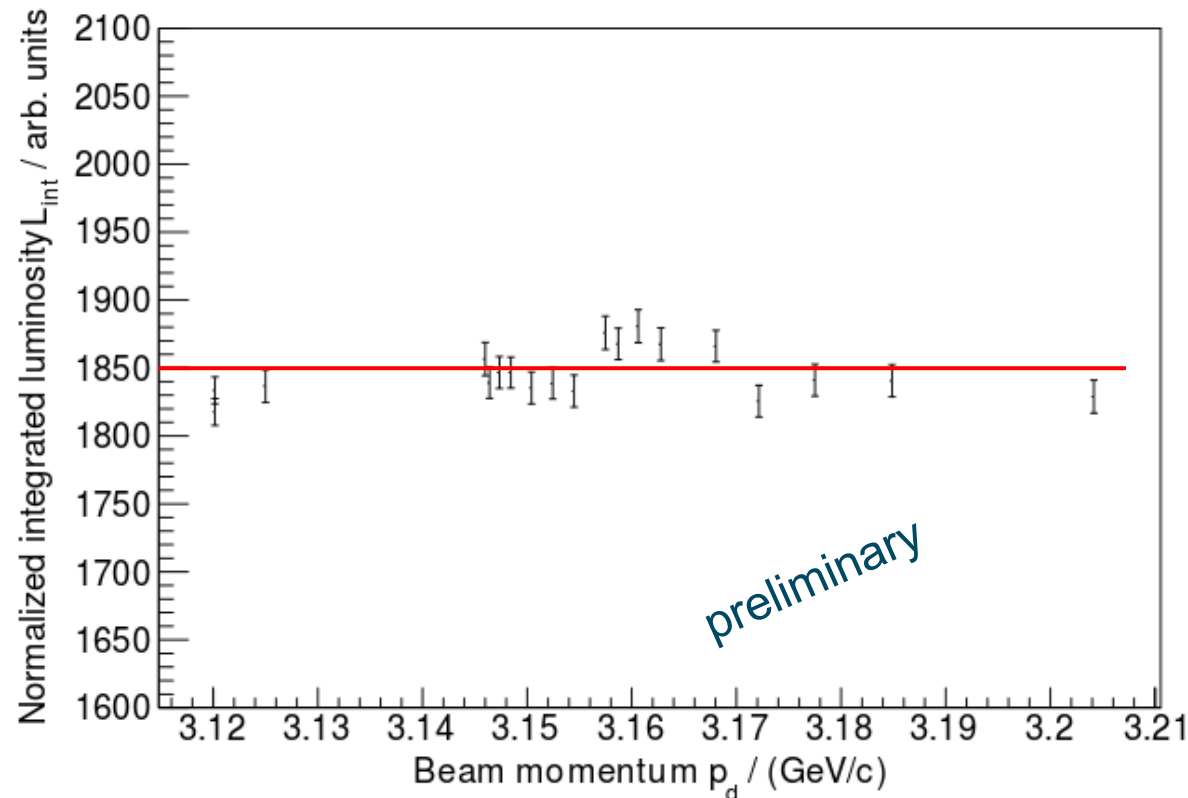
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- luminosity should be independent of momentum transfer
- determination performed for 18 momentum transfer bins ( $\Delta t = 0.01 \text{ (GeV/c)}^2$ )  
⇒ for each of the 19 beam momenta
- luminosity precision achieved:

$$\Delta L_{\text{stat}} = 1\%$$

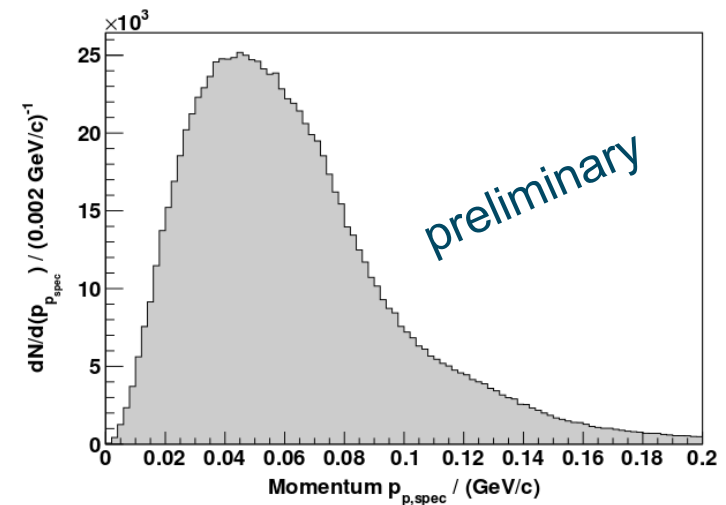
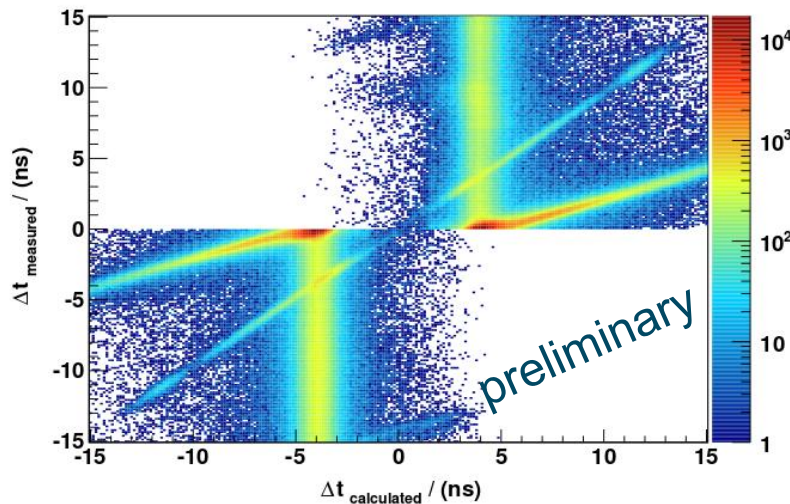
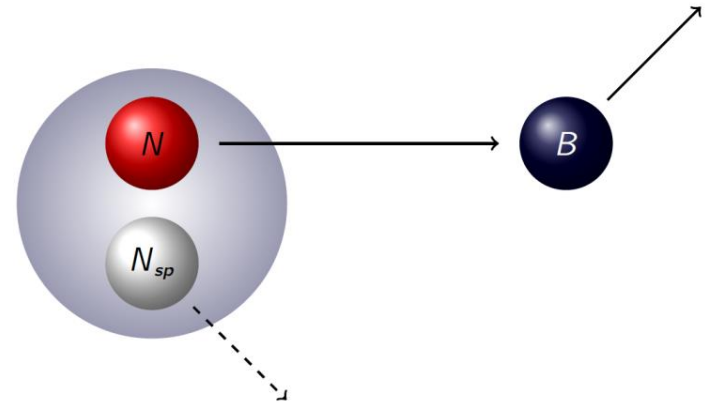
$$\Delta L_{\text{sys}} = 6\%$$

- improvement by at least a factor of two
- relative normalization via the reactions  $dp \rightarrow p_{\text{spec}}X$  by determination of the number of spectator protons in the Fd-system

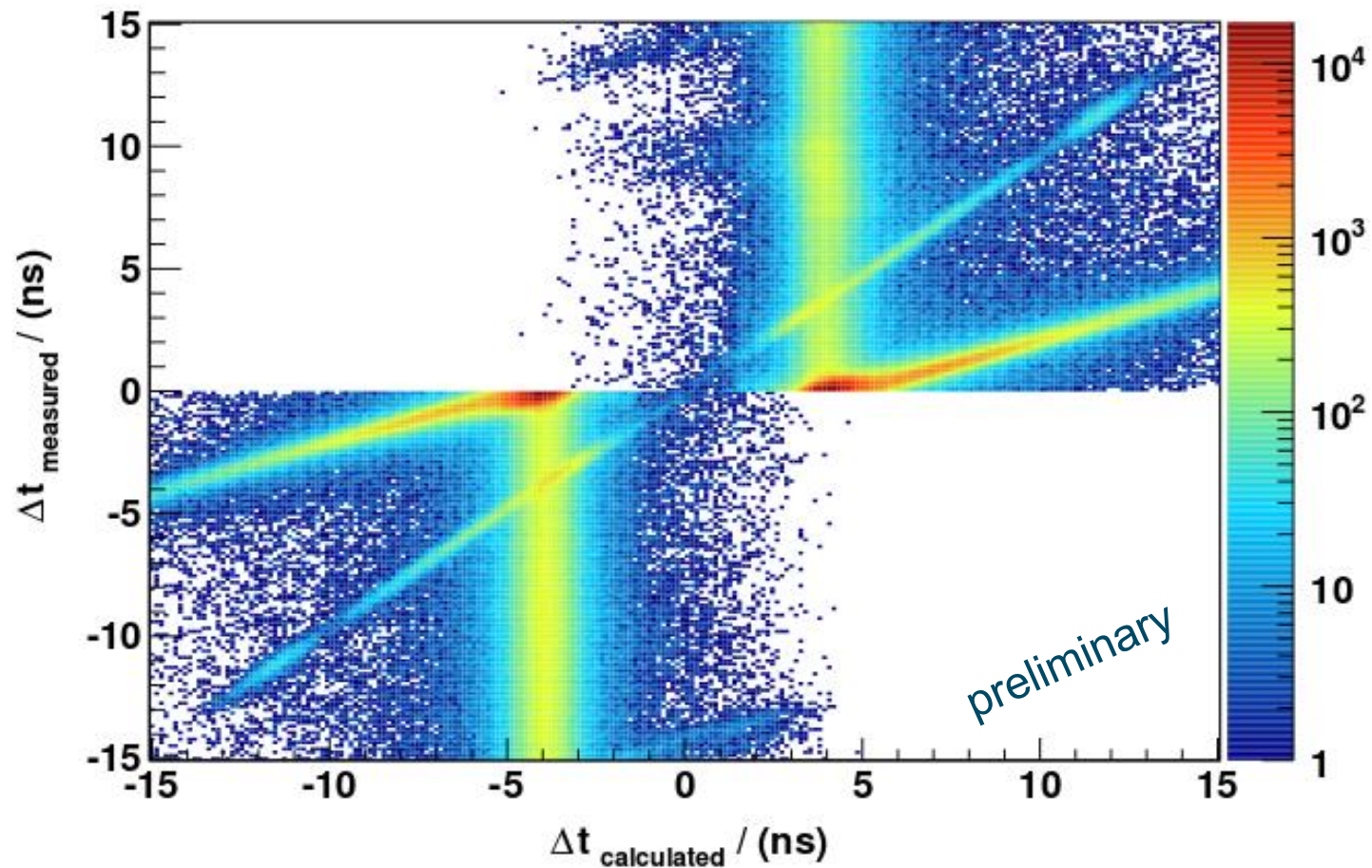


# Luminosity via $dp \rightarrow dp_{spec}\pi^0$

- to verify the results of the  $dp$ -elastic scattering  
⇒ independent normalization required  
⇒ reaction used:  $dp \rightarrow dp_{spec}\pi^0$
- $p_{spec}$  = half of beam momentum + fermi motion
- identification of the  $dp$  pairs in the Fd-system  
⇒ time-of-flight measurement

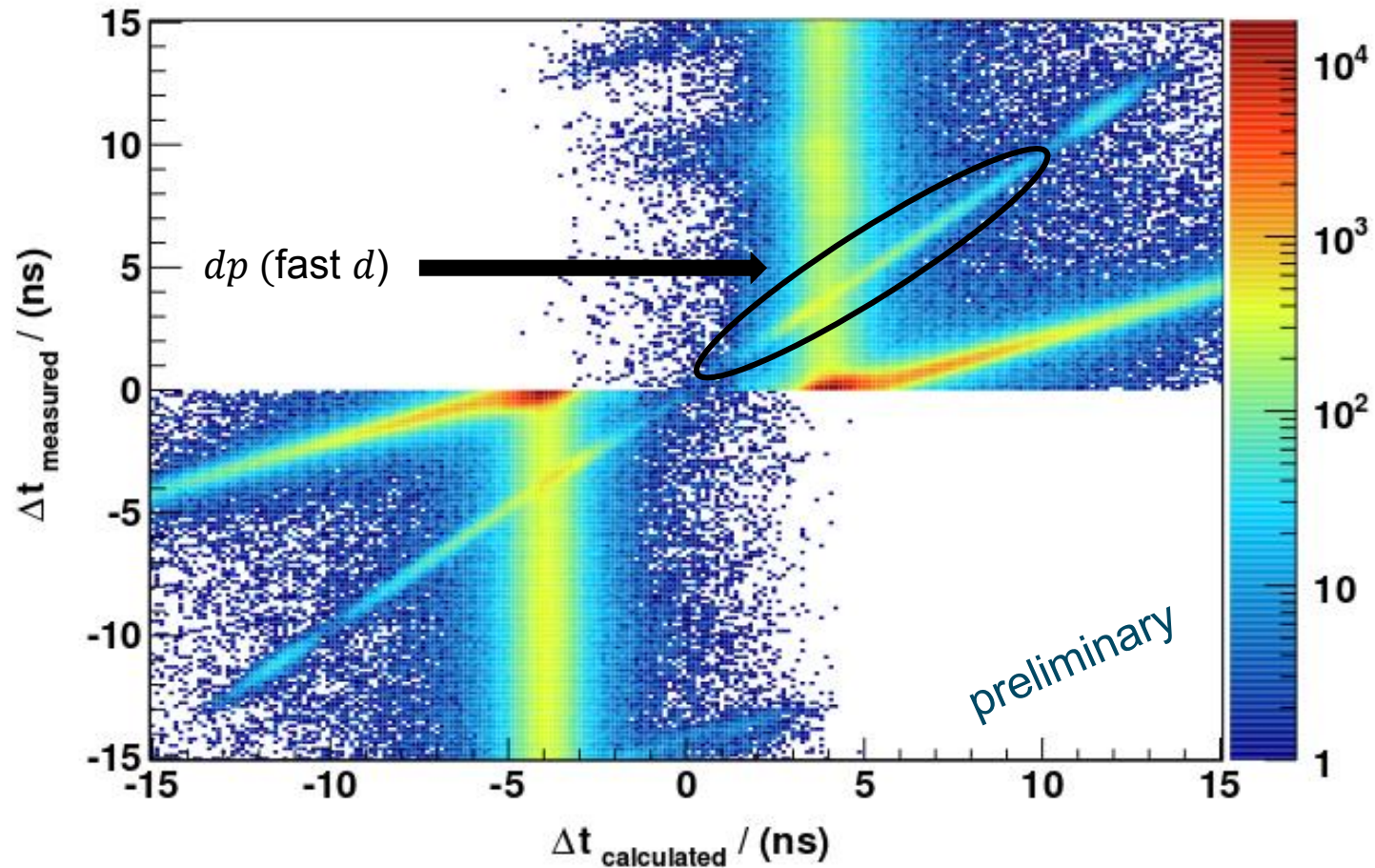


# Identification of $dp_{spec}$ pairs

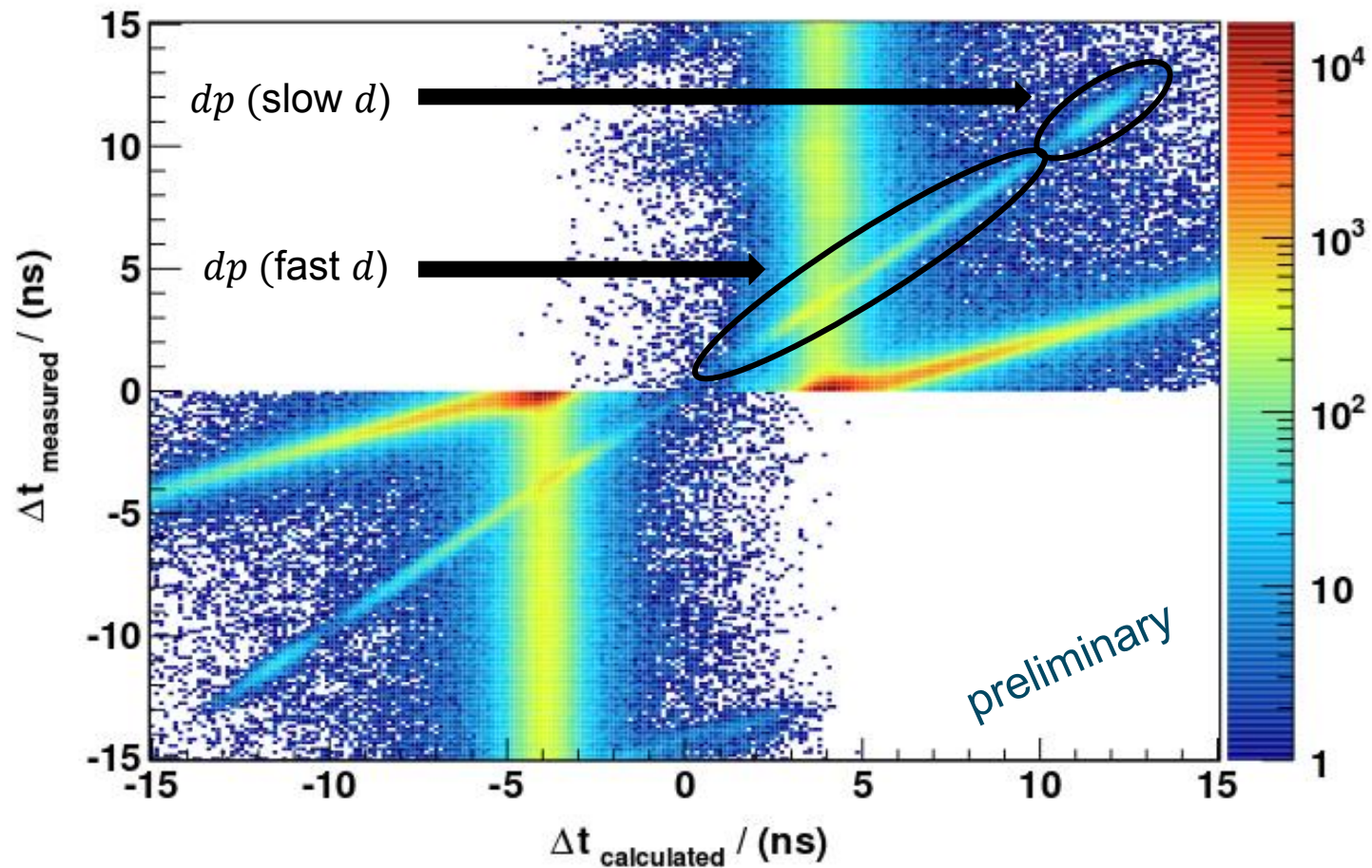




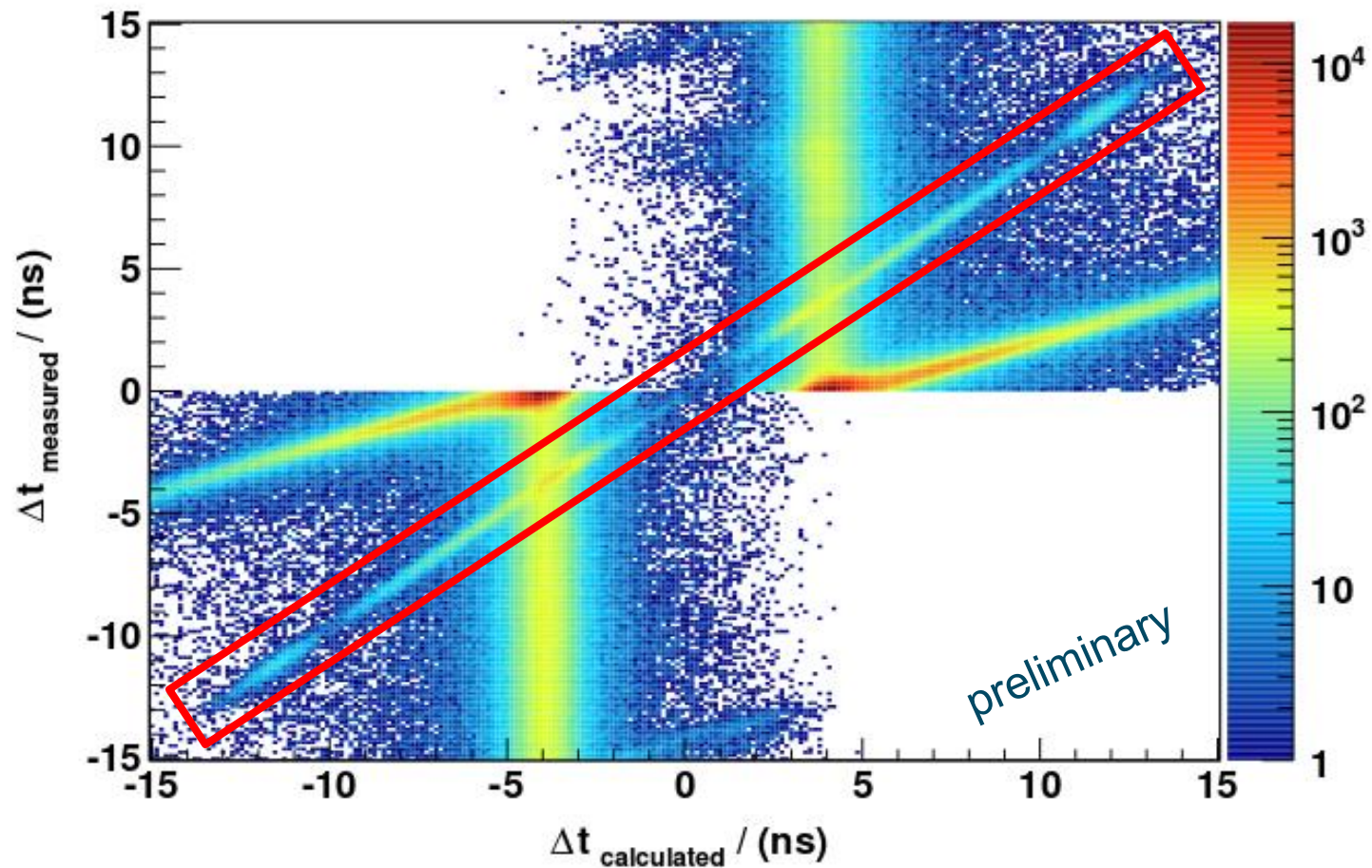
# Identification of $dp_{spec}$ pairs



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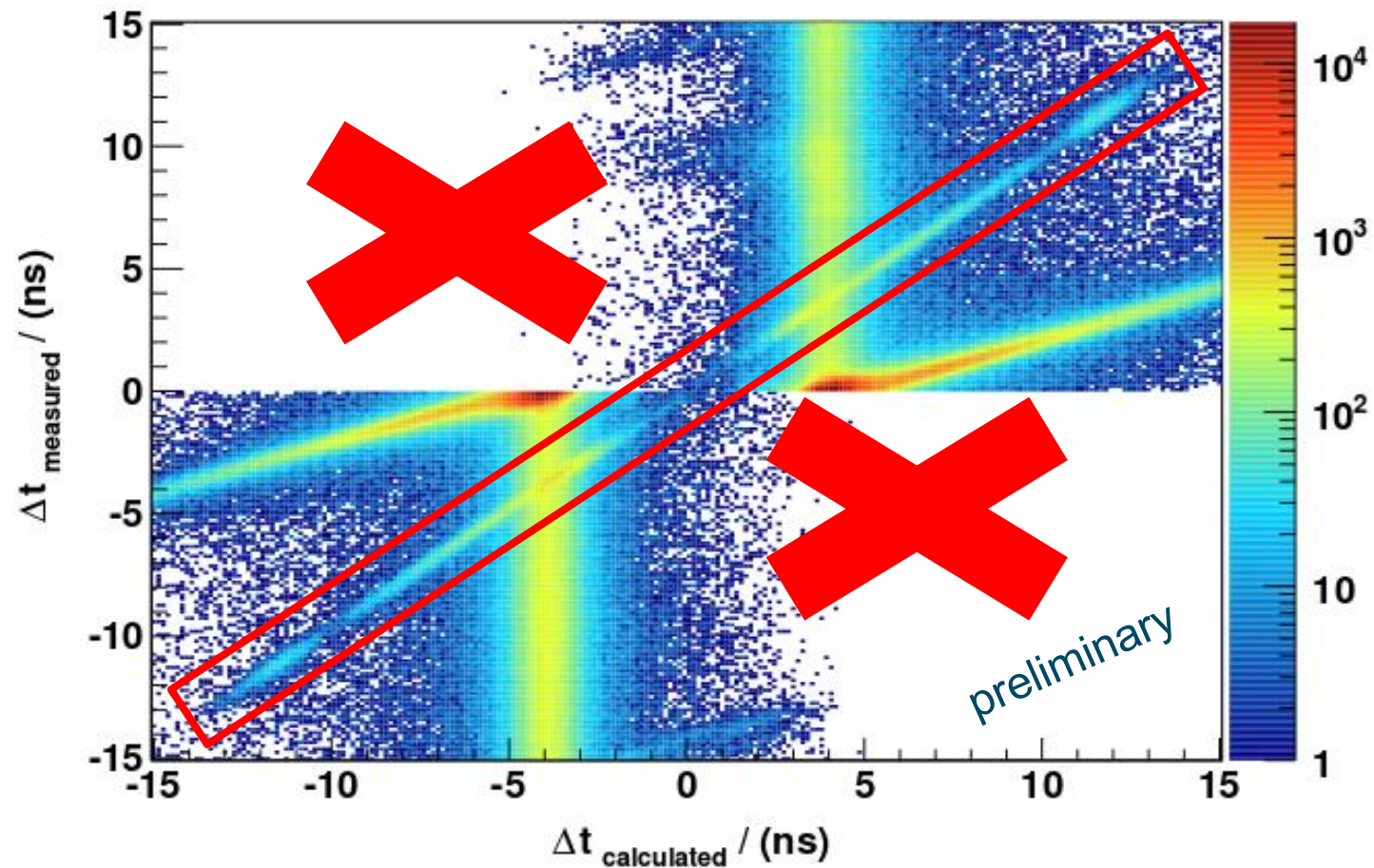


# Identification of $dp_{spec}$ pairs



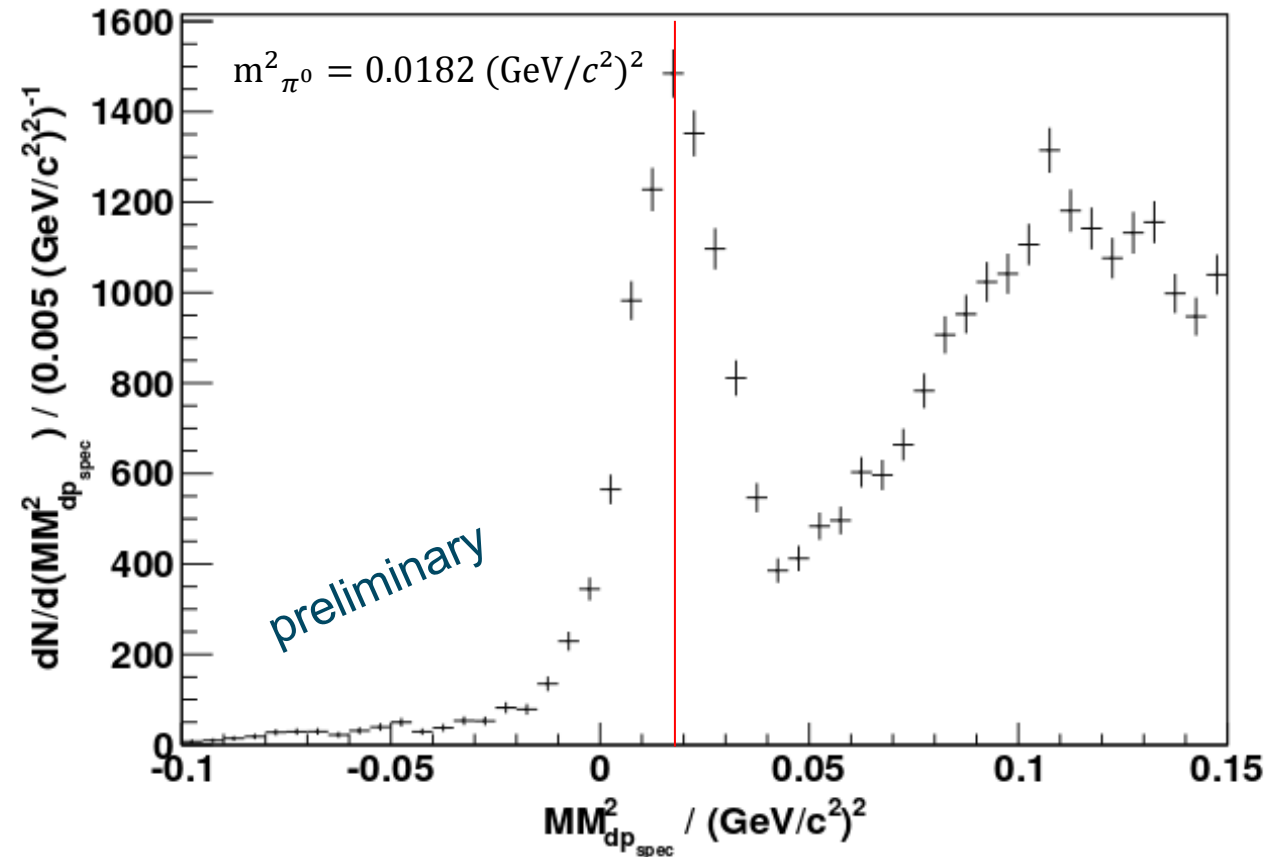
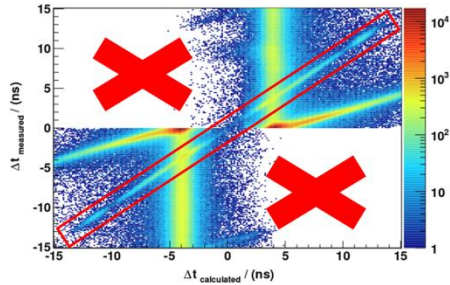


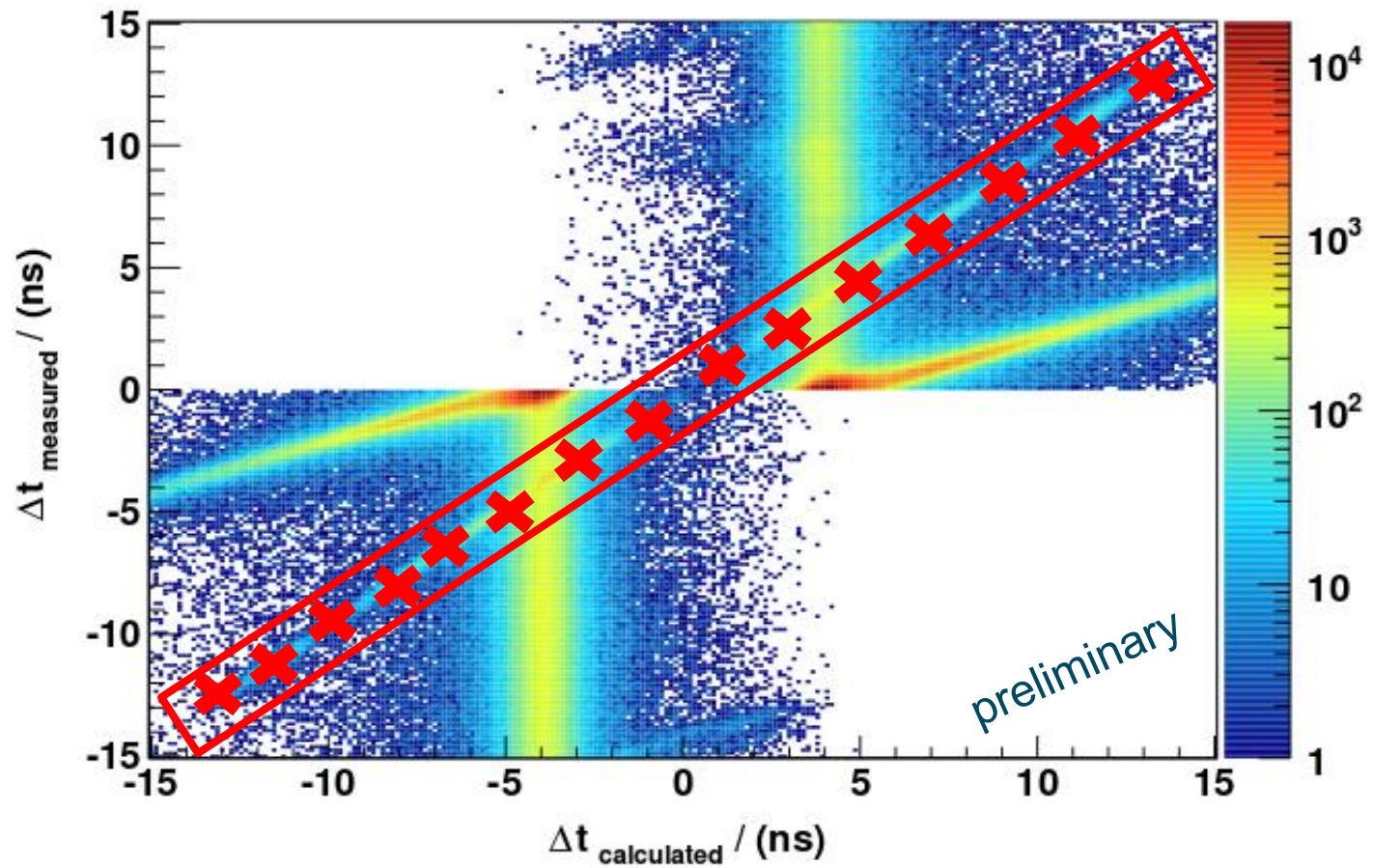
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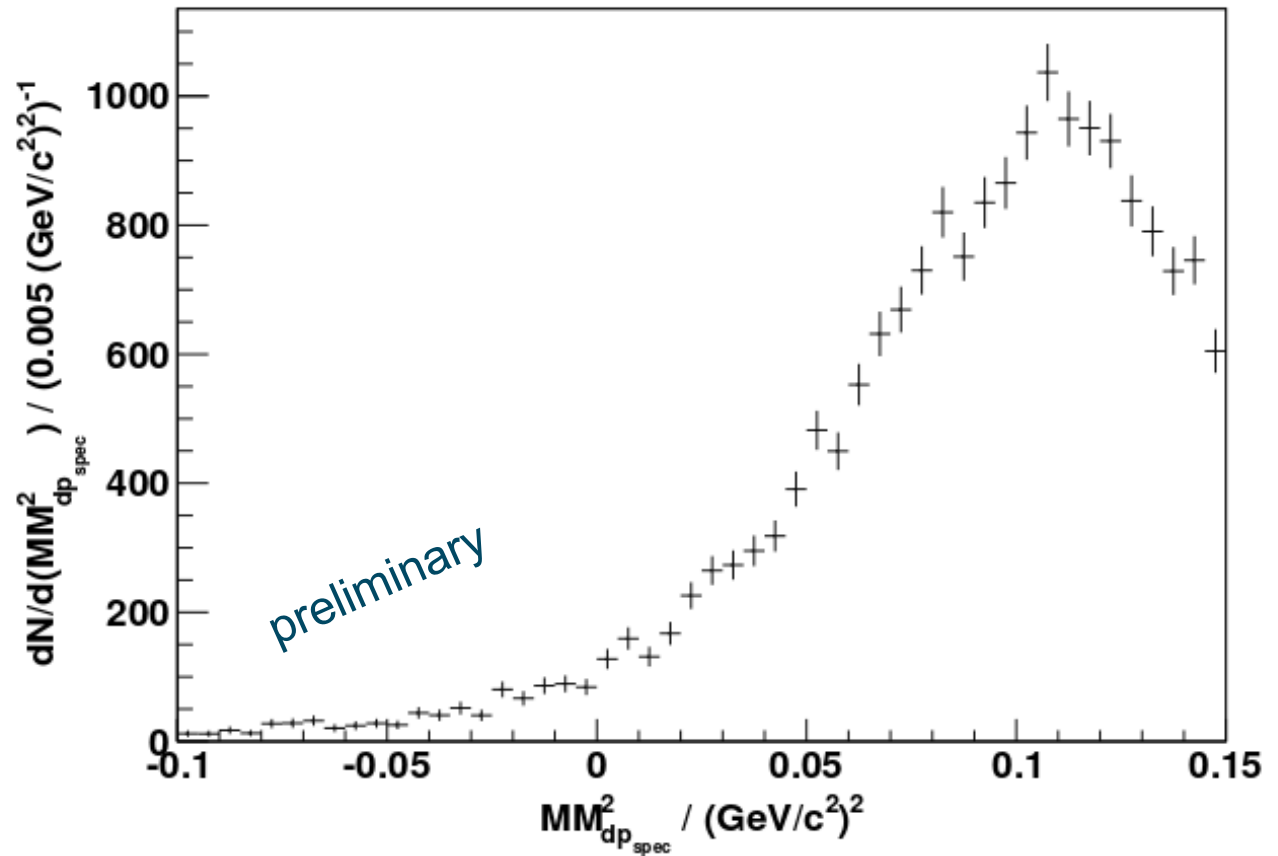
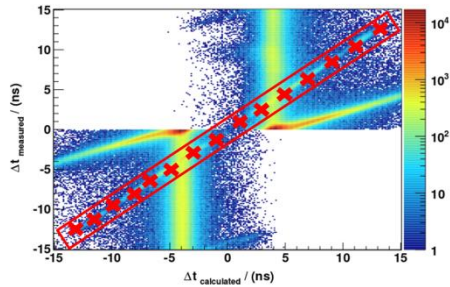


# Identification of $dp_{spec}$ pairs

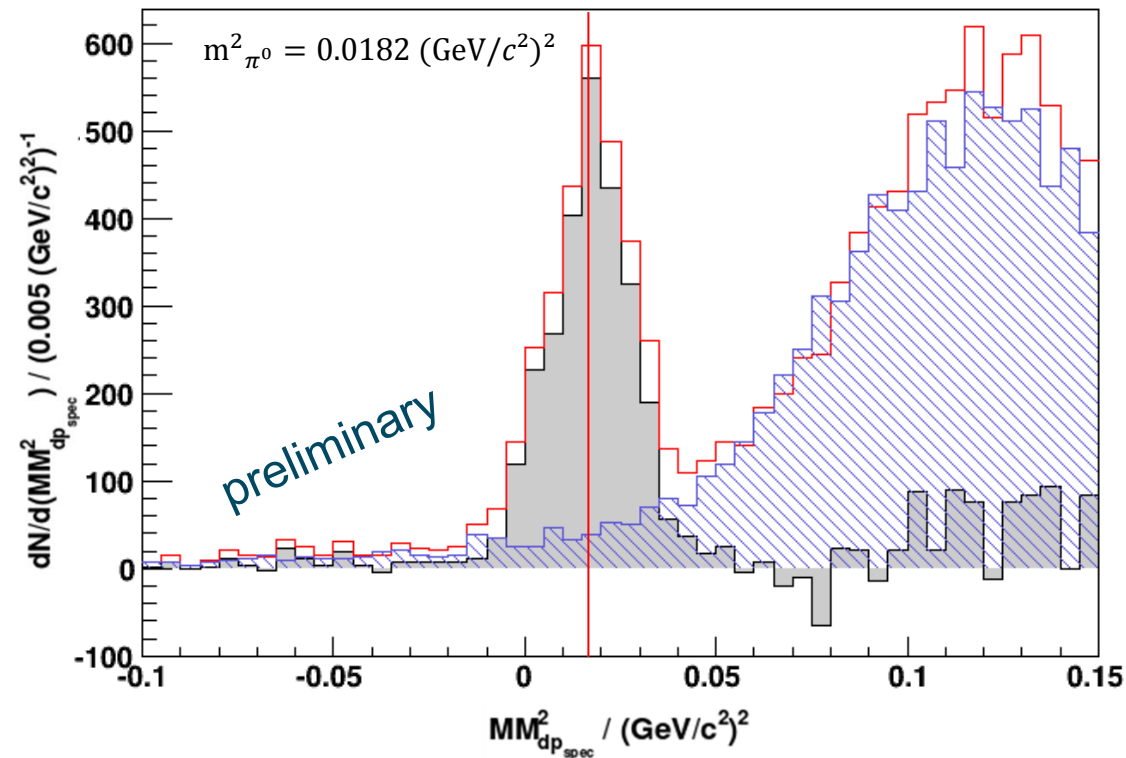




# Background description



- accurate signal background subtraction
- next steps:
  - determination of correction terms
    - acceptance
    - dead time
    - shadow effect
    - ...
- calculation of uncertainties
- luminosity determination for every beam momentum
- compare results with  $dp$ -elastic scattering

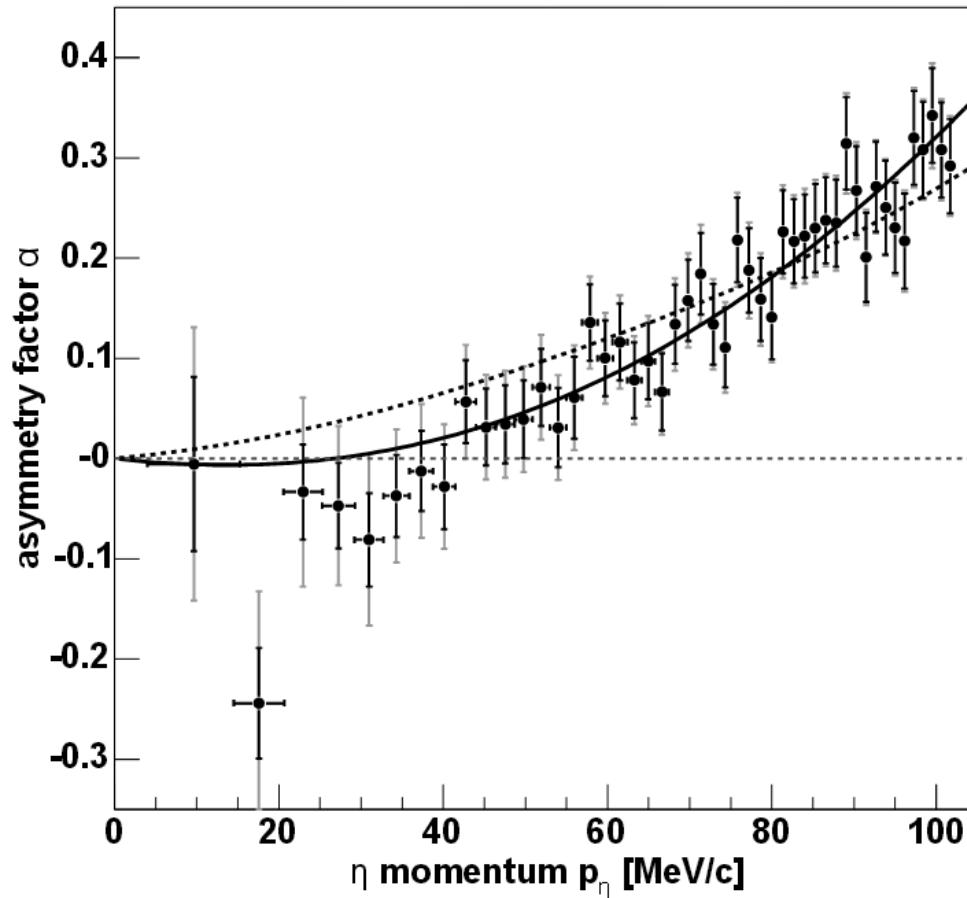


- luminosities were determined via  $dp$ -elastic scattering for 18 momentum transfer bins for each of the 19 beam momenta
- precision of  $\Delta L_{\text{stat}} = 1\%$  and  $\Delta L_{\text{sys}} = 6\%$  achieved
  - ⇒ improvement by at least a factor of two compared to previous calculations
- luminosities already used to determine differential and double differential cross sections for  $dp \rightarrow {}^3\text{He}\pi^+\pi^-$  M. Mielke et al., European Journal of Physics J. A50 (2014) 102
- relative normalization via spectator protons of reactions of the type  $dp \rightarrow p_{\text{spec}}X$
- ongoing analysis:
  - independent luminosity determination using the reaction  $dp \rightarrow dp_{\text{spec}}\pi^0$
  - identification of  $dp_{\text{spec}}$  pairs and background description by time-of-flight measurements
  - next steps:
    - determination of correction terms
    - luminosity determination for every beam momentum





Thank you for your attention!



$$\alpha = \frac{d}{d(\cos \theta_\eta)} \ln \left( \frac{d\sigma}{d\Omega} \right) \Big|_{\cos \theta_\eta = 0}$$